## Breakdown of Random-Matrix Universality in Persistent Lotka-Volterra Communities

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The eigenvalue spectrum of a random matrix often only depends on the first and second moments of its elements, but not on the specific distribution from which they are drawn. The validity of this universality principle is often assumed without proof in applications. In this Letter, we offer a pertinent counterexample in the context of the generalized Lotka-Volterra equations. Using dynamic mean-field theory, we derive the statistics of the interactions between species in an evolved ecological community. We then show that the full statistics of these interactions, beyond those of a Gaussian ensemble, are required to correctly predict the eigenvalue spectrum and therefore stability. Consequently, the universality principle fails in this system. We thus show that the eigenvalue spectra of random matrices can be used to deduce the stability of "feasible" ecological communities, but only if the emergent non-Gaussian statistics of the interactions between species.

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The theory of disordered systems enables one to deduce the behavior of collections of many interacting constituents, whose interactions are assumed to be random, but fixed in time [1]. A related discipline, random matrix theory (RMT), is concerned with the eigenvalue spectra of matrices with entries drawn from a joint probability distribution. Both fields have found numerous applications in physics [2,3] (the study of spin glasses in particular [1]), and in other disciplines such as neural networks [4–9], economics [10,11], and theoretical ecology [12–19].

It is frequently assumed that the distribution of the randomness in RMT or disordered systems is Gaussian, possibly with correlations between different interaction coefficients or matrix entries. Reasons cited for this assumption include analytical convenience, maximumentropy arguments, and the observation that higher-order moments often do not contribute to the results of calculations [1,20,21].

In random matrix theory, this latter observation is referred to as the principle of *universality* [22–24]. The principle states that results obtained for the spectra of Gaussian random matrices frequently also apply to matrix ensembles with non-Gaussian distributions. The conditions for universality to apply are usually mild (higher-order moments of the distribution must fall off sufficiently quickly with the matrix size [22,23]), and it is often tacitly assumed that these conditions will hold.

In this Letter, we offer a pertinent counterexample to the universality principle in RMT. We focus on the ecological community resulting from the dynamics of the generalized Lotka-Volterra equations with random interaction coefficients. The stability of this community is governed by the interactions between species that survive in the long run [25,26]. This is a submatrix of the original interactions, which we will refer to as the "reduced interaction matrix."

Firstly, using dynamic mean-field theory [27], we obtain the statistics of the elements in the reduced interaction matrix. These turn out to be non-Gaussian (even when the original interaction matrix is Gaussian). Secondly, we analytically calculate the leading eigenvalue of this non-Gaussian ensemble of random matrices. We show that this eigenvalue is different from the one that we would obtain from a Gaussian ensemble with the same first and second moments as in the reduced interaction matrix. This demonstrates that the principle of universality fails, and it indicates that the Gaussian assumption should not be made lightly.

Our findings have relevance to the random matrix approach to ecosystem stability, introduced by Robert May [12,13]. This approach assumes a random interaction structure between species in the community. One line of criticism of May's model is the observation that such interactions do not necessarily describe a feasible equilibrium (that is, an equilibrium for which all species' abundances are positive) [25,28–31]. The community of surviving species in the generalized Lotka-Volterra model on the other hand is feasible by construction, and we derive the statistics of the emergent random matrix ensemble that describes this community [26,32-34]. From this ensemble, we then recover the stability criteria that have previously been derived from the dynamic Lotka-Volterra model [16,35]. We thus show that one can construct a random matrix ensemble (in the sense of May) that correctly reflects the stability of a feasible community of coexistent species. This ensemble is

non-Gaussian and quite intricate. In May's words, our work contributes to "elucidating the devious strategies of nature which make for stability in enduring natural systems" [36].

