Thermodynamic, Raman Spectroscopic and UV-Visible Optical Characterization of the Deltic, Squaric and Croconic Cyclic Oxocarbon Acids

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ABSTRACT: A precise and complete thermodynamic, Raman spectroscopic and ultraviolet-visible optical characterization of the deltic, squaric and croconic cyclic oxocarbon acids is obtained using theoretical solid-state methods employing very demanding calculation parameters. The computed fundamental thermodynamic properties include the isobaric specific heat, the entropy, the enthalpy and the Gibbs free energy as a function of temperature. The calculated specific heats at 298.15 K of the deltic, squaric and croconic acids are 89.7, 111.2 and 133.2 J \cdot mol⁻¹ \cdot K⁻¹, respectively, and the corresponding entropies are 98.3, 117.3 and 136.5 J \cdot $mol^{-1} \cdot K^{-1}$. The only value of these properties known from experimental measurements is the specific heat of the squaric acid which differs from the computed value at 315 K by about 4.9%. The calculated values of the thermodynamic properties are then used to determine the thermodynamic properties of formation of these materials in terms of the elements. As an application of the calculated thermodynamic properties of formation, the Gibbs free energies of reaction and associated reaction constants are evaluated for the reactions of thermal decomposition and complete combustion of the squaric and croconic acids and the reaction of interconversion between them. The only available experimental values of these properties, namely the enthalpies of combustion of squaric and croconic acids at room temperature, are reproduced theoretically with high accuracy. The Raman spectra of these materials are also computed using Density Functional Perturbation Theory. The analysis of the theoretical Raman spectra of these materials points out to significant differences with respect to their usual empirical assignment. Therefore, the Raman spectra of these materials is fully reassigned. Finally, the ultraviolet-visible (UV-Vis) optical properties of the deltic, squaric and croconic acids are computed. The UV-Vis absorption spectrum of the croconic acid in the spectral region 225-425 nm and the UV absorption spectrum of the squaric acid in the region 200-350 nm which had previously been measured experimentally are well reproduced. The corresponding spectrum for the deltic acid and the reflectivity, optical conductivity, dielectric, refractive index and loss optical functions of the three materials, which had never been published as far as we know, are reported as a function of the wavelength of incident radiation in the range 200-750 nm. The origin of the peaks in the absorption spectra which had not been analyzed so far is unveiled here by examining the inter-band electronic transitions in these materials.

I. INTRODUCTION

The cyclic oxocarbon acids¹⁻⁶ ($C_n O_n H_2$, n = 3, 4, 5, 6) have long been studied in the last decades and have been employed in a wide range of scientific, technological and biomedical applications. A review of the literature has been reported in a previous work.⁶ However, the analysis of the literature revealed the existence of several significant pieces of knowledge concerning these materials which are almost completely unknown. The first one, the mechanical behavior of these materials, was addressed in the previous work.⁶ The second one is their fundamental thermodynamic properties and the third one is their ultraviolet-visible (UV-Vis) optical spectra and properties. Finally, although the Raman spectra of these materials have been recorded experimentally in many previous works,⁷⁻²³ the band assignments were founded on empirical arguments only. The results obtained in the present theoretical solid-state calculations of the deltic, squaric and croconic acids (n = 3, 4 and 5, respectively), pointed out to significant problems in their assignment. Therefore, the Raman spectra of these materials were reassigned. Of course, this feature is not a critique to the excellent spectroscopic works performed so far⁷⁻²³ for these materials but to the empirical method of assignment which is too rough to provide the correct spectral interpretations in many occasions.²⁴⁻ ²⁵ The extrapolation of spectroscopic data from gas phase mol-

ecule to solid state materials is a very crude approximation indeed. Likewise, the use of normal-coordinate treatments based on approximate force field methods^{8,10,11,14,15,21} may lead to significant errors in the calculated spectra of solid materials and their assignment.

The fundamental thermodynamic properties computed in this work include the isobaric specific heat, the entropy, the enthalpy and the Gibbs free energy as well as their temperature dependence. The results obtained for the thermodynamic properties where then used to determine the thermodynamic properties of formation of these materials in terms of the elements. Despite their importance, the majority of the thermodynamic properties of these materials and their temperature dependence are unknown.²⁶⁻³³ In a subsequent step, the calculated thermodynamic properties of formation were used to evaluate the Gibbs free energies of reaction and associated reaction constants for the reactions of thermal decomposition of the squaric and croconic acids:

 $\begin{array}{l} (A) \ H_2 C_4 O_4 (cr) \rightarrow 2 \ CO(g) + 1/2 \ CO_2 (g) + 3/2 \ C(cr) + \ H_2 O(l) \\ (B) \ H_2 C_5 O_5 (cr) \rightarrow 3 \ CO(g) + 1/2 \ CO_2 (g) + 3/2 \ C(cr) + \ H_2 O(l) \end{array}$

their reactions of complete combustion:

 $\begin{array}{l} (C) \ H_2 C_4 0_4 (cr) + 5/2 \ 0_2 (g) \rightarrow 4 \ C 0_2 (g) + H_2 0 (l) \\ (D) \ H_2 C_5 0_5 (cr) + 3 \ 0_2 (g) \rightarrow 5 \ C 0_2 (g) + H_2 0 (l) \end{array}$

and the reaction of conversion between them in the presence of carbon monoxide:

(E) $H_2C_4O_4(cr) + CO(g) \rightarrow H_2C_5O_5(cr)$

as a function of temperature. The thermodynamic parameters of organic compounds are extraordinarily important not only in

several scientific fields such as organic reactivity but also for many specific applications as the development of novel energetic materials.³⁴⁻³⁵ Croconic acid, in particular, has been found to be useful for the next generation of gun propellants which substantially increase the muzzle velocity.³⁴⁻³⁵ The oxocarbon acids are very relevant materials in organic reactivity studies.³⁶

The good results obtained for the structural properties of these materials in the previous work⁶ encouraged the determination of their mechanical properties. These acids in the solid state were shown to display negative Poisson ratios (NPR).³⁷⁻³⁹ Besides, the croconic acid also exhibits the phenomenon of negative linear compressibility (NLC)⁴⁰⁻⁴² for applied pressures larger than 0.4 GPa directed along the direction of minimum Poisson ratio. The materials displaying the NPR and NLC phenomena have an enormous range of applications owing to their exclusive mechanical properties.⁶ These facts emphasize the need of determining the corresponding thermodynamic data of the cyclic oxocarbon acids required to optimize its production and to investigate their chemical behavior in the presence of different chemical reagents.

While the UV-Vis absorption spectrum of the croconic acid in the spectral region 225-425 nm and the UV absorption spectrum of the squaric acid in the region 200-350 nm have been measured experimentally,⁴³⁻⁴⁴ the corresponding spectrum for the deltic acid and the reflectivity, optical conductivity, dielectric, refractive index and loss optical functions⁴⁵⁻⁴⁷ of the three materials are unknown. The great relevance of the small organic acids in the description of global atmospheric chemistry, the influence of their presence in the optical characteristics of the atmospheres⁴³ and the use of compounds closely related to the cyclic oxocarbon acids, as squaraines and croconines, as dyes⁴⁸⁻⁵³ which may be employed as optical sensors underlines the im-

portance of the precise knowledge of their optical properties. For example, it has been proposed that the croconic acid is a contaminant in the sulfuric acid cloud droplets in the Venus clouds and may represent the absorber responsible for the ultraviolet cloud features and the pale-yellow color of the Venus clouds.⁴³ The origin of the peaks in the absorption spectra of the deltic, squaric and croconic acids which has not been analyzed so far is unveiled here by examining the inter-band electronic transitions in these materials.

The main purpose of this study is to complement the relatively poor current experimental characterization of these important materials. This paper is organized as follows. In Section II, the theoretical solid-state methods used in this work are described. Then, in Section III, the main results of this work are given and discussed. This section contains three main Subsections in which the thermodynamic, optical and spectroscopic results are studied separately. Finally, in Section IV, the conclusions of this work are presented.

II. THEORY

II.1 Crystal structures and Raman Spectra

Periodic Density Functional Theory (DFT) methods based on plane wave basis sets and pseudopotentials⁵⁴ were employed in order to model the crystal structures of the cyclic oxocarbon acids.6 Thus, the generalized gradient approximation (GGA) together with Perdew-Burke-Ernzerhof functional⁵⁵ supplemented with Grimme empirical dispersion correction,⁵⁶ was used. The computations were carried out with very exigent computational parameters in order to obtain very precise crystal structures and values of the thermodynamic, optical and Raman spectroscopic properties of these materials. The technical details of the corresponding computational treatments may be found in the previous paper.⁶ The optimized crystal structures of deltic, squaric, and croconic acids were very accurate and the computed X-ray powder diffraction patterns were in excellent agreement with their experimental counterparts.⁶ The computed crystal structures are displayed in Figure 1. For the calculations of the Raman spectra, the linear response density functional perturbation theory (DFPT),⁵⁷⁻⁵⁹ implemented in the CASTEP code,⁶⁰ a module of Materials Studio package,⁶¹ was used in the same way as in previous works.^{24-25,62-67}The theoretical solidstate methodology used in this paper has been employed not only for organic and metal-organic materials 6.24-25,68-71 but also for a large series of uranyl containing compounds as oxides,⁷² oxyhydroxides,⁷³⁻⁷⁷ peroxides,⁶³ carbonates^{62,78-79} and silicates.⁶⁴⁻⁶⁵ In all these cases the results were very satisfactory.

