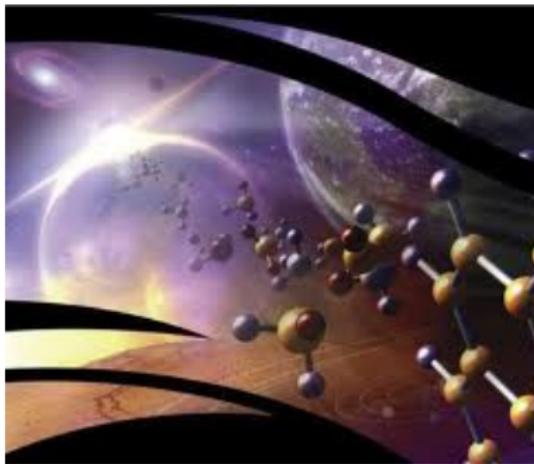


Reactions involving the formation and transformation of H_3^+

Octavio Roncero

Inst. Física Fundamental, CSIC

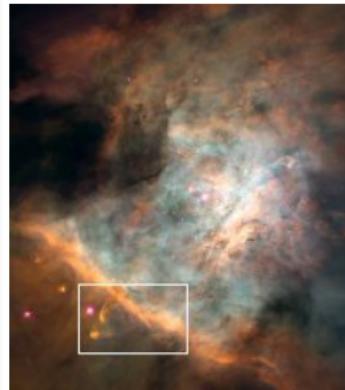
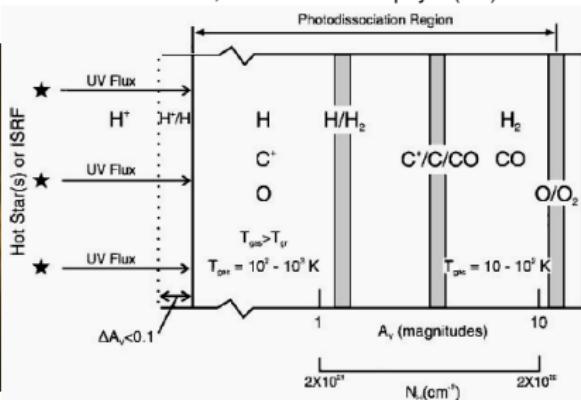
octavio.roncero@csic.es



Outline

- 1 Introduction
 - 2 $\text{H}_2 + \text{A}^+$ & $\text{HA}^+ + \text{H}$
 - 3 H_n^+ PES's
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 - 5 $\text{H}_2 + \text{H}_2^+$ direct reaction
 - 6 Conclusions

Molecular Universe



Stellar atmosphere:

atomic lines

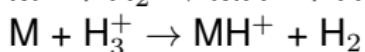
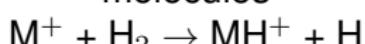
Atom+Atom collisions

Magnetic fields

Polarization of lines

Hanle effect

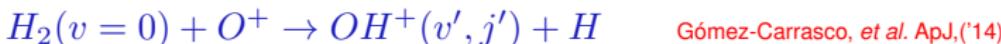
PDR
Formation of H₂
act as shield for other
molecules



Molecular clouds
low temperatures
varying densities
formation of complex
molecules

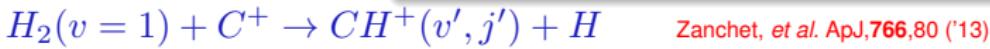
Chemical pumping in molecular clouds and PDR's

- Exothermic reactions of H₂ with atoms and ions



- Chemistry of $H_2(v > 0)$ in ISM Agúndez, et al. ApJ, 713,662 ('10)

Initial state dependent chemistry



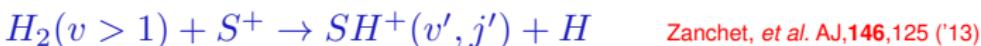
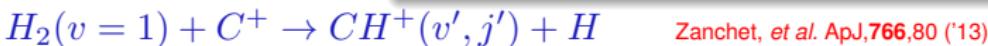
Chemical pumping in molecular clouds and PDR's

- Exothermic reactions of H₂ with atoms and ions



- Chemistry of H₂(v > 0) in ISM Agúndez, et al. ApJ, 713, 662 ('10)