We start from the generalized Lotka-Volterra equations (GLVEs) [16,35,37,38]

$$\dot{x}_i = x_i \left( 1 - x_i + \sum_{ij} a_{ij} x_j \right),\tag{1}$$

where the  $x_i \ge 0$  describe the abundances of species i = 1, ..., N. The interaction matrix elements in Eq. (1)  $a_{ij}$  are quenched random variables. We refer to these as the "original interaction matrix" elements. We assume that the mean of each matrix element is  $\overline{a_{ij}} = \mu/N$  (we use an overbar to denote averages over the ensemble of interaction matrices), and that they have variance  $\operatorname{Var}(a_{ij}) = \sigma^2/N$ . We also allow for correlations between diagonally opposed matrix elements,  $\operatorname{Corr}(a_{ij}, a_{ji}) = \Gamma$ ,  $(-1 \le \Gamma \le 1)$  where  $\operatorname{Corr}(a, b) = (\overline{ab} - \overline{a} \overline{b})/\sqrt{\operatorname{Var}(a)\operatorname{Var}(b)}$ .

The scaling with *N* of the moments of  $a_{ij}$  follows the standard conventions in disordered systems [1] and guarantees a well-defined thermodynamic limit  $N \rightarrow \infty$ . All our results are independent of the higher moments of  $a_{ij}$  as long as these moments decay sufficiently quickly with *N*. Further details can be found in Sec. S1 of the Supplemental Material [39].

Previous analyses of this system [16,35] in the thermodynamic limit have shown that there is a range of parameter combinations  $\mu$ ,  $\sigma^2$ , and  $\Gamma$  for which the dynamics reaches a unique stable fixed point, independently of the starting conditions. This is the case in the region to the left and below the instability lines in the phase diagram in Fig. 1.



FIG. 1. Stability diagram [16,35] of the GLVE system in the plane spanned by  $\mu$  and  $\sigma^2$  for fixed values of the correlation parameter  $\Gamma$ . Solid lines indicate the  $M \rightarrow \infty$  transition, dashed horizontal lines the linear instability. These lines were produced using Eqs. (S22) and (S28) in the Supplemental Material [39], respectively. Vertical lines mark the values of  $\mu$  used in the two panels of Fig. 3. The system has a unique stable fixed point below the dashed lines and to the left of the solid lines.

When a fixed-point solution is reached, not all species survive, i.e., there are some species for which  $x_i^* > 0$  and others with  $x_i^* = 0$  (we use an asterisk to denote the fixed point). Using dynamic mean-field theory (DMFT) [15,40–42], one can deduce these statistics of the species' abundances at the fixed point.

From the DMFT analysis, one can also find the combinations of system parameters at which the system is no longer able to support a unique stable fixed point. There are two types of transitions: (1) the average species abundance can diverge [i.e.,  $M \rightarrow \infty$ ], or (2) the fixed-point solution can become linearly unstable to perturbations. Closed-form expressions for the critical sets of parameters ( $\sigma$ ,  $\Gamma$ , and  $\mu$ ) at which each of these transitions occurs were derived in Refs. [16,35]. A selection of phase lines for different values of the correlation parameter  $\Gamma$  are shown in Fig. 1.

We now examine an alternative approach to analyzing the stability of the GLVEs in Eq. (1). Namely, we consider the reduced interaction matrix (the interaction matrix between the species in the surviving subcommunity). More precisely, this is defined by

$$a_{ij}' = a_{ij} - \delta_{ij},\tag{2}$$

where  $i, j \in S$  (with S the set of surviving species), and where the shift in the diagonal elements reflects the  $-x_i$ term inside the brackets of Eq. (1). It can be shown that a fixed point of the GLVEs is stable if and only if all of the eigenvalues of the reduced interaction matrix have negative real parts [17,25,26] (see also Sec. S2 in the Supplemental Material).

We note that the statistics of the reduced interaction matrix elements are determined by the extinction dynamics in the GLVE system, and are consequently vastly different from those of the original interaction matrix [33,43]. For instance, they are non-Gaussian (even when the  $a_{ij}$  are Gaussian), and there are correlations between elements sharing only one index (see the Supplemental Material [39], Sec. S6). This makes the calculation of the eigenvalue spectrum of the reduced interaction matrix a nontrivial task.