II.2. Fundamental thermodynamic properties

In order to determine the thermodynamic properties of these materials, phonon calculations were carried out at their optimized structures. The phonon spectrum at different points of the Brillouin zone were obtained employing the Density Functional Perturbation Theory (DFPT) technique as second order derivatives of the total energy.⁸⁰ The phonon dispersion curves and density of states were computed from the phonon spectra and, using these results, several important thermodynamic properties in the quasi-harmonic approximation, namely the Gibbs free energies, enthalpies, entropies, and specific heats, were evaluated.^{80,24-25,72,78} The thermodynamic properties of formation and reaction will be evaluated in terms of the calculated fundamental thermodynamic properties using the methods described in the next subsections.^{72-73,81}

II.3. Thermodynamic properties of formation

The enthalpies and Gibbs free-energies of formation in terms of the elements were determined, from the computed enthalpy and entropy functions, $(H_T - H_{298})^{calc}$ and S_T^{calc} , by means of the relationships:^{82,72}



Figure 1. Calculated crystal structures of (A) deltic, (B) squaric and (C) croconic acids: Views of the corresponding unit cells from [010] crystallographic direction. Color code: Grey-C; Red-O; White-H.

$$\Delta_{f}H(T) = \Delta_{f}H^{0} + (H_{T} - H_{298})^{calc} + \sum_{i}^{elements} n_{i}(H_{T} - H_{298})_{i}^{exp}$$
(1)

$$\Delta_f G(T) = \Delta_f H(T) - T \left\{ S_T^{calc} - \sum_i^{elements} n_i \left(S_T \right)_i^{exp} \right\}$$
(2)

In these equations, $\Delta_f H^0$ is the enthalpy of formation at the standard state of the material being considered, and $(H_T (H_{298})_i^{exp}$ and $(S_T)_i^{exp}$ are the enthalpy and entropy functions of the elements forming part of this material with stoichiometric coefficients n_i , respectively. For a pure substance, the concept of standard state applies to the substance in a well-defined state of aggregation at a well-defined but arbitrarily chosen standard pressure (usually chosen as 1 bar). Therefore, the definition of standard state makes no reference to a fixed temperature. However, it is usually convenient to complete the definition of the standard state in a particular context by choosing a value for the reference temperature.⁸³ The most usually recommended value is 298.15 K. The precise values used for $\Delta_f H^0$ will be given below. The thermodynamic functions for carbon, oxygen and hydrogen were taken from JANAF (Joint Army-Navy-Air Force) thermochemical tables.⁸² The reaction constants of the formation reactions were calculated from the corresponding Gibbs free energies of formation employing the relationship:⁸²

$$\Delta G(T) = -R T Ln K \tag{3}$$

II.4. Thermodynamic properties of reaction

The enthalpies and Gibbs free energies of reaction at different temperatures were obtained, using our computed values of the Gibbs free energies and entropies of formation functions, $\Delta_f G(T)^{calc}$ and $\Delta_f S(T)^{calc}$, by means of the expressions:^{82,73}

$$\Delta_r G(T) = \sum_i^{\text{products}} n_i \Delta_f G^i(T) - \sum_j^{\text{reactants}} n_j \Delta_f G^j(T) \qquad (4)$$

$$\Delta_r H(T) = \Delta_r G(T) + T \cdot \Delta_r S(T)$$
(5)

where,

$$\Delta_r S(T) = \sum_i^{\text{products}} n_i \Delta_f S^i(T) - \sum_j^{\text{reactants}} n_j \Delta_f S^j(T)$$
(6)

In these relations, $\Delta_f G^i(T)$ and $\Delta_f S^i(T)$, are the Gibbs free energy and entropy of formation at temperature T of compound *i* participating in the reaction with stoichiometric coefficient n_i . The thermodynamic functions for CO(g), CO₂(g), H₂O(l), O₂(g) and C(cr) were taken from JANAF thermochemical tables.⁸² The reaction constants were calculated from the Gibbs free energies of reaction by using equation (3).

II.5. Optical properties

The propagation of an electromagnetic wave through a certain material can be described by a complex refractive index,⁴⁷⁻⁴⁹ $N = n + i \cdot k$. For propagation in vacuum N is real and equal to unity. The imaginary part k is related to the absorption coefficient α by $\alpha = 2 k\omega/c$, where ω is the frequency of the electromagnetic wave and c is the speed of light. The absorption coefficient gives the fraction of wave energy lost when it passes through the material. For normal incidence of the radiation wave on the material surface, the radiation intensity at the distance x from the surface is $I(x) = I(0) \cdot \exp(-2\pi\alpha x)$, where I(0) is the initial intensity of the incident light. For normal incidence of the radiation is related to the refractive index:

$$R = \left|\frac{1-N}{1+N}\right|^2 = \frac{(n-1)^2 + k^2}{(n+1)^2 + k^2} \tag{7}$$

Theoretically, to obtain the optical properties of a material, the complex dielectric constant is obtained in the first place and then the remaining optical properties are determined from it. The complex dielectric constant ε is related to the refractive index by $\varepsilon = \varepsilon_1 + i \cdot \varepsilon_2 = N^2$. Hence,

$$\varepsilon_1 = n^2 - k^2 ; \ \varepsilon_2 = 2 \ n \ k \tag{8}$$

The optical properties are also usually expressed in terms of the optical conductivity σ

$$\sigma_1 = \sigma_1 + i \cdot \sigma_2 = i \frac{\omega}{4\pi} (\varepsilon - 1) \tag{9}$$

The loss function ξ describing the energy lost by an electron crossing the considered material is expressed in terms of the dielectric function as

$$\xi = Im\left(-\frac{1}{\varepsilon}\right) \tag{10}$$

The interaction of a photon with a material may be described as a time dependent perturbation of the ground state electronic state. The approach implemented in CASTEP with this purpose is the time dependent density functional theory (TD-DFT).⁸⁴ The main approximation performed in this approach is the neglect of local field effects, arising from the fact that the electric field experienced at a particular site in the material is screened by the polarizability of the system itself. While these effects might also be accounted for, the corresponding computations could become prohibitively expensive.

III. RESULTS AND DISCUSION

III.1. Thermodynamic characterization

III.1.1. Fundamental thermodynamic properties

The calculated isobaric specific heats (C_p) , entropies (S), enthalpies (H) and Gibbs free energies (G) of the deltic acid are displayed in Figures 2.A, 2.B, 2.C and 2.C, respectively. Similarly, the computed thermodynamic properties of the squaric and croconic acids are shown in Figures 3.A-3.D and 4.A-4.D, respectively. The values of enthalpy and Gibbs free energy were divided by the temperature to express these properties using the same units as for C_p and S (J · K⁻¹ · mol⁻¹). The precise values of the computed thermodynamic properties over the temperature range 0-1000 K for the deltic and croconic acids and 0-350 K for the squaric acid (this material undergoes an antiferroelectric phase transition near 373 K)⁸⁵⁻⁸⁸ are provided in Tables S.1 to S.12 of the Supporting Information.

Since the most part of the thermodynamic properties of these materials have never been measured experimentally, their values were predicted. The calculated specific heats and entropies at 298.15 K are reported in Table 1. The calculated specific heats at 298.15 K of the deltic, squaric and croconic acids are 89.7, 111.2 and 133.2 $J \cdot K^{-1} \cdot mol^{-1}$, respectively, and the corresponding entropies are 98.3, 117.3 and 136.5 J · K⁻¹ · mol⁻¹. The only thermodynamic property which is known from experiment is the specific heat of squaric acid.89-90,30 The experimental and computed values at 315 K ($C_p^{exp} = 121.6 \text{ J} \cdot \text{K}^{-1} \cdot \text{mol}^{-1}$ and $C_p^{calc} = 115.6 \text{ J} \cdot \text{K}^{-1} \cdot \text{mol}^{-1}$, respectively) agree satisfactorily, the values differing by about 4.9%. The calculated specific heats of deltic and croconic acids at the highest temperature considered in the present work (1000 K), 166.6 and 249.6 J \cdot K⁻¹ \cdot mol⁻¹, respectively, are 16.5 and 16.6% below the corresponding Dulong-Petit asymptotic limits (199.5 and 299.3 $[\cdot K^{-1} \cdot mol^{-1})$.

Table 1. Calculated specific heats and entropies of the deltic, squaric and croconic acids at 298.15 K. All values are given in units of $J \cdot K^{-1} \cdot mol^{-1}$.



Figure 2. Calculated thermodynamic properties of the deltic acid: A) isobaric specific heat; B) entropy; C) enthalpy; and D) Gibbs free energy, as functions of temperature.



Figure 3. Calculated thermodynamic properties of the squaric acid: A) isobaric specific heat; B) entropy; C) enthalpy; and D) Gibbs free energy, as functions of temperature.



Figure 4. Calculated thermodynamic properties of the croconic acid: A) isobaric specific heat; B) entropy; C) enthalpy; and D) Gibbs free energy, as functions of temperature.

The expected accuracy for the computed isobaric specific heats and entropies given in Table 1 may be assessed from the analysis of the previous results obtained using the same methodology for materials in which the corresponding experimental values are available. This analysis has already been performed for uranyl containing materials.⁸¹ The expected accuracy is better than 3-5% for the computed specific heats and entropies at 298 K.81 The fundamental thermodynamic properties were also determined using the same methodology for non-uranyl containing materials as natroxalate and oxammite minerals.²⁴⁻²⁵ The results were found to be very good and to be within the above mentioned expected accuracy. The difference found in this work for the computed specific heat of squaric acid with respect to the experimental value at 315 K (4.9%) is also within this estimated accuracy. The temperature dependence of these properties is also expected to be well described. For example, the computed specific heat and entropy of uranium trioxide^{72,81} were in excellent agreement with experimental data in the full range of temperatures considered 0-1000 K. The differences were 3.9% and 1.8% at 100 K and 6.1% and 3.6% at 1000 K, respectively.