Initial state dependent chemistry

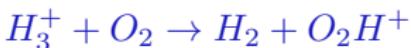


- The most abundant ion is H_3^+

- H_2^+ disappears rapidly in the exothermic $H_2 + H_2^+ \rightarrow H_3^+ + H$ reaction

- However $H_2 + H_3^+ \rightarrow H_3^+ + H_2$ (ortho/para conversion, deuteration)

- H_3^+ is very reactive with other species:



Chemical pumping in molecular clouds and PDR's

- Exothermic reactions of H₂ with atoms and ions



Gómez-Carrasco, et al. ApJ ('14)

- Chemistry of H₂(v > 0) in ISM Agúndez, et al. ApJ, 713, 662 ('10)

Initial state dependent chemistry

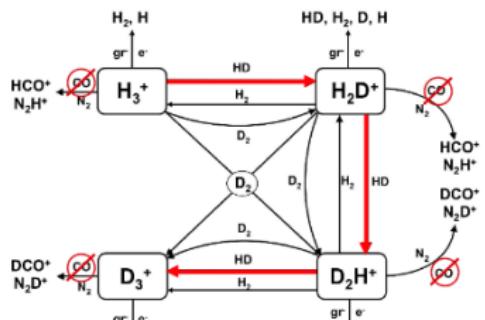


Zanchet et al. ApJ 766:80 ('13)



Zanchet et al. AJ 146:125 (13)

- The most abundant ion is H_3^+



- at low T molecules freeze
 - H_3^+ “lives” longer time
 - H_3 Deuteration occurs
 - Key process for high D/H

Outline

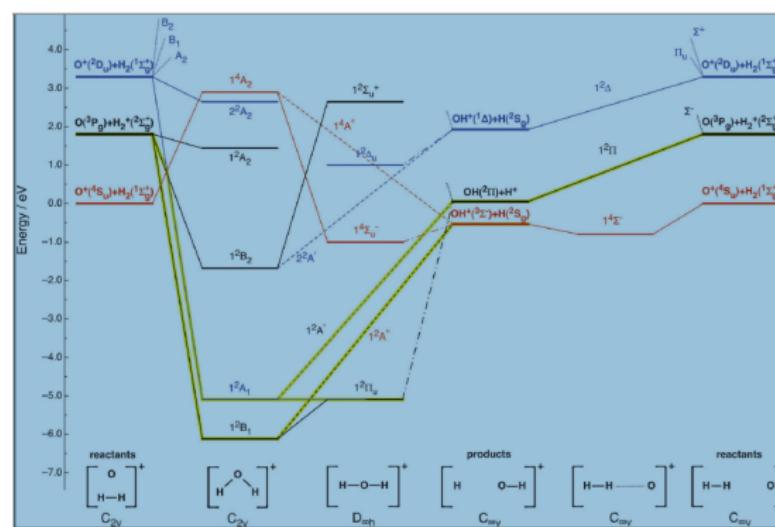
- 1 Introduction
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$\text{OH}^+(X^3\Sigma^-) + \text{H}$ inelastic collisions

- Need of accurate description:
 - Beyond mass scaling
 - Validity of rigid rotor
 - Exchange reaction:



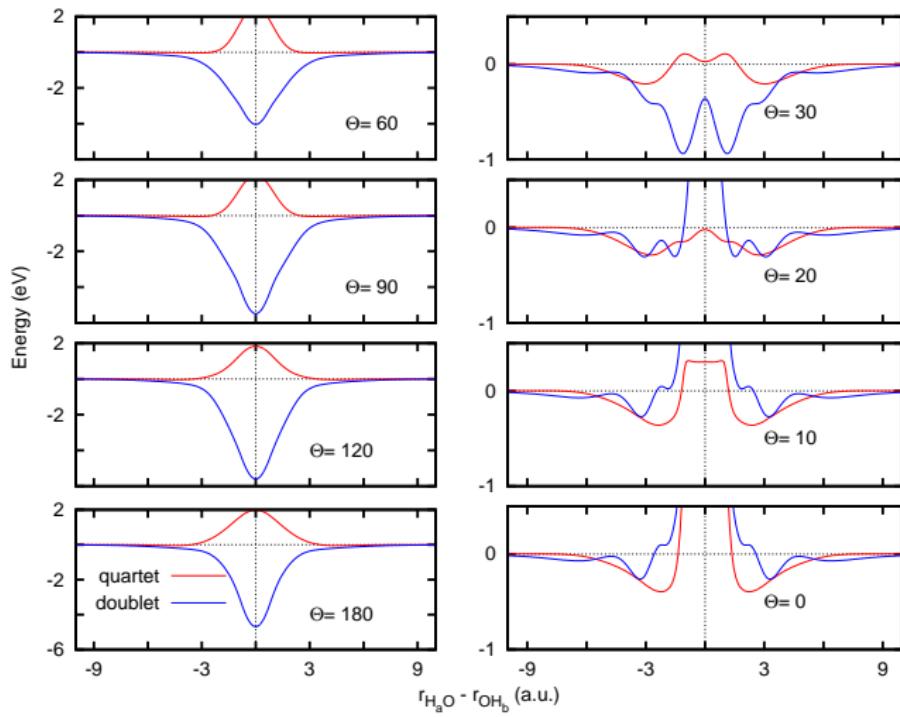
- 2 Open shell systems:
 $\text{OH}^+({}^3\Sigma^+) + \text{H}({}^2S)$
doublet and quadruplet states