As is illustrated in Fig. 2, the spectrum of the reduced interaction matrix consists of a bulk set of eigenvalues and a single outlier. Writing  $z_{ij} = a_{ij} - \mu N^{-1}$  (where once again  $i, j \in S$ ), both the outlier eigenvalue  $\lambda_{\text{outlier}}$  and the bulk spectral density  $\rho_{\text{bulk}}(\lambda)$  can be obtained from the resolvent matrix  $\underline{G} = [\omega \underline{1} - \underline{z}]^{-1}$ . The bulk density is calculated from the trace of  $\underline{G}$  via well-known relations [44]. The outlier eigenvalue in turn fulfills [45–47]

$$\mathcal{G}(1+\lambda_{\text{outlier}}) = \frac{1}{\mu\phi},\tag{3}$$

where  $\mathcal{G}(\omega) \equiv (N\phi)^{-1} \overline{\sum_{i,j \in S} G_{ij}(\omega)}$ , and where  $\phi$  is the fraction of surviving species at the fixed point.



FIG. 2. The eigenvalues of the reduced interaction matrix. Results from a computer simulation of the GLVE are shown as markers. The solid red curve and the hollow circle show the theoretical predictions for the bulk region and outlier eigenvalue in Eq. (4) and Eqs. (S71)–(S73) of the Supplemental Material [39], respectively. Two naive predictions for the outlier that do not take the full statistics of the reduced interaction matrix into account are shown as a yellow triangle ( $\lambda_0$  in the text) and an orange square ( $\lambda_1$  in the text). System parameters are  $\sigma = 1.1$ ,  $\mu = 0.9$ ,  $\Gamma = -0.5$ , and simulation data is from a single realization with N = 10000.

We first briefly discuss the bulk spectrum, for which the results do not run counter to the universality principle. We use a series expansion for a Hermitized version of the resolvent of the reduced interaction matrix. This standard approach accounts for the nonanalytic nature of the resolvent in the bulk region [48–50].

We find that the resulting series for the trace of the resolvent matrix is identical to that of a Gaussian random matrix in the limit  $N \to \infty$ . That is, we show that the higher-order statistics of the reduced interaction matrix do not contribute to this series and, therefore, that the universality principle holds for the bulk region. The only statistics of the reduced interaction matrix that contribute are  $(\sigma')^2 \equiv N_S \operatorname{Var}(a'_{ij}) = \phi \sigma^2$  and  $\Gamma' \equiv \operatorname{Corr}(a'_{ij}, a'_{ji}) = \Gamma$  where  $N_S$  is the number of surviving species (we calculate these statistics in Sec. S6 of the Supplemental Material). One obtains the familiar elliptic law [51]

$$\rho_{\text{bulk}}(\lambda) = \begin{cases} \frac{1}{\pi(\sigma')^2 [1 - (\Gamma')^2]} & \text{if } \frac{(1+x)^2}{(1+\Gamma')^2} + \frac{y^2}{(1-\Gamma')^2} < (\sigma')^2, \\ 0 & \text{otherwise,} \end{cases}$$
(4)

where  $\lambda = x + iy$ . We can show (see the Supplemental Material [39], Sec. S5C) that the bulk of the eigenvalue spectrum crossing the imaginary axis corresponds to the linear instability of the GLVEs, represented by the dashed horizontal lines in Fig. 1. This is verified in Fig. 3(a).

We now move on to the outlier eigenvalue, which is a far less trivial matter. We first discuss two candidate expressions for the outlier eigenvalue based upon calculations for Gaussian random matrix ensembles. We show that neither of these expressions is accurate, and that the universality principle fails to predict the outlier eigenvalue. We subsequently derive an accurate expression for the outlier, which we show correctly predicts stability.

Noting previous work [14,24,45,52], one might perhaps expect that  $\mu' = N_S \overline{a'_{ij}}$   $(i \neq j)$ , together with  $(\sigma')^2$  and  $\Gamma'$ would be sufficient to predict the outlier eigenvalue of the reduced interaction matrix. Using an established formula for the outlier eigenvalues of Gaussian random matrices [45,52], one then obtains  $\lambda_0 = -1 + \mu' + \Gamma' \sigma'^2 / \mu'$ .