III.1.2. Thermodynamic properties of formation

The enthalpies of formation of the squaric and croconic acids at the standard state are known from experimental data.^{91,34-35} The corresponding values are $\Delta_f H^0(\text{H}_2\text{C}_4\text{O}_4(\text{cr})) = -598.2 \pm 0.4$ kJ · mol⁻¹ and $\Delta_f H^0(\text{H}_2\text{C}_5\text{O}_5(\text{cr})) = -817.41 \pm 1$ kJ · mol⁻¹.^{91,34-35} It is important to note that the value reported by Sellers⁹¹ for squaric acid, was unnoticed in the review works by Domalsky *et al.*²⁸⁻³⁰ Using these values, the computed enthalpy and entropy functions of the squaric and croconic acids and the experimental thermodynamic properties of the corresponding elements from Chase *et al.*,⁸² the enthalpies and Gibbs free energies of formation in terms of the elements were derived. The calculated values are given in Table 2. The Gibbs free energies of formation are shown in Figure 5. For the deltic acid, since the experimental value of the enthalpy of formation appears not to have been measured, the values given in Table 2 are those of $(\Delta_f H(T) - \Delta_f H^0)$ and $(\Delta_f G(T) - \Delta_f H^0)$. If the experimental value of $\Delta_f H^0(H_2C_3O_3(cr))$ is determined from accurate experimental measurements in the future, the full temperature dependent thermodynamic properties of formation of the deltic acid could be evaluated directly by adding this parameter to the values reported in Table 2.

The calculated temperature dependent thermodynamic properties of formation given in Table 2 are also expected to be accurate. In this case, the expected accuracy is about 3%. This estimate was obtained⁸¹ from the analysis of the results obtained using the same methodology for the thermodynamic functions of formation of rutherfordine, uranium trioxide and metaschoepite, the latter being a very complex uranyl containing material.⁷⁶ For these materials, the differences of the calculated and experimental values of the Gibbs free energies of formation were lower than 1% at ambient temperature and the differences remained very small at high temperatures. These differences were 1.6%, 1.0%, and 2.0% at 700, 900, and 800 K, respectively.⁸¹

III.1.3. Thermodynamic properties of the reactions of thermal decomposition of squaric and croconic acids [(A)-(B)].

Among the different possible thermal decomposition reactions, only the reactions (A) and (B) have been shown to have an stoichiometry compatible with the experimentally observed percent of residual carbon resulting from the thermal treatment of the these acids.⁹² These reactions state that under heating, one formula unit of these acids decompose into 2 (3) formula units of carbon monoxide, 1/2 of carbon dioxide, 3/2 of carbon and one of water. From the thermodynamic properties of formation in terms of the elements reported in the previous subsection III.2 and the experimental thermodynamic properties of CO(g), $CO_2(g) H_2O(l)$, and C(cr) taken from JANAF thermochemical tables,⁸² the Gibbs free energies of reaction and associated reaction constants of the reactions of thermal decomposition of

the squaric and croconic acids may be determined. The calculated values are given in Table 3 and shown in Figure 6.

III.1.4. Thermodynamic properties of the reactions of complete combustion of squaric and croconic acids [(C)-(D)]. Enthalpies of combustion

Following the same scheme as in Subsection III.1.3, the corresponding reaction parameters for reactions (C) and (D) were calculated. Thus, from the computed thermodynamic properties of formation in terms of the elements and the experimental thermodynamic properties for $CO_2(g)$, $H_2O(l)$, and $O_2(cr)$ taken from JANAF thermochemical tables,⁸² the Gibbs free energies of reaction and associated reaction constants of the reactions of complete combustion of the squaric and croconic acids, (C) and (D), were obtained. The calculated thermodynamic properties of reaction are provided in Table 4 and plotted in Figure 7.

The computed value of the combustion enthalpies at room temperature are $\Delta_r^c H^0(H_2C_4O_4(cr)) = -1261.70 \text{ kJ} \cdot \text{mol}^{-1}$ and $\Delta_r^c H^0(H_2C_5O_5(cr)) = -1436.01 \text{ kJ} \cdot \text{mol}^{-1}$ for the squaric and croconic acids, respectively. These values agree very well with the experimental values reported by Sellers,⁹¹ -1261.67±0.33

kJ · mol⁻¹ and Ciezak-Jenkins *et al.*³⁴⁻³⁵, -1436 kJ · mol⁻¹, for the squaric and croconic acids, respectively. The excellent agreement between these values gives strong support to the computed values of the thermodynamic properties of the cyclic oxocarbon acids reported in this article because the combustion enthalpies of organic compounds are generally determined experimentally with high accuracy.

III.1.5. Thermodynamic properties of the reaction of conversion between squaric and croconic acids [(E)]

Again, from the computed thermodynamic properties of formation in terms of the elements of the squaric and croconic acids and the experimental thermodynamic properties for CO(g), taken from JANAF thermochemical tables,⁸² the Gibbs free energies of reaction and associated reaction constants of the reaction (E) were calculated. These properties are listed in Table 5 and displayed in Figure 8. Since the enthalpies and free energies of this reaction are negative for the full range of temperatures considered, this reaction should be exothermic and spontaneous. This suggests a way to obtain the croconic acid from the squaric acid submitted to high carbon monoxide gas pressures.

Table 2. Calculated enthalpies ($\Delta_f H$) and free energies ($\Delta_f G$) of formation and reaction constants (Log K) for the deltic, squaric and croconic acids. The values of $\Delta_f H$ and $\Delta_f G$ are in units of kJ/mol.

TAD	Deltic acid		Squaric acid			Croconic acid		
1(K)	$\Delta_f H - \Delta_f H^0$	$\Delta_f G - \Delta_f H^0$	$\Delta_{\mathbf{f}}\mathbf{H}$	$\Delta_{\mathbf{f}}\mathbf{G}$	Log K	$\Delta_{\mathbf{f}}\mathbf{H}$	$\Delta_{\mathbf{f}}\mathbf{G}$	Log K
298.15	0.00	106.54	-598.20	-465.04	81.47	-817.41	-657.68	115.22
50	-274.98	-262.20	-933.82	-917.85	958.85	-1212.69	-1193.58	1246.89
100	-115.41	-82.05	-739.91	-698.16	364.67	-985.18	-935.11	488.44
200	-29.68	40.68	-634.92	-546.92	142.84	-861.44	-755.81	197.39
300	0.38	107.59	-597.73	-463.72	80.74	-816.85	-656.11	114.24
350	8.78	134.22	-587.27	-430.51	64.25	-804.28	-616.30	91.97
400	14.75	158.25	-	-	-	-795.34	-580.32	75.78
450	18.87	180.28	-	-	-	-789.15	-547.31	63.53
500	21.55	200.72	-	-	-	-785.13	-516.66	53.97
550	23.07	219.85	-	-	-	-782.86	-487.96	46.34
600	23.65	237.92	-	-	-	-782.03	-460.85	40.12
650	23.43	255.05	-	-	-	-782.41	-435.13	34.97
700	22.51	271.37	-	-	-	-783.85	-410.65	30.64
750	20.97	286.93	-	-	-	-786.28	-387.29	26.97
800	18.84	301.79	-	-	-	-789.62	-364.99	23.83



Figure 5. Calculated free energies of formation and associated reaction constants for the deltic, squaric and croconic acids as a function of temperature.

Table 3. Calculated enthalpies ($\Delta_r H$) and free energies ($\Delta_r G$) of reaction and associated reaction constants (Log K) for the reactions of thermal decomposition of the squaric and croconic acids [(A) and (B)]. The values of $\Delta_r H$ and $\Delta_r G$ are in units of kJ/mol.

	Squaric acid. Reaction (A)				Croconic acid. Reaction (B)			
I (K)	$\Delta_{\mathbf{r}}\mathbf{H}$	$\Delta_{\mathbf{r}}\mathbf{G}$	Log K	I(K)	$\Delta_{\mathbf{r}}\mathbf{H}$	$\Delta_{\mathbf{r}}\mathbf{G}$	Log K	
298.15	-243.64	-105.45	42.68	298.15	3.22	-188.17	32.97	
310	-252.03	-107.71	42.46	320	-1.85	-208.57	34.05	
330	-259.10	-109.62	42.29	340	-5.79	-226.52	34.80	
330	-265.69	-111.05	42.05	360	-9.24	-243.97	35.40	
340	-272.64	-112.84	41.89	380	-11.98	-260.68	35.83	
350	-279.28	-114.32	41.68	400	-14.36	-277.03	36.17	
360	-285.79	-115.68	41.47	420	-16.33	-292.94	36.43	
-	-	-	-	440	-17.93	-308.46	36.62	
-	-	-	-	460	-19.18	-323.63	36.75	
-	-	-	-	480	-20.11	-338.47	36.83	
-	-	-	-	500	-20.75	-353.02	36.88	



Figure 6. Calculated free energies and associated reaction constants for the reactions of thermal decomposition of squaric and croconic acids [(A) and (B)] as a function of temperature.