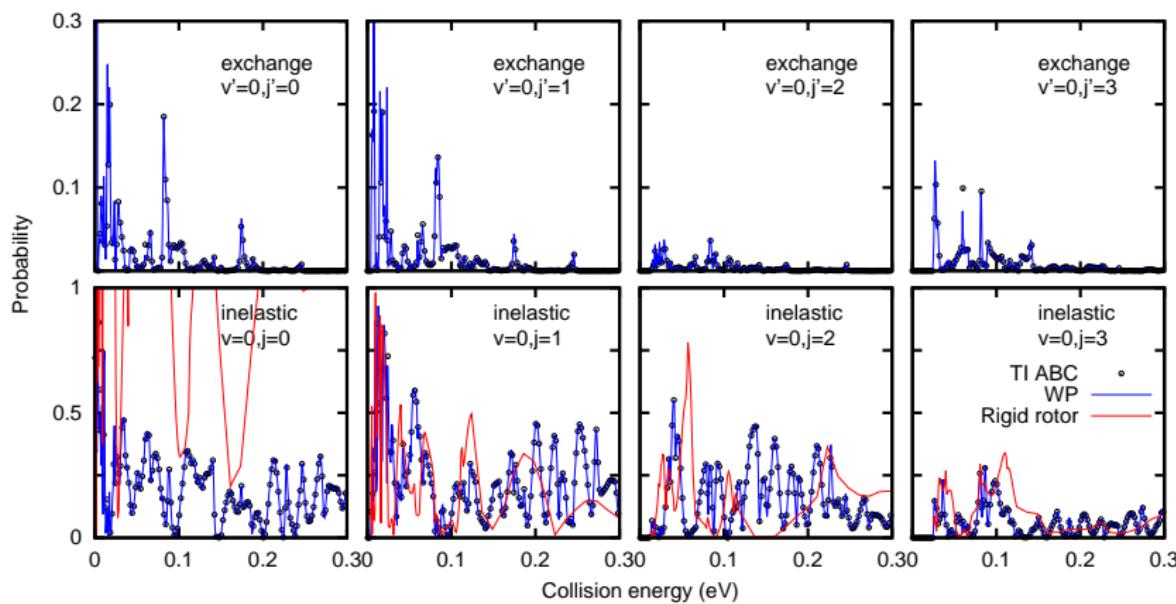
Bulut, Lique & Roncero
J.Chem. Phys A. (2015)