FIG. 3. (a) Right edge of the bulk of the eigenvalue spectrum of the reduced interaction matrix versus  $\sigma^2$  for different values of the system parameter  $\Gamma$  and fixed  $\mu = -0.5$ . Markers are the result of averaging the results of ten simulations of the GLVE with N = 4000. The dashed colored lines are given by  $\lambda_{edge} = -1 + \sigma \sqrt{\phi}(1 + \Gamma)$ , and the vertical dot-dashed lines are the points where the linear instability occurs in the GLVE (see the dashed lines in Fig. 1). (b) Outlier eigenvalue of the reduced interaction matrix versus  $\sigma^2$  at fixed  $\mu = 0.6$  and for the same values of  $\Gamma$  as in (a). Markers are the result of averaging the results of ten simulations with N = 4000. The solid lines are the analytical result in Eqs. (S71)–(S73) of the Supplemental Material, and the vertical dot-dashed lines are the points where  $M \rightarrow \infty$  in the GLVE (see the solid lines in Fig. 1).

If we also include the effects of correlations between elements sharing only one index  $\gamma' = N^2 \text{Corr}(a'_{ij}, a'_{ki})$ (where  $k \neq i$ ), we arrive at (using results from our previous work [47])

$$\lambda_1 = -1 + \mu' + \frac{\mu'}{2} \left( 1 + \frac{\Gamma'}{\gamma'} \right) \left[ \sqrt{1 + \frac{4\gamma' \sigma'^2}{(\mu')^2}} - 1 \right].$$
 (5)

The approach leading to Eq. (5) takes into account all possible correlations for a Gaussian random matrix with statistical symmetry between different species. We note that correlations between elements in the same row or column also exist in the reduced interaction matrix (see the Supplemental Material [39], Sec. S6A), but these do not affect the location of the outlier [47].

If the universality principle were to apply to the reduced interaction matrix, then the Gaussian prediction  $\lambda_1$  and the true outlier eigenvalue would coincide, whether or not the elements of the reduced interaction matrix were also Gaussian distributed. As can be seen in Fig. 4,  $\lambda_1$  is a better approximation than  $\lambda_0$ , but neither expression correctly predicts the outlier.

We now take into account the full statistics of the matrix elements  $a'_{ij}$ , as we did when calculating the bulk eigenvalue spectrum, and deduce the correct expression for the outlier eigenvalue. In the region of the complex plane outside the bulk (where the outlier resides), the resolvent can be expanded as a series in  $1/\omega$  [Eq. (S36) in the Supplemental Material]. We evaluate each term in this series in terms of the statistics of species abundances, which are available to us via DMFT. This is accomplished via a generating-functional approach (Supplemental Material [39], Sec. S4).



FIG. 4. Outlier eigenvalue of the reduced interaction matrix as a function of  $\sigma^2$ , at fixed  $\mu = 0.6$ ,  $\Gamma = -0.2$ . Markers indicate the results of computer simulations (N = 1000, averaged over ten trials). The solid line is from Eqs. (S71)–(S73) of the Supplemental Material, whereas the dashed line and dot-dashed lines are the two naive predictions  $\lambda_0$  and  $\lambda_1$  (respectively) given in the text. The vertical dot-dashed line marks the point at which  $M \rightarrow \infty$  in the GVLE (see the solid lines in Fig. 1).

Using diagrammatic techniques to recognize the selfsimilarity of the resulting series [48–50,53], we arrive at a compact formula for the resolvent [Supplemental Material [39], Eq. (S69)]. Using Eq. (3), we then obtain an implicit set of equations for the outlier eigenvalue in terms of the statistics of the surviving species abundances [see Eqs. (S71)–(S73) in the Supplemental Material]. We emphasize that in finding our final expression for the outlier, no approximations have been made other than assuming the thermodynamic limit. The simulation data in Figs. 3 and 4 verify that the expression in Eqs. (S71)–(S73) accurately predicts the outlier eigenvalue.

We also demonstrate analytically (see the Supplemental Material [39], Sec. S4D) that this prediction for the outlier eigenvalue correctly predicts instability of the fixed point of the GLVE system. That is,  $\lambda_{\text{outlier}}$  crosses the imaginary axis precisely at locations in parameter space where the  $M \rightarrow \infty$  transition occurs in the GLVEs. This is also verified in Figs. 3 and 4.