	Squaric	acid. React	ion (C)	Croconic acid. Reaction (D)			
$I(\mathbf{K})$	Δ _r H	$\Delta_{\mathbf{r}}\mathbf{G}$	Log K	$\Delta_{\mathbf{r}}\mathbf{H}$	$\Delta_{\mathbf{r}}\mathbf{G}$	Log K	
298.15	-1261.70	-1349.65	236.45	-1436.01	-1551.40	271.79	
310	-1264.03	-1356.04	228.49	-1441.53	-1566.22	255.65	
320	-1266.18	-1361.42	222.22	-1445.82	-1579.03	242.58	
330	-1267.44	-1366.32	216.27	-1449.58	-1591.33	230.89	
340	-1269.64	-1371.57	210.71	-1452.60	-1602.89	220.33	
350	-1271.23	-1376.51	211.47	-1455.20	-1614.06	210.77	
360	-1272.69	-1381.32	200.42	-1457.36	-1624.78	202.07	
-	-	-	-	-1459.11	-1635.11	194.11	
-	-	-	-	-1460.49	-1645.07	186.80	
-	-	-	-	-1461.52	-1654.70	180.06	
-	-	-	-	-1462.24	-1664.02	173.84	

Table 4. Calculated enthalpies $(\Delta_r H)$ and free energies $(\Delta_r G)$ of reaction and associated reaction constants (Log K) for the reactions of complete combustion of the squaric and croconic acids [(C) and (D)]. The values of $\Delta_r H$ and $\Delta_r G$ are in units of kJ/mol.



Figure 7. Calculated free energies and associated reaction constants for the reactions of complete combustion of the squaric and croconic acids [(C) and (D)] as a function of temperature.

Table 5. Calculated enthalpies $(\Delta_r H)$ and free energies $(\Delta_r G)$ of reaction and associated reaction constants (Log K) for the reaction of conversion between the squaric and the croconic acid in the presence of carbon monoxide [(E)]. The values of $\Delta_r H$ and $\Delta_r G$ are in units of kJ/mol.

T(K)	$\Delta_{\mathbf{r}}\mathbf{H}$	$\Delta_{\mathbf{r}}\mathbf{G}$	Log K
298.15	-108.67	-55.47	9.72
310	-108.18	-52.79	8.89
320	-107.77	-50.52	8.25
330	-107.45	-48.36	7.65
340	-107.06	-46.12	7.09
350	-106.73	-43.96	6.56
360	-106.44	-41.83	6.07



Figure 8. Calculated free energies and associated reaction constants for the reaction of conversion between the squaric and the croconic acid [(E)] in the presence of carbon monoxide as a function of temperature.

III.2. Raman spectroscopic characterization

III.2.1. Deltic acid

The experimental Raman spectrum of the deltic acid was recorded by Lautié *et al.*⁷ from a synthetic sample obtained using the synthetic method of Serratosa.⁹³ The Raman spectra was recorded employing a Raman spectrometer MOLE (Jobin-Yvon technology) using the 514.5 nm output of an Ar⁺ laser as excitation source. The spectral resolution was $\pm 1 \text{ cm}^{-1}$. The theoretical Raman spectrum was computed at T=298 K, λ =532 nm, and with a full width at half maximum (FWHM) of 5 cm⁻¹. The experimental and computed spectra are compared in Figure 9. In Figure 9, the Raman spectrum have been split in four different spectral ranges in order to improve their visualization: (A) 0-1400 cm⁻¹; (B) 1400-1600 cm⁻¹; (C) 1700-2000 cm⁻¹; and (D) 2000-3000 cm⁻¹. As can be appreciated the agreement is very satisfactory. Therefore, a normal mode analysis was carried out in order to assign the bands of the Raman spectrum. The experimental Raman band shifts and assignments performed by Lautie *et al.*⁷ are shown in Table 6, where the Raman shifts and assignments of the spectrum of free deltate ion, $C_3O_3^{2-}$, are also given.⁹⁻¹⁰ The band wavenumbers of the experimental and calculated spectra along with the corresponding calculated intensities and assignments are given in Table 7. Pictures of the atomic motions in the Raman active vibrational modes are depicted in Figure S.1 of the Supporting Information. The most important motion types of the atoms of the C_3 carbon skeleton in the Raman active vibrational normal modes are shown in Figure 10.



Figure 9. Experimental⁷ and calculated Raman spectra of the deltic acid in the spectral regions: (A) 0-1400 cm⁻¹; (B) 1400-1600 cm⁻¹; (C) 1700-2000 cm⁻¹; (D) 2000-3000 cm⁻¹.

Table 6. Observed band wavenumbers (cm^{-1}) and assignments of the Raman spectrum of the deltic acid.⁷⁻¹⁰ The letters s, m and w stand for strong, medium and weak bands.

Band	Solid delt	ic acid ⁷⁻⁸	Free C ₃ O ²⁻ ₃ ion ⁹⁻¹⁰		
name	Raman shift	Assignment	Raman shift	Assignment	
а	63	-	-	-	
b	109	-	-	-	
С	133	-	-	-	
d	269	$\nu(OH \cdots O)$	258/236 (m)	$\delta^{oop}(CO)$	
е	315	$\delta(C - OH)$	321/346 (m)	-	
f	379	$\gamma(C - OH)$	-	-	
g	641	$\gamma(C=0)$	-	-	
h	653	$\gamma(C=0)$	689/696 (w)	δ_{ring}	
i	799	Cycle	786/803 (s)	Ring breathing	
j	957	$\delta(C=0)$	-	-	
k	964	$\delta(C=0)$	996/992 (m)	-	
l	1350	$\nu(C - OH)$	-	-	
m	1458	$\delta(0-H)$	1432/1446 (w)	$\nu(CC)$	
n	1569	Cycle	-	-	
0	1940	$\nu(C=0)$	1818/1835 (m)	$\nu(CO)$	

Table 7. Experimental and calculated Raman band wavenumbers (cm^{-1}) of the Raman spectrum of the deltic acid, irreducible representations, calculated intensities and assignments. The meaning of the acronyms T, R, Pr, Di, Ro and Co is given in Figure 10.

Band	F 78	Calc.	Irr. Rep.	Int.	Assignment
name	Exp. ⁷⁻⁸	(This work)	(D_{2h})	(Å4)	(This work)
а	63	86.1	A_g	9.97	$T(C = C) + T(OH) + \delta^{oop}(C - C = O)$
b	109	127.4	B_{1g}	3.24	R(C=C)+T(OH)
		131.0	A_g	4.74	$T(C_3) + \delta^{oop}(C - C = 0)$
		149.1	B_{3g}	13.15	$Def(C_3) + T(OH)$
С	133	169.3	A_g	2.66	$T(C_3) + \nu(C = 0) + \delta(C - H - 0)$
d	269	268.1	B_{2g}	0.67	$T(C_3) + \delta^{oop}(C - C = 0) + \delta^{oop}(C - 0 - H)$
е	315	326.3	A_g	62.19	<i>T</i> (<i>OH</i>)
f	379	404.1	B_{1g}	64.57	$T(OH) + \delta(C - C = 0)$
g	641	622.0	B_{3g}	25.12	R(C = C)
h	653	627.5	A_g	9.48	$Pr(C_3) + \delta^{oop}(C - C = 0)$
i	799	796.5	A_g	416.88	$Di(C_3) + \nu(C = 0) + T(OH)$
j	957	941.9	B_{3g}	19.94	$Ro(C_3) + \delta^{oop}(C - O - H)$
k	964	952.2	B_{1g}	33.74	"
l	1350	1370.6	A_g	184.97	$\delta(C-H-O)$
m	1458	1472.6	B_{3g}	12.68	"
n	-	1510.4	B_{1g}	11.34	"
0	1569	1559.8	B_{2g}	13.21	"
		1576.2	A_g	97.01	"
p	1940	1930.2	A_g	79.14	$Co(C_3) + \nu(C - 0) + \delta(C - 0 - H)$
q	-	2369.4	B_{3g}	287.6	ν(<i>0H</i>)
r	-	2383.2	B_{1g}	1000.2	"
S	-	2537.5	A_g	3044.2	"



Figure 10. Motion types of the atoms of the C_3 carbon skeleton in the Raman active vibrational normal modes of the deltic acid: (A) Translation of the C = C bond, T(C = C) [86.1 cm⁻¹]; (B) Rotation of the C = C bond, R(C = C) [127.4, 622.0 cm⁻¹]; (C) Translation T(C₃) [131.0, 169.3, 268.1 cm⁻¹]; (D) Plane rotation, Pr(C₃) [627.5 cm⁻¹]; (E) Dilatation, Di(C₃) [796.5 cm⁻¹]; (F) In-plane rotation, Ro(C₃) [941.9, 952.2 cm⁻¹]; (G) Contraction, Co(C₃) [1930.2 cm⁻¹]

As may be observed from Tables 6 and 7, the assignment of the Raman spectrum performed by Lautié *et al.*⁷ is incomplete and differs significantly from the present assignment. The empirical assignment derived from the spectrum of the free deltate ion is also incomplete and very different from that of Lautié *et al.*⁷ and that of the present work. The origin of the bands (*a*) to (*c*) and (*q*) to (*s*) is completely unknown

The lowest wavenumber band (a) having an observed wavenumber of 63 cm⁻¹ corresponds to the calculated one at 86 cm⁻¹. This band is assigned to a combination of translations of the C = C double bonds (see Fig. 10.A), translations of the hydroxyl OH^- ions and out of plane C - C = O bending vibrations. As observed in Table 7 and Figure 9.A, the next band (b), placed in the experimental spectrum at 109 cm⁻¹ contains in fact three different peaks corresponding to the calculated bands at 127, 131 and 149 cm⁻¹. The first of these bands is assigned to rotations of the C = C double bonds (see Fig. 10.B) and translations of the hydroxyl ions. The second one is attributed to C_3 carbon skeleton translations (Fig. 10.C) and C - C = 0out of plane bending vibrations. Finally, the third one is ascribed to C₃ carbon skeleton deformations and hydroxyl translations. The experimental band placed at 133 cm^{-1} (c) corresponds to the calculated band at 169 cm⁻¹ and it is assigned to a combination of C_3 ring translations, C = 0 stretching and C -H - 0 bending vibrations. The three bands (a) to (c), were left without assignment in the experimental works.