Paniagua *et al.*, PCCP, (2014), **16**, 23594

MEP's for exchange: two mechanisms

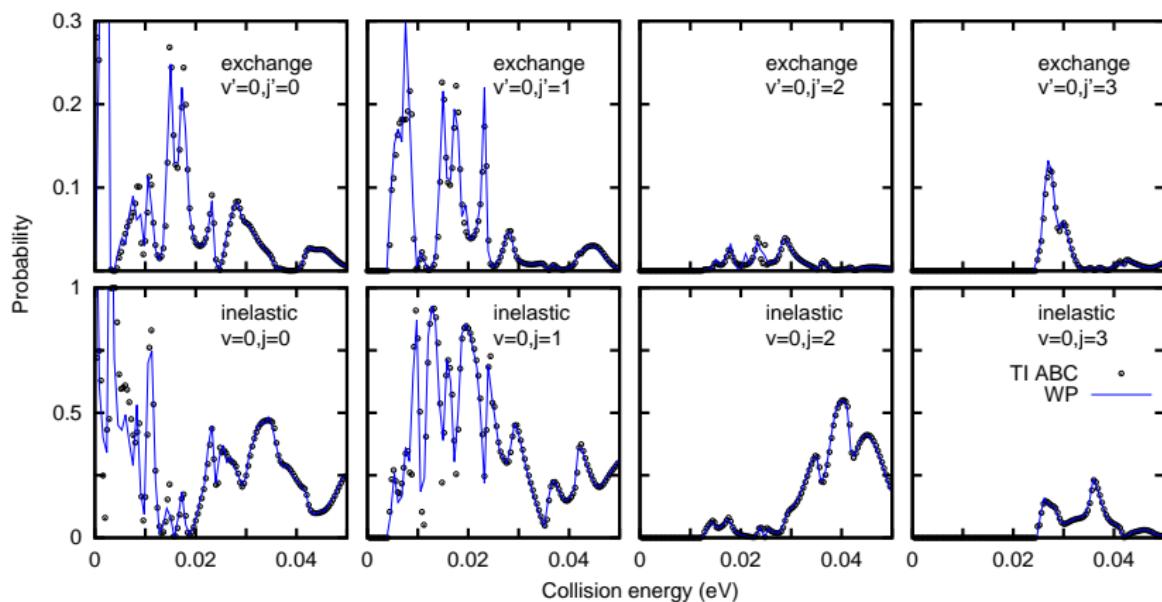


Quadruplet: H + OH⁺(v = 0, j = 0, J = 0)



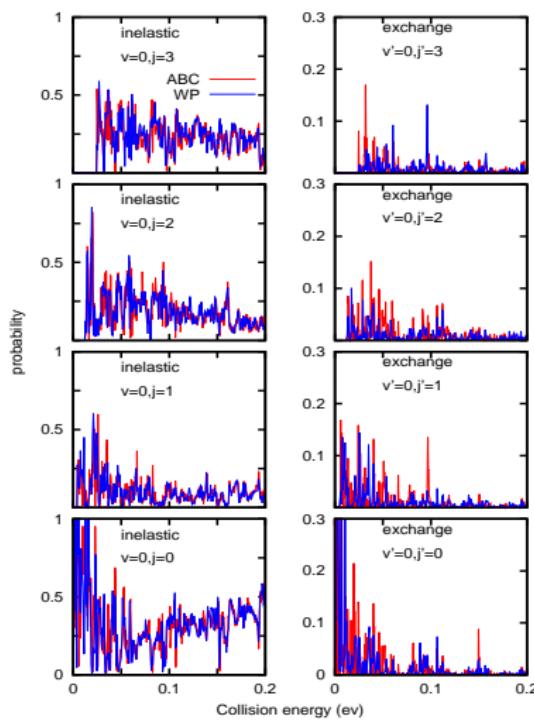
- Good agreement WP vs. TI-ABC
- Rigid rotor of the same order

Quadruplet: H + OH⁺(v = 0, j = 0, J = 0)



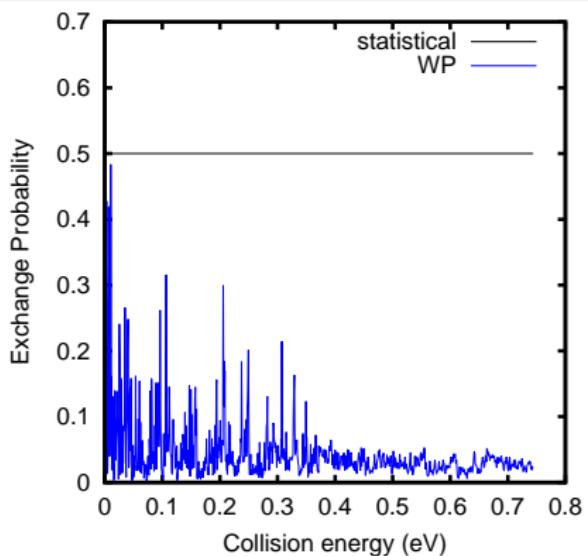
- Good agreement WP vs. TI-ABC, even at rather low energies!!
 - Rigid rotor of the same order