We thus conclude that stability cannot be predicted from the reduced interaction matrix using Gaussian random matrix results, even if all correlations are accounted for. This indicates that the extinction dynamics leads to some more intricate structure to the interactions in the surviving community.

Advancing ideas in Refs. [26,33,54], we show in the Supplemental Material [39] (Sec. S10) how one can generate the ensemble of reduced interaction matrices "from scratch" (i.e., without running the Lotka-Volterra dynamics and eliminating extinct species). This is achieved by first drawing a set of mock abundances from the known distribution of GLVE fixed-point abundances [16,33]. Subsequently, one then draws interaction matrices from a carefully constructed distribution, which is dependent on the mock abundances. We verify in the Supplemental Material [39] that this bottom-up construction leads to non-Gaussian matrices with the same statistical properties and leading eigenvalue as the ensemble of true reduced interaction matrices.

Having constructed the reduced interaction matrix ensemble in this way, we can thus see more clearly why universality fails to capture stability. The ensemble is manifestly non-Gaussian with complex interdependencies between matrix elements. By making a simple Gaussian assumption and ignoring the higher-order moments, one does not correctly take into account this intricate underlying structure.

Finally, we perform some additional tests of our results to demonstrate their robustness. For example, realistic ecological communities might be composed of only a relatively small number of species. We have verified that our expression for the outlier in Eqs. (S71)–(S73) of the Supplemental Material [39] is also a better predictor of stability than the more naive theories when N = 50, leading to communities of about 25 surviving species (Fig. S4 in the Supplemental Material). It has also been pointed out that heterogeneity of carrying capacities across species can significantly affect ecological equilibria [33,55]. We show in Sec. S9 of the Supplemental Material [39] that our conclusions continue to hold in such situations.

To conclude, we have deduced the stability of the generalized Lotka-Volterra system by calculating the eigenvalue spectrum of the interaction matrix of the surviving species. We have shown that results that are derived for Gaussian random matrices, which are often assumed also to apply to non-Gaussian ensembles, fail in this case. Instead, higher-order statistics of the reduced interaction matrix must be taken into account. We have therefore found a noncontrived class of random matrices for which the universality principle of RMT is not applicable. This demonstrates that there are limitations to results in RMT that are derived making an assumption of Gaussian interactions. Universality should therefore not be invoked without careful consideration.

Our results also have immediate relevance for the field of theoretical ecology. In the widely used approach pioneered by Robert May [12,13], one supposes that the Jacobian governing small deviations of species abundances about a fixed point can be represented by a random matrix. May does not say what the dynamics are that lead to this Jacobian. One particular objection to this approach is hence that the statistics of May's random matrices do not necessarily correspond to "feasible" equilibria [24,25,28,30].

The fixed point of the GLVEs is feasible by construction. Therefore, our work shows that the stability of a feasible equilibrium in a complex ecosystem can be found by studying the eigenvalues of a random interaction matrix. Feasibility is reflected in the higher-order statistics of the interactions between species. Crucially, we find that these intricate statistics cannot be ignored if one is to correctly predict stability.

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 M. Mézard, G. Parisi, and M. Virasoro, *Spin Glass Theory* and Beyond: An Introduction to the Replica Method and Its Applications (World Scientific Publishing Company, London, 1987), Vol. 9.