The band at 269 cm⁻¹ (d) was calculated to be at 268 cm⁻¹ and was assigned by Lautié *et al.*⁷ to OH ··· O stretching vibrations. However, it should be attributed to C₃ ring translations and C – C = O and C – H – O out of plane bending vibrations. The pair of bands located at 315 and 379 cm⁻¹, (*e*) and (*f*), are very intense and correspond to the theoretical ones at 326 and 404 cm⁻¹. The first one must be reassigned to hydroxyl ion translations plus C – C = O bending vibrations. The bands at 641 and 653 cm⁻¹, (*g*) and (*h*), were calculated to be at 622 and 627 cm⁻¹. The first of these bands should be ascribed to C = C double bond rotations and the second one C₃ ring plane rotations (see Fig. 10.D) and C – O – H bending vibrations.

The very intense band at 799 cm⁻¹ (*i*) is very well reproduced by the theoretical calculations (796 cm⁻¹). It must be assigned to a combination of C₃ ring dilatations (Fig. 10.E), C = 0 bond stretching and hydroxyl translations. While the empirical assignment of this band attributes correctly this band to the C_3 ring dilatations (called "cycle" by Lautié *et al.*⁷ or ring breathing vibrations by Junqueira *et al.*¹⁰) the remaining motions are neglected. However, as it can be seen in Figure S.1 of the Supporting Information, this motion has the smaller weight in the description of the corresponding normal mode, the weights of the other two motions being much larger. We prefer the use of the term dilatation instead of breathing, which is very commonly used, because we need to distinguish among ring dilatation and contraction (see Figure 10). The same term is frequently used for these two different types of motion.

The bands at 957 and 964 cm⁻¹ (bands (j) and (k), respectively), are reproduced theoretically at 942 and 952 cm⁻¹. Both bands must be reassigned to in-plane rotations of the C₃ ring (Fig. 10.F) and C - O - H bending vibrations. The experimental bands located at 1350 and 1458 cm⁻¹, (l) and (m), are reproduced theoretically at 1371 and 1473 cm⁻¹. The band (n), whose existence is obvious in Figure 9.B, was apparently not seen by Lautié *et al.*,⁷ and is situated at 1510 cm⁻¹ in the theoretical spectrum. The band at 1569 cm⁻¹ (*o*) is resolved into two components in the theoretical Raman spectrum placed at 1560 and 1576 cm⁻¹. The last five bands must be assigned solely to C - H - O bending vibrations. The band at 1940 cm⁻¹ is located at 1930 cm⁻¹ in the theoretical spectrum. It is attributed to a combination of contractions of the C₃ rings (Fig. 10.G), C - O stretching and C - O - H bending vibrations.

Finally, it must be noted that the Raman spectrum bands in the long wavenumber region, 2000-3000 cm⁻¹, are very broad. In fact, for this reason, Lautié *et al.*,⁷ do not attempted to resolve the bands in this region. As it can be seen in Figure 9.D, the broad band located near 2500 cm⁻¹ in the experimental spectrum is the result of the broadening and overlap of the three bands (q), (r) and s) which appear in the theoretical spectrum at 2369, 2383 and 2537 cm⁻¹, respectively. These three bands are assigned to OH stretching vibrations. The reason for the large broadening of these bands is clearly the strong hydrogen bonding present in the crystal structure of deltic acid.⁶

III.2.2. Squaric acid

The experimental Raman spectrum of the squaric acid has been obtained in many occasions and may be found in several published articles.¹¹⁻¹⁹ For example, this spectrum was recorded at

309 K by Baglin and Rose¹¹ in 1970 with an estimated resolution of ±4 cm⁻¹. Similarly, Nakashima and Balkanski¹² recorded the Raman spectrum of squaric acid at 77 K in 1976 with an improved resolution of about $\pm 2 \text{ cm}^{-1}$. A more recent version was recorded at room temperature by Georgopoulos et al.¹³ in 2013 using a Bruker FT-Raman spectrometer having an estimated resolution of $\pm 2 \text{ cm}^{-1}$. The experimental Raman band shifts from these works and the band assignments performed by Georgopoulos et al.¹³ are shown in Table 8. In Table 8 the Raman shifts and assignments of the spectrum of the free squarate ion are also given.¹⁴⁻¹⁵ The theoretical Raman spectrum was computed at T=298 K, λ =532 nm, and with a full width at half maximum (FWHM) of 20 cm⁻¹. The experimental spectrum of Georgopoulos et al.¹³ and the computed spectra are compared in Figure 11. As in the case of deltic acid, the agreement is very satisfactory and, therefore, a normal mode analysis was carried out in order to assign the bands in the Raman spectrum. The band wavenumbers of the experimental and calculated spectra along with the corresponding calculated intensities and assignments are given in Table 9. Images of the atomic motions in the Raman active vibrational modes are shown in Figure S.2 of the Supporting Information. The motion types of the atoms of the C4 carbon skeleton in the Raman active vibrational normal modes are shown in Figure 12.

As shown in Table 9, the band with the lowest wavenumber (α), at 85 cm⁻¹, really corresponds to two contributing bands which are very close and are calculated theoretically to be at 89 and 90 cm⁻¹. These two bands must be assigned to translations of the squaric acid molecules along the a and b directions, respectively. The second band which was found at 153 cm⁻¹ in the works by Baglin and Rose¹¹ and Georgopoulos et al.¹³ was shown by Nakashima and Balkanski¹² to be resolved into three different contributions at 153, 158 and 164 cm⁻¹. This feature is in complete agreement with the findings from the theoretical spectrum, since three near bands at 152, 163 and 172 cm⁻¹ were obtained. The theoretical methods allow going far beyond of predicting this feature since permit to assign them. The first one must be attributed to out-of-plane C - O - H bending vibrations and hydroxyl translations, the second one to out of plane C - O - H and C - C = O bending vibrations and the last one to out-of-plane C - C = O bending vibrations. The weak band (b) located at 230 cm⁻¹ found by Baglin and Rose¹¹ was not encountered by Nakashima and Balkanski¹² and Georgopoulos et al.¹³ and it is absent in the theoretical spectrum. It is identified as a combination of the (α) and (α) bands as indicated in the assignment column of Table 9.

The band at 241 cm⁻¹ is calculated to be at 260 cm⁻¹ and it must be assigned to a combination of translations of the four membered carbon rings plus C = 0 and C - 0 bond rotations and hydroxyl translations. The band at 274 cm⁻¹ corresponds to the computed one at 272 cm⁻¹ which is ascribed to hydroxyl translations, OH stretching and C - 0 - H bending vibrations. This is the first band for which we could try to assign the band empirically from the weak band of the squarate ion at 294 cm⁻¹. However, this band was assigned to out of plane CO bending vbrations. This vibrational motion associated to this assignment is completely different to the combination of atomic motions found in the present work. The band found at 308 cm⁻¹ is reproduced theoretically at 306 cm⁻¹ and should be assigned to hydroxyl translations and C - C = 0 bending vibrations. Band (*f*), at 383 cm⁻¹, corresponds to the computed one at 415 cm⁻¹ is assigned to C-O-H and C-C=O bending vibrations coupled with OH stretching vibrations. The weak band (g) at 445 cm⁻¹, absent in the computed spectrum, is a combination band arising from bands (*a*) and (*e*).

The experimental band at 626 cm^{-1} is, as some previous bands, shown to contain two contributions which were theoretically found at 626 and 629 cm⁻¹. These two bands having similar intensities in the calculated spectrum are attributed to rotations of the C₄ carbon ring plane (see Fig. 12.B). The next three experimental bands, at 639, 735 and 858 cm⁻¹ are reproduced theoretically at 620, 721 and 852 cm⁻¹, respectively. These three bands must be assigned to a combination of C = 0 and C - 0bond stretching, C - O - H bending vibrations and three different atomic motions of the C_4 carbon skeleton: $Rh(C_4)$, $Co(C_4)$ and $Ro(C_4)$ (see Figs. 12.C, D, E). The first one is the square to rhombus deformation, the second one is the ring contraction and the last one is the C₄ ring rotation in the plane. With the exception of the assignment of the second band (assigned in the experimental works as a ring breathing mode), the other bands must be reassigned and even for the second one, the most part of the vibrational contributions are neglected in the empirical assignment. The bands at 1053, 1062 and 1174 cm⁻¹ correspond to the calculated bands at 1037, 1047 and 1141 cm⁻¹. Both bands are attributed to square to trapezium deformations, C - O - H bending vibrations and, for the first two bands, C - OO bond stretching vibrations. Again, the empirical assignment should be corrected.