Doublet: H + OH⁺(v = 0, j = 0, J = 0)



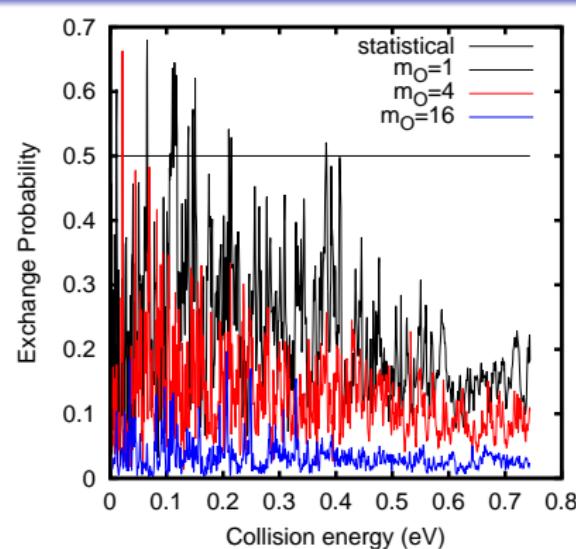
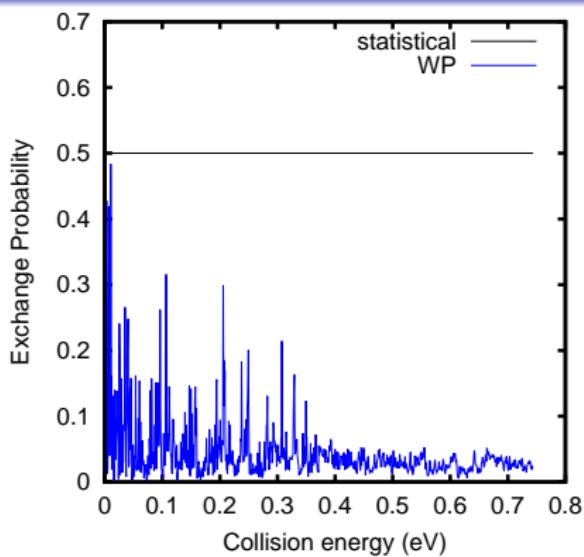
- Many resonances
- Comparison difficult using different coordinates
- Even ABC has problems when $D_{H-OH^+} \leq D_{OH^+}$

Is the exchange statistical in the doublet state?



**Non statistical
energy transfer inefficient**

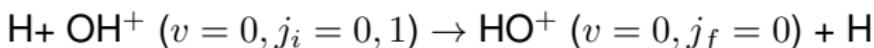
Is the exchange statistical in the doublet state?



**Non statistical
energy transfer inefficient**

- as energy decreases
- and mass difference reduces

Cross sections: inelastic and exchange for quadruplet



For quadruplet:

$8 \cdot 10^4$ iterations

$J=0, 5, 10, 15, 20, 25, 30, 40, \dots, 110$

$\Omega_{max} = 15$

$j_i = 0, 1$

For doublet:

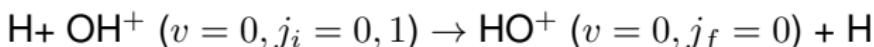
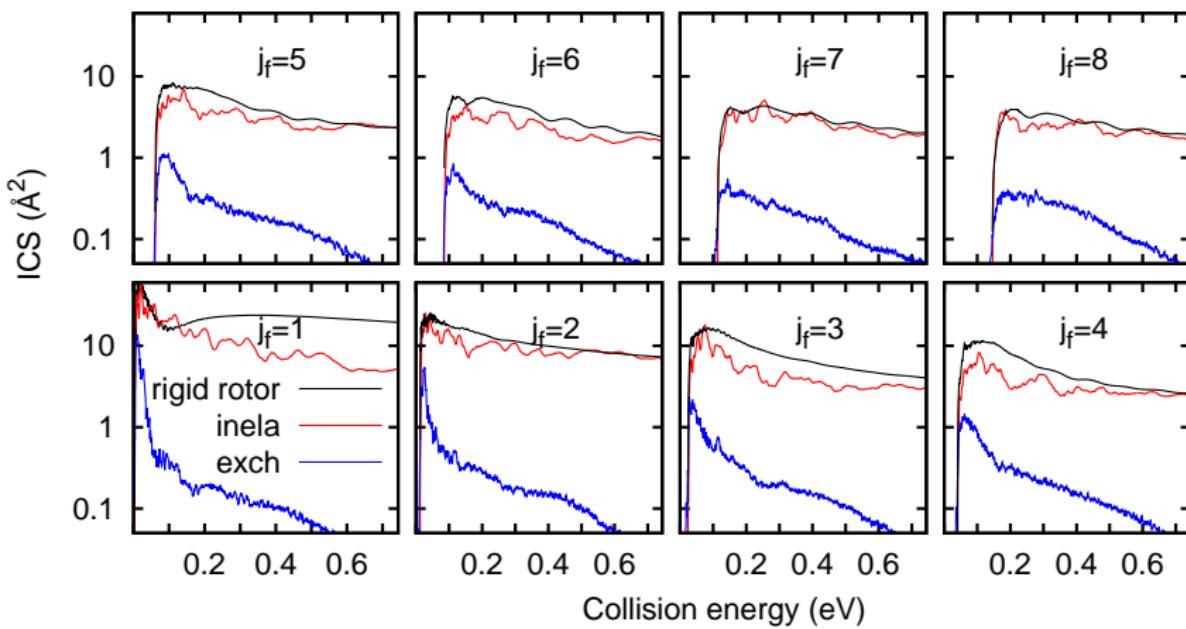
$2.5 \cdot 10^5$ iterations and denser grids

Calculations still in progress

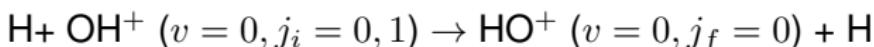
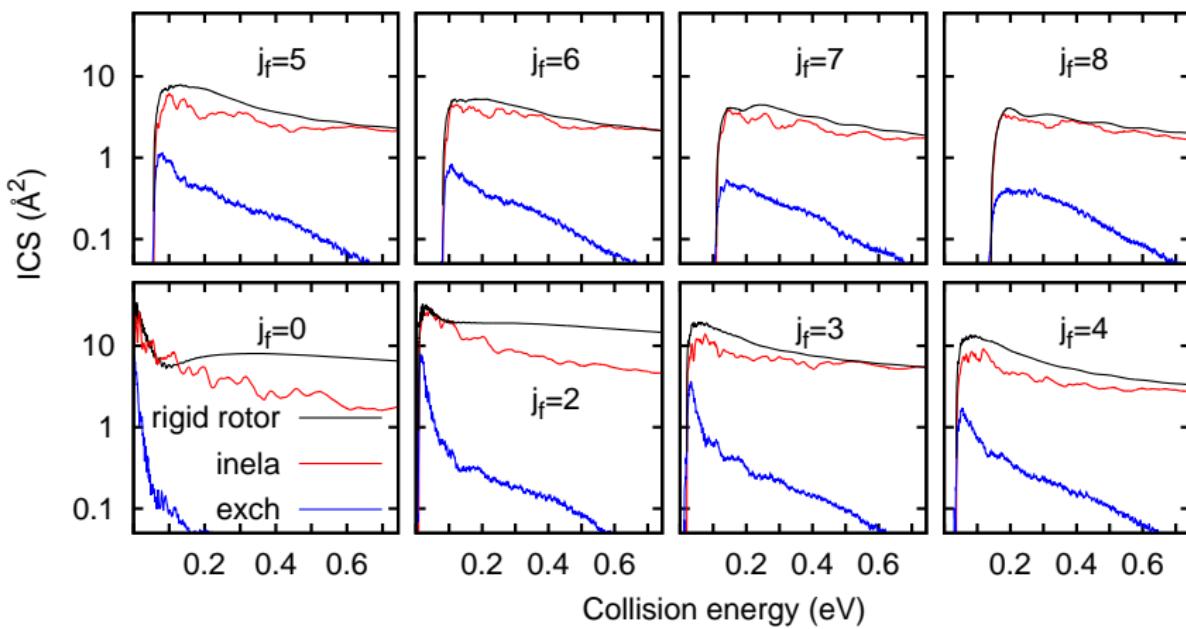
$\Omega_{max} = 19$

$j_i = 0$

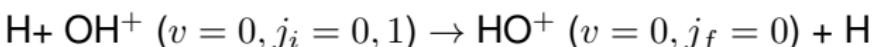