- [2] E. P. Wigner, On the distribution of the roots of certain symmetric matrices, Ann. Math. **67**, 325 (1958).
- [3] E. P. Wigner, Random matrices in physics, SIAM Rev. 9, 1 (1967).
- [4] J. Aljadeff, M. Stern, and T. Sharpee, Transition to Chaos in Random Networks with Cell-Type-Specific Connectivity, Phys. Rev. Lett. **114**, 088101 (2015).
- [5] A. Kuczala and T. O. Sharpee, Eigenvalue spectra of large correlated random matrices, Phys. Rev. E 94, 050101(R) (2016).
- [6] A. C. Coolen, P. Sollich, and R. Kühn, *Theory of Neural Information Processing Systems* (Oxford University Press, Oxford, UK, 2005).
- [7] K. Rajan and L. F. Abbott, Eigenvalue Spectra of Random Matrices for Neural Networks, Phys. Rev. Lett. 97, 188104 (2006).
- [8] C. Louart, Z. Liao, and R. Couillet, A random matrix approach to neural networks, Ann. Appl. Probab. 28, 1190 (2018).
- [9] A. Kuczala, Dynamics and Information Processing in Recurrent Networks (University of California, San Diego, 2019).
- [10] L. Laloux, P. Cizeau, M. Potters, and J.-P. Bouchaud, Random matrix theory and financial correlations, Int. J. Theor. Appl. Finance 3, 391 (2000).
- [11] J.-P. Bouchaud and M. Potters, in *The Oxford Handbook of Random Matrix Theory* (Oxford University Press, Oxford (UK), 2008).
- [12] R. M. May, Will a large complex system be stable?, Nature (London) **238**, 413 (1972).
- [13] R. M. May, Stability in multispecies community models, Math. Biosci. 12, 59 (1971).
- [14] S. Allesina and S. Tang, Stability criteria for complex ecosystems, Nature (London) **483**, 205 (2012).
- [15] M. Opper and S. Diederich, Phase Transition and 1/f Noise in a Game Dynamical Model, Phys. Rev. Lett. 69, 1616 (1992).
- [16] T. Galla, Dynamically evolved community size and stability of random Lotka-Volterra ecosystems, Europhys. Lett. 123, 48004 (2018).
- [17] G. Biroli, G. Bunin, and C. Cammarota, Marginally stable equilibria in critical ecosystems, New J. Phys. 20, 083051 (2018).
- [18] A. Altieri, F. Roy, C. Cammarota, and G. Biroli, Properties of Equilibria and Glassy Phases of the Random Lotka-Volterra Model with Demographic Noise, Phys. Rev. Lett. 126, 258301 (2021).
- [19] J. W. Baron and T. Galla, Dispersal-induced instability in complex ecosystems, Nat. Commun. 11, 1 (2020).
- [20] S. F. Edwards and P. W. Anderson, Theory of spin glasses, J. Phys. F 5, 965 (1975).
- [21] T. Galla and J. D. Farmer, Complex dynamics in learning complicated games, Proc. Natl. Acad. Sci. U.S.A. 110, 1232 (2013).
- [22] T. Tao and V. Vu, Random matrices: Universality of local eigenvalue statistics up to the edge, Commun. Math. Phys. 298, 549 (2010).
- [23] T. Tao, V. Vu, M. Krishnapur *et al.*, Random matrices: Universality of ESDs and the circular law, Ann. Probab. 38, 2023 (2010).

- [24] S. Allesina and S. Tang, The stability–complexity relationship at age 40: A random matrix perspective, Popul. Ecol. 57, 63 (2015).
- [25] L. Stone, The feasibility and stability of large complex biological networks: A random matrix approach, Sci. Rep. 8, 8246 (2018).
- [26] M. Barbier, C. de Mazancourt, M. Loreau, and G. Bunin, Fingerprints of High-Dimensional Coexistence in Complex Ecosystems, Phys. Rev. X 11, 011009 (2021).
- [27] C. De Dominicis, Dynamics as a substitute for replicas in systems with quenched random impurities, Phys. Rev. B 18, 4913 (1978).
- [28] M. E. Gilpin, Stability of feasible predator-prey systems, Nature (London) 254, 137 (1975).
- [29] T. Namba, Multi-faceted approaches toward unravelling complex ecological networks, Popul. Ecol. 57, 3 (2015).
- [30] T. Gibbs, J. Grilli, T. Rogers, and S. Allesina, Effect of population abundances on the stability of large random ecosystems, Phys. Rev. E 98, 022410 (2018).
- [31] J. Grilli, M. Adorisio, S. Suweis, G. Barabás, J. R. Banavar, S. Allesina, and A. Maritan, Feasibility and coexistence of large ecological communities, Nat. Commun. 8, 14389 (2017).
- [32] B. Goh and L. Jennings, Feasibility and stability in randomly assembled Lotka-Volterra models, Ecol. Model. 3, 63 (1977).
- [33] G. Bunin, Interaction patterns and diversity in assembled ecological communities, arXiv:1607.04734.
- [34] C. A. Serván, J. A. Capitán, J. Grilli, K. E. Morrison, and S. Allesina, Coexistence of many species in random ecosystems, Nat. Ecol. Evol. 2, 1237 (2018).
- [35] G. Bunin, Ecological communities with Lotka-Volterra dynamics, Phys. Rev. E 95, 042414 (2017).
- [36] R. M. May, Stability and complexity in model ecosystems, in *Stability and Complexity in Model Ecosystems* (Princeton University Press, Princeton, NJ, 2019).
- [37] O. Malcai, O. Biham, P. Richmond, and S. Solomon, Theoretical analysis and simulations of the generalized Lotka-Volterra model, Phys. Rev. E 66, 031102 (2002).
- [38] L. Brenig, Complete factorisation and analytic solutions of generalized Lotka-Volterra equations, Phys. Lett. A 133, 378 (1988).
- [39] See Supplemental Material at http://link.aps.org/ supplemental/10.1103/PhysRevLett.130.137401 for further details of the DMFT calculation of the statistics and eigenvalue spectrum of the reduced interaction matrix, as