The band (o) at 1251 cm⁻¹ is clearly an overtone band of the (h) band (2 v_1 where v_1 =626 cm⁻¹). This fact has not been noticed in any of the previous spectroscopic works. The bands at 1307 and 1362 and 1514 cm⁻¹ are reproduced theoretically at 1324, 1331 and 1551 cm⁻¹. These bands are assigned to C -0 - H bending vibrations. The band at 1385 cm⁻¹, absent in the computed spectrum, is recognized as a combination of (i) and (j) bands. The band at 1532 cm⁻¹ is reproduced at 1570 cm⁻¹ and has the same assignment that the (j) band, which as mentioned in the previous paragraph is mainly ascribed to square to rhombus deformations, $Rh(C_4)$. The bands (u) and (v), at 1630 and 1677 cm⁻¹, are combination bands of the (n) and (g) bands and (l) and (h) bands, respectively. Finally, the band 1829 cm^{-1} is reproduced at 1784 cm⁻¹ and assigned to a combination of ring contractions and C = 0, C - 0 and 0 - H bond stretching vibrations.

While Baglin and Rose²¹ identified a low wavenumber combination band in the infrared spectrum of squaric acid at 346 cm⁻¹ and another in the Raman spectrum at 1677 cm⁻¹ (band (v) see above), the theoretical calculations have allowed to recognize one overtone band (o) and four additional combination bands, (b), (g), (r) and (u), in the Raman spectrum. The situation is similar to that found in the Raman spectrum of natroxalate²⁴ and oxammite²⁵ minerals in which a large number of bands of these types were identified. To conclude this section is very important to note that in the theoretical Raman spectrum there are two additional bands at 2047 and 2054 cm⁻¹ which are assigned to OH stretching vibrations. These bands must obviously exist in the experimental spectrum but, as in the case of the deltic acid, they must be relatively weak or too broad to be resolved experimentally.

Dand		Solid squ	aric acid ¹¹⁻¹³		Free C	$L_4 O_4^2 - ion^{14-15}$
name	Raman shift ¹¹ (309 K)	Raman shift ¹² (77 K)	Raman shift ¹³ (298 K)	Assignment ¹³	Raman shift	Assignment
a	85	89	-	-	99(w)	_
u	05	90	-	-)) (w)	
а	153	153.5	-	-	-	-
		158.5	-	-	-	-
		164	-	-	-	-
b	230	-	-	-	-	-
С	241	-	-	-	259 (s)	$\delta_{oop}(CO)$
d	274	245	-	-	294 (w)	δ(CO)
е	308	307	-	-	-	-
f	383	383	380 (s)	δ(CO)	350 (m)	v(CO)
g	445	-	-	-	-	-
h	626	626	-	-	(() (,,,,,)	δ (CO)
		629	-	-	002 (VW)	$0_{00p}(CO)$
i	639	635	635 (s)	Ring torsion	647 (s)	v(CC)
j	735	729	726 (s)	Ring breathing	723 (s)	Ring breathing
k	858	856	-	-	844	$\nu(CC)+\nu(CO)$
l	1053	1051	1049 (w)	v(CC)	1000(a)	u(CC)
т	1062	-	-	-	1090 (8)	V(CC)
n	1174	1177	1172 (m)	v(CC)	1123 (vs)	$\nu(CC)$
0	1251	1250	1297 (w)	v(CO)	-	-
p	1307	-	-	-	-	-
q	1362	-	-	-	-	-
r	1385	-	-	-	-	-
S	1514	1510	1510 (w)	ν(CO)	1520 (vg)	u (CO)
t	1532	-	1579 (w)	v(CO)	1330 (vs)	V(CO)
и	1630	-	1616 (w)	v(CO)	1593 (s)	v(CO)
v	1677	-	-	1380+308=168821	-	-
w	1829	-	1824(vw)	v(CO)	1794 (w)	$\nu(CC)+\nu(CO)$

Table 8. Observed band wavenumbers (cm^{-1}) and assignments of the Raman spectrum of the squaric acid.¹¹⁻¹⁵ The letters vs, s, m, w and vw stand for very strong, strong, medium, weak and very weak bands.



Figure 11. Experimental¹³ and calculated Raman spectra of the squaric acid. The overtone band (*o*) and the combination bands (*b*), (*g*), (*r*), (*u*) and (*v*) are highlighted in red.

Table 9. Observed and calculated Raman band wavenumbers (cm^{-1}) in the Raman spectrum of the squaric acid, irreducible representations, calculated intensities and assignments. The meaning of the acronyms T, Pr, Rh, Co, Ro and Tp is given in Figure 12.

Band	Exn. ¹¹	Exn. ¹²	Exn. ¹³	Calc.	Irr. Rep.	Int.	Assignment
name	пур.	пур.	Ехр:	(This work)	(\mathcal{C}_{2h})	(Å ⁴)	(This work)
α	85	89	-	76.4	Ag	0.40	$T^b(H_2C_4O_4)$
		90	-	99.2	Ag	0.41	$T^a(H_2C_4O_4)$
а	153	153.5	-	151.9	Bg	10.26	$\delta^{oop}(C - O' - H') + T(OH^{-})$
		158.5	-	162.8	Bg	40.02	$\delta^{oop}(C - O - H) + \delta^{oop}(C - C = O)$
		164	-	171.5	Bg	28.00	$\delta^{oop}(C-C=0)$
b	230	-	-	-	-	-	$v_1 + v_2 = 238 (v_1 = 153, v_2 = 85)$
С	241	-	-	260.1	Ag	0.23	$T(C_4) + R(C = 0) + R(C - 0) + T(OH^{-})$
d	274	245	-	272.2	Ag	20.22	$T(OH^{-}) + \nu(O' - H') + \delta(C - O' - H')$
е	308	307	-	306.9	Ag	15.97	$T(OH^{-}) + \delta(C - C = 0)$
f	383	383	-	415.2	Ag	162.35	$\delta(C - C = 0) + \nu(0 - H) + \delta(C - 0 - H)$
g	445	-	-	-	-	-	$\nu_1 + \nu_2 = 461 (\nu_1 = 308, \nu_2 = 153)$
h	626	626	-	613.2	Bg	29.70	$Pr(C_4)$
		629	-	614.6	Bg	25.60	"
i	639	635	635 (s)	620.1	Ag	176.62	$Rh(C_4) + \nu(C = 0) + \nu(C - 0) + \delta(C - 0 - H)$
j	735	729	726 (s)	721.1	Ag	679.47	$Co(C_4) + \nu(C = 0) + \nu(C - 0) + \delta(C - 0 - H)$
k	858	856	-	851.6	Ag	6.53	$Ro(C_4) + \nu(C = 0) + \nu(C - 0) + \delta(C - 0 - H)$
l	1053	1051	1049 (w)	1036.8	Ag	102.86	$Tp(C_4) + \nu(C - 0) + \delta(C - 0 - H)$
m	1062	-	-	1047.5	Ag	7.54	"
n	1174	1177	1172 (m)	1140.8	Ag	593.75	$Tp(C_4) + \delta(C - O - H)$
0	1251	1250	1297 (w)	-	-	-	$2 v_1 = 1252 (v_1 = 626)$
p	1307	-	-	1324.4	Ag	507.87	$\delta(C - O - H)$
q	1362	-	-	1331.2	Ag	1077.13	"
r	1385	-	-		-	-	$v_1 + v_2 = 1374 (v_1 = 735, v_2 = 639)$
S	1514	1510	1510 (w)	1551.0	Ag	151.97	$\delta(C - O - H)$
t	1532	-	1579 (w)	1569.9	Ag	522.81	$Rh(C_4) + \nu(C = 0) + \nu(C - 0) + \delta(C - 0 - H)$
и	1630	-	1616 (w)	-	-	-	$v_1 + v_2 = 1619 (v_1 = 1174, v_2 = 445)$
v	1677	-	-	-	-	-	$v_1 + v_2 = 1679 (v_1 = 1053, v_2 = 626)$
w	1829	-	1824(vw)	1783.8	Ag	110.04	$Co(C_4) + \nu(C = 0) + \nu(C - 0) + \nu(0 - H)$



Figure 12. Motion types of the atoms of the C₄ carbon skeleton in the Raman active vibrational normal modes of the squaric acid: (A) Translation, $T(C_4)$ [260.1 cm⁻¹]; (B) Plane rotation, $Pr(C_4)$ [613.2, 614.6 cm⁻¹]; (C) Square to rhombus deformation, $Rh(C_4)$ [620.1, 1569.9 cm⁻¹]; (D) Contraction, $Co(C_4)$ [721.1, 1783.8 cm⁻¹]; (E) In-plane rotation, $Ro(C_4)$ [851.6 9 cm⁻¹]; (F) Square to trapezium deformation, $Tp(C_4)$ [1036.8, 1047.5, 1140.8 cm⁻¹].

III.2.3. Croconic acid

The experimental Raman spectrum of the croconic acid used as reference in this work, was recorded very recently by Milan-Garces *et al.*²⁰ using a Bruker Senterra Raman spectrometer coupled to an Olympus microscope. The excitation source used

was the 632.8 nm output of a He–Ne laser. The spectral resolution was 3–5 cm⁻¹. The measured experimental Raman band shifts and the band assignments given by these authors are listed in Table 10 where the Raman shifts and assignments of the spectrum of the free croconate ion²¹⁻²³ are also given. The theoretical Raman spectrum was computed at T=298 K, λ =532 nm,

and with a FWHM of 10 cm⁻¹. The experimental spectrum and the computed spectra are compared in Figure 13. Again, the agreement between these spectra is very good. The band wavenumbers of the experimental and calculated spectra along with the corresponding calculated intensities and the assignments performed in this work are given in Table 11. The atomic motions in the Raman active vibrational modes of croconic acid are displayed in Figure S.3 of the Supporting Information. The motion types of the atoms of the C5 carbon skeleton in the Raman active vibrational normal modes are shown in Figure 14.