well as our method for producing "imitation" reduced interaction matrices.

- [40] H. Sompolinsky and A. Zippelius, Dynamic Theory of the Spin-Glass Phase, Phys. Rev. Lett. 47, 359 (1981).
- [41] T. R. Kirkpatrick and D. Thirumalai, *P*-spin-interaction spin-glass models: Connections with the structural glass problem, Phys. Rev. B 36, 5388 (1987).
- [42] F. Roy, G. Biroli, G. Bunin, and C. Cammarota, Numerical implementation of dynamical mean field theory for disordered systems: Application to the Lotka–Volterra model of ecosystems, J. Phys. A 52, 484001 (2019).
- [43] J. Fraboul, G. Biroli, and S. De Monte, Artificial selection of communities drives the emergence of structured interactions, arXiv:2112.06845.
- [44] H.-J. Sommers, A. Crisanti, H. Sompolinsky, and Y. Stein, Spectrum of Large Random Asymmetric Matrices, Phys. Rev. Lett. 60, 1895 (1988).
- [45] S. O'Rourke, D. Renfrew *et al.*, Low rank perturbations of large elliptic random matrices, Electron. J. Pro 19 (2014).
- [46] F. Benaych-Georges and R. R. Nadakuditi, The eigenvalues and eigenvectors of finite, low rank perturbations of large random matrices, Adv. Math. 227, 494 (2011).
- [47] J. W. Baron, T. J. Jewell, C. Ryder, and T. Galla, Eigenvalues of Random Matrices with Generalised Correlations: A Path Integral Approach, Phys. Rev. Lett. **128**, 120601 (2022).
- [48] R. A. Janik, M. A. Nowak, G. Papp, and I. Zahed, Nonhermitian random matrix models, Nucl. Phys. B501, 603 (1997).
- [49] J. Feinberg and A. Zee, Non-hermitian random matrix theory: Method of hermitian reduction, Nucl. Phys. B504, 579 (1997).
- [50] E. Brézin and A. Zee, Correlation functions in disordered systems, Phys. Rev. E 49, 2588 (1994).
- [51] J. Ginibre, Statistical ensembles of complex, quaternion, and real matrices, J. Math. Phys. (N.Y.) 6, 440 (1965).
- [52] S. F. Edwards and R. C. Jones, The eigenvalue spectrum of a large symmetric random matrix, J. Phys. A **9**, 1595 (1976).
- [53] G. 't Hooft, A planar diagram theory for strong interactions, Nucl. Phys. **B72**, 461 (1974).
- [54] F. Vrins, Sampling the multivariate standard normal distribution under a weighted sum constraint, Risks 6, 64 (2018).
- [55] C. Song and S. Saavedra, Will a small randomly assembled community be feasible and stable?, Ecology 99, 743 (2018).