The band with the lowest wavenumber (a) in the Raman spectrum of the croconic acid is calculated to be at 121 cm⁻¹. It must be assigned to translations of the C_5 ring, C = 0 bond and hydroxyl ions. Similarly, the bands (b) to (f) are found in the theoretical Raman spectrum at 151, 243, 383, 407 and 506 cm⁻¹ and are assigned as detailed in Table 11. These low wavenumber bands, (a) to (f), do not appear in the experimental work of Milan-Garces et al,²⁰ their analysis beginning at the band (g). The band (g) is observed at 521 cm⁻¹ and calculated at 522 cm⁻¹. It must be assigned to a combination of deformations of the five membered rings, C - C bond stretching, and C - C = O bending vibrations and hydroxyl translations. The bands placed at 534 and 553 cm⁻¹ are reproduced theoretically at 543 and 547 $\rm cm^{-1}$ and are ascribed to $\rm C_5$ ring elongations (see Fig. 14.C), C = 0 bond and hydroxyl translations. The band at 565 cm⁻¹ is calculated to be at 561 cm⁻¹ and attributed to a combination of C_5 ring widenings (see Fig. 14.D), C = 0bond and hydroxyl translations. The peak at 634 cm⁻¹ is found at 631 cm⁻¹ and assigned to ring contractions (see Fig. 14.K), C = 0 and C - 0 bond stretching and C - 0 - H bending vibrations. While this band was correctly assigned to ring breathing vibrations in the experimental works, this description is very incomplete since it neglects the remaining vibrational contributions to the corresponding mode. The bands at 731 and 806 cm⁻¹ are found at 734 and 783 cm⁻¹. Both are assigned to ring bending vibrations (see Fig. 14.E).

The three bands found experimentally at 1074, 1082 and 1150 cm^{-1} were found theoretically at 1063, 1081 and 1152 cm^{-1} . The three bands are assigned to C - O - H bending vibrations. The band at 1169 cm⁻¹ is shown to be composed of two bands which were determined at 1165 and 1171 cm⁻¹. These bands are assigned to a combination of $C_2 - C_3$ and $C_4 - C_5$ bond elongation, $C_3 - C_4$ bond shortening and C - O - H bending vibrations (See Fig. 14.F). The band at 1236 cm⁻¹ also contains two contributing bands which were determined at 1245 and 1257 cm⁻¹ which are attributed to a combination of $C_5 - C_1$ bond shortening and C - O - H bending vibrations (See Fig. 14.G). The bands (s), (t) and (v) at 1302, 1384 and 1522 cm⁻¹ are reproduced at 1310, (1378, 1397) and 1517 cm⁻¹ and are assigned to C - O - H bending vibrations. A band corresponding to the theoretical band (u) at 1477 cm^{-1} was not reported by Milan-Garces et al.²⁰ although it is clearly observed in the experimental spectrum given in Figure 13. The band (w) at 1652 cm⁻¹ contains three contributing bands whose theoretical wavenumbers are 1552, 1564 and 1576 cm⁻¹. The first is assigned to $C_3 - C_4$ bond shortening, $C_4 - C_5$ bond rotation, C =O bond stretching and C - O - H bending vibrations and the other two ones to $C_3 - C_4$ bond elongation, the approach of the C_2 and C_5 non bonded carbon atoms, C = 0 bond stretching and C - O - H bending vibrations. The band at 1722 cm⁻¹ is found at 1632 cm⁻¹ and assigned to $C_3 - C_4$ bond rotation, C = 0bond stretching and C - O - H bending vibrations. The last band in the experimental spectrum, placed at 1757 cm⁻¹ is reproduced theoretically at 1694 cm⁻¹ and assigned to ring contractions (see Fig. 14.K), C = 0 and C - 0 bond stretching and C - O - H bending vibrations.

As for the deltic and squaric acids, several Raman bands attributable to OH stretching vibrations appear in the theoretical Raman spectrum (at 2282, 2313, 2334 and 2439 cm⁻¹) and were not reported in the experimental works.

Table 10. Observed band wavenumbers (cm ⁻¹) and assignments of the Raman spectrum of the croconic acid. ²⁰⁻²³ The letters s, m and w and
vw stand for strong, medium, weak and very weak bands.

Band	Solid cr	oconic acid ²⁰	Free $C_5 O_5^{2-}$ ion ²¹⁻²³		
name	Raman shift	Assignment	Raman shift	Assignment	
d	-	-	360 (vw)	δ(CO)	
g	521	-	-	-	
ĥ	534	$\delta_{oop}(CO)$	537 (m)	$\delta_{oop}(CO)$	
i	553	δ_{ring}	556 (s)	δ_{ring}	
j	565	-	-	-	
k	634	Ring breathing	637 (m)	Ring breathing	
l	731	C_2 - C_3 - C_4 out of plane ring twisting	-	-	
m	806	δ(CO)	-	-	
n	1074	v(CC)	-	-	
0	1082	$\nu(CC)$	-	-	
p	1150	$\nu(CC)$	-	-	
q	1169	$\nu(CC)$	-	-	
r	1236	v(CC)	1246 (m)	v(CC)	
S	1302	-	-	-	
t	1384	-	-	-	
ν	1522	v(CC) + v(CO)	1591 (s)	ν(CO)	
w	1652	v(CC) + v(CO)	-	-	
x	1722	v(CO)	1721 (w)	ν(CO)	
y	1757	v(CO)	-	-	

Table 11. Observed²⁰ and calculated Raman band wavenumbers (cm^{-1}) in the Raman spectrum of the croconic acid, irreducible representations, calculated intensities and assignments. The meaning of the acronyms T, Pr, El, Wi, Co, Be, Sh, R and Ap is given in Figure 14.

Band name	Exp. ²⁰ Solid 298 K	Calc. (This work)	Irr. Rep. (C2v)	Int. (Å ⁴)	Assignment (This work)	
а	-	120.9	A ₁	77.16	$T(C_5) + T(C = 0) + T(OH^{-})$	
b	-	151.5	A_1	231.13	$\delta^{\text{oop}}(C - O - H) + \delta^{\text{oop}}(C - C = 0)$	
С	-	242.7	A_1	35.69	$T(C_5) + \delta^{oop}(C - C = 0) + T(OH^{-})$	
d	-	382.6	A_2	20.84	$\delta(C - C = 0) + T(OH^{-})$	
е	-	407.3	A ₁	178.21	"	
f	-	506.4	A_1	366.70	$Pr(C_5) + \delta^{oop}(C - C = 0) + \delta^{oop}(C - 0 - H)$	
g	521	522.0	A_1	329.00	$Def(C_5) + \nu(C_2 - C_3) + \delta^{oop}(C - C = 0) + T(OH^-)$	
h	534	543.1	A_1	536.01	$EI(C_5) + T(C = 0) + T(OH)$	
i	553	547.3	B ₁	302.02	"	
j	565	561.3	B ₂	290.76	$Wi(C_5) + T(C = 0) + T(OH)$	
k	634	630.6	A_1	2200.18	$Co(C_5) + \nu(C = 0) + \nu(C - 0) + \delta(C - 0 - H)$	
l	731	733.6	A ₁	36.51	$Be(C_5)$	
m	806	783.0	B_1	5.0	$Be(C_5)$	
п	1074	1062.7	A_1	286.78	$\delta(C - O - H)$	
0	1082	1080.9	A ₁	86.17	"	
p	1150	1151.7	A_1	271.30	"	
q	1169	1165.2	B_1	152.37	$El(C_2 - C_3) + Sh(C_3 - C_4) + El(C_4 - C_5) + \delta(C - O - H)$	
		1171.0	B ₂	128.12	"	
r	1236	1244.7	A_1	670.63	$Sh(C_5 - C_1) + \delta(C - O - H)$	
		1257.2	A ₂	216.75	"	
S	1302	1310.0	A_1	4291.03	$\delta(C - O - H)$	
t	1384	1377.7	B_1	155.62	"	
		1397.0	A ₂	242.34	"	
и	-	1477.2	A_1	448.02	"	
v	1522	1517.0	A_1	1383.14	"	
W	1652	1552.2	A ₁	1148.94	$Sh(C_3 - C_4) + Ro(C_4 - C_5) + v(C = 0) + \delta(C - 0 - H)$	
		1563.8	B_2	1041.10	$El(C_3 - C_4) + Ap(C_2 - C_5) + v(C = 0) + \delta(C - 0 - H)$	
		1576.5	B_1	1183.85	"	
x	1722	1632.1	A ₁	3583.24	$R(C_3 - C_4) + \nu(C = 0) + \delta(C - 0 - H)$	
у	1757	1693.9	A_1	3061.12	$Co(C_5) + v(C = 0) + v(C - 0) + \delta(C - 0 - H)$	



Figure 13. Experimental²⁰ and calculated Raman spectra of the croconic acid.



Figure 14. Motion types of the atoms of the C₅ carbon skeleton in the Raman active vibrational normal modes of the croconic acid: Vibrational motions involving the C₅ carbon skeleton: (A) Translation, T(C₅) [120.9, 242.7 cm⁻¹]; (B) Plane rotation, Pr(C₅) [506.4 cm⁻¹]; (C) Elongation, El(C₅) [547.3 cm⁻¹]; (D) Widening, Wi(C₅) [561.3 cm⁻¹]; (E) Bending Be(C₅) [733.6, 783.0, cm⁻¹]; (F) Elongation of C₂ – C₃ and C₄ – C₅ pentagon sides and C₃ – C₄ side shortening, El(C₂ – C₃) + Sh(C₃ – C₄) + El(C₄ – C₅) [1165.2, 1171.0 cm⁻¹]; (G) C₅ – C₁ side shortening, Sh(C₅ – C₁) [1244.7, 1257.2 cm⁻¹]; (H) C₃ – C₄ pentagon side shortening and C₄ – C₅ side rotation, Sh(C₃ – C₄) + Ro(C₄ – C₅) [1552.2 cm⁻¹]; (I) C₃ – C₄ side elongation and C₂ … C₅ approach, El(C₃ – C₄) + Ap(C₂ … C₅) [1563.8, 1576.5 cm⁻¹]; (J) C₃ – C₄ side rotation, R(C₃ – C₄) [1632.1 cm⁻¹]; (K) Contraction, Co(C₅) [630.6, 1693.9 cm⁻¹]

III.3. UV-Visible Optical Characterization

III.3.1. Band gaps

The croconic, squaric and deltic acids are correctly predicted to be insulator materials in the present calculations, the computed band gaps being 1.80, 2.90 and 5.04 eV, respectively. However, as it is well-known, the band gaps are generally underestimated by the Density Functional Theory (DFT) methods in the generalized gradient approximation.94-99 The accurate estimation of the band gaps would need to go beyond DFT-PBE using for example hybrid functionals. As it has been shown computationally,⁹⁴ for insulators, most of the differences between the DFT energy eigenvalues and the true excitation energies can be accounted for by a uniform shift of the unoccupied band energy levels. This way to deal with the band gap problem is so-called scissor correction. The shifting values for the croconic and squaric acids were 1.11 and 0.59 eV, respectively. The final values of the band gaps were 2.91 and 3.49 eV. These values were determined by adjusting the wavelength of the most intense peak in the computed UV-Vis absorption spectrum (see the next Sections) to the observed experimental wavelength for this peak. However, for deltic acid, the corresponding absorption spectrum has not been measured experimentally and, therefore, an average value of shifting constants obtained for the croconic and squaric acids, 0.85 eV, was used. When this shifting was applied the final band gap for value for deltic acid was 5.89 eV. It must be emphasized that the application of this correction

does not modify the shape of the computed optical functions reported in the next sections.

III.3.2. Croconic acid

The computed UV-Vis absorption spectrum and the reflectivity, optical conductivity, dielectric, refractive index and loss optical functions of the croconic acid are shown in Figure 15. The calculated UV-Vis absorption spectrum, covering the wavelength range from 200 to 450 nm, is displayed in Figure 15.A where it is compared with the experimental spectrum of Hartley et al.43 The remaining optical functions are plotted in the wavelength range 200-750 nm. As can be seen in Figure 15.A, the agreement between the calculated absorption spectrum and the experimental one is quite good. The position of the four main peaks in the calculated absorption spectrum are given in Table 12. The origin of these peaks was obtained by analyzing the inter-band electronic transitions in the croconic acid. As it can be seen in Figure S.4 of the Supplementary Information, the peak a located at λ_a =299 nm results from the excitation of one electron from the occupied band with number 95 ($\sigma(C = C)$ + $p_v(0)$) to the unoccupied band number 107 $(p_z(C) + p_x(0))$. Similarly, the peaks b, c and d at 314, 333 and 351 nm are due to excitations from the occupied bands 97, 99 and 97 to the unoccupied bands 109, 109 and 107, respectively, and represent mainly charge transfers from oxygen p_x orbitals to carbon p_z orbitals.



Figure 15. Computed and experimental⁴³ UV-Vis absorption spectra and computed optical functions of the croconic acid.

Table 12. Calculated peaks in the calculated UV-Vis absorption spectra of the croconic, squaric and deltic acids.

Peak name	Wavelength (nm)	Transition	Excitation energy (eV)	Absorption (cm ⁻¹)					
Croconic acid (highest occupied and lowest unoccupied bands: ϕ_{104} and ϕ_{105} ; band gap: 2.9									
а	299±1	ϕ_{95} - ϕ_{107}	4.15±0.01	135917.8					
b	314±1	ϕ_{97} - ϕ_{109}	3.95 ± 0.01	137852.9					
С	333±1	ϕ_{99} - ϕ_{109}	3.72±0.01	155047.6					
d	351±1	ϕ_{97} - ϕ_{105}	3.53 ± 0.01	156317.5					
Squaric aci	id (highest occupied a	nd lowest uno	ccupied bands: ϕ_{42} and ϕ_{43}	; band gap: 3.49 eV)					
а	269±1	ϕ_{40} - ϕ_{45}	4.61 ± 0.02	231153.8					
b	278±1	ϕ_{39} - ϕ_{46}	4.46 ± 0.02	192574.4					
Deltic acid	Deltic acid (highest occupied and lowest unoccupied bands: ϕ_{64} and ϕ_{65} ; band gap: 5.89 eV)								
а	179±1	$\phi_{64} - \phi_{66}$	6.93±0.04	289221.3					

III.3.3. Squaric acid

The computed UV absorption spectrum and the reflectivity, optical conductivity, dielectric, refractive index and loss optical functions of the squaric acid are shown in Figure 16. The calculated UV absorption spectrum, covering the wavelength range from 220 to 350 nm, is displayed in Figure 16.A compared with the experimental spectrum from the Bio-Rad Laboratories Spectral Database.⁴⁴ As for croconic acid, the remaining optical functions are plotted in the wavelength range 200-750 nm. As can be seen in Figure 16.A, the shape of the calculated and experimental absorption spectra is quite similar. The position of two peaks of the calculated absorption spectrum are given in Table 12. The peak *a* at λ_a =269 nm (see Figure S.4 of the Supplementary Information), results from the excitation of one electron from the occupied band with number 40 to the unoccupied band number 45 and represents a charge transfer from π orbitals to π^* orbitals. Similarly, the peak *b* at λ_b =378 nm is due to an excitation from the occupied band 39 to the unoccupied band number 46 and represents a similar type of charge transfer ($\pi \rightarrow \pi^*$).



Figure 16. Computed and experimental⁴⁴ UV absorption spectra and optical functions of the squaric acid.



Figure 17. Computed UV-Vis absorption spectra and optical functions of the deltic acid.

III.3.4. Deltic acid

The computed UV-Vis absorption spectrum and optical functions of the deltic acid are shown in Figure 17. The calculated UV absorption spectrum, covering the wavelength range from 220 to 350 nm, is displayed in Figure 17.A. The position of the main peak in the calculated absorption spectrum is found at 179 nm (see Table 12). This peak (see Figure S.4 of the Supplementary Information), results from the excitation of one electron from the occupied band with number 64 to the unoccupied band number 66. This transition represents mainly a ($\pi \rightarrow p_z(C)$) charge transfer.

IV. CONCLUSIONS

A complete characterization of the thermodynamic, Raman spectroscopic and ultraviolet-visible optical properties of the

deltic, squaric and croconic cyclic oxocarbon acids has been performed by employing a computational approach based on theoretical solid-state methods. The computational study was carried out using very demanding calculation parameters, better that usually used in this kind of studies, in order to provide a realistic characterization of these materials. While these important materials have been studied at length experimentally, their characterization was relatively poor. One of the most surprising features of this study is that the characterization of their Raman spectra, being very accurately recorded and published in a large number of excellent spectroscopic studies, ⁷⁻²³ was not reliable. This is due to the use of the empirical method of assignment as the only source of information, surely for lack of relevant theoretical solid-state calculations. The assumption of the transferability of the information from gas or liquid phase spectroscopic data to crystalline solids is very crude and leads

to significant problems. The assignment of the bands in their Raman spectrum was also very incomplete. Most bands were reassigned. This study points to the need of using more elaborate methods of spectral assignment.

The fundamental thermodynamic properties, including the specific heat, the entropy, the enthalpy and the Gibbs free energy as well as their temperature dependence were also obtained. The calculated specific heat of the squaric acid, one of the only properties known from experimental measurements, was in good agreement with the corresponding experimental value. The thermodynamic properties obtained for the three materials where then used to determine the thermodynamic properties of formation of these materials in terms of the elements as a function of temperature. In turn, the calculated thermodynamic properties of formation were used to evaluate the Gibbs free energies of reaction and associated reaction constants for the reactions of thermal decomposition, complete combustion and interconversion of the squaric and croconic acids. The experimental enthalpies of combustion of the croconic and squaric acids at room temperature were also well reproduced.

Finally, the ultraviolet-visible absorption spectrum and the reflectivity, optical conductivity, dielectric, refractive index and loss optical functions of the three materials were determined. The UV-Vis absorption spectra calculated for the croconic and squaric acids were in good agreement with their experimental counterparts and that of deltic acid, which has not been obtained experimentally, was predicted. The origin of the peaks in the absorption spectra of these materials was found by examining the inter-band electronic transitions in these materials.

SUPPORTING INFORMATION

Supporting Information associated with this article contain: (a) Computed thermodynamic functions of the deltic, squaric and croconic acids; (b) Pictures of the atomic motions associated to active Raman normal modes of these materials; (c) Calculated UV-Vis absorption spectra and optical functions.

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TOC Graphic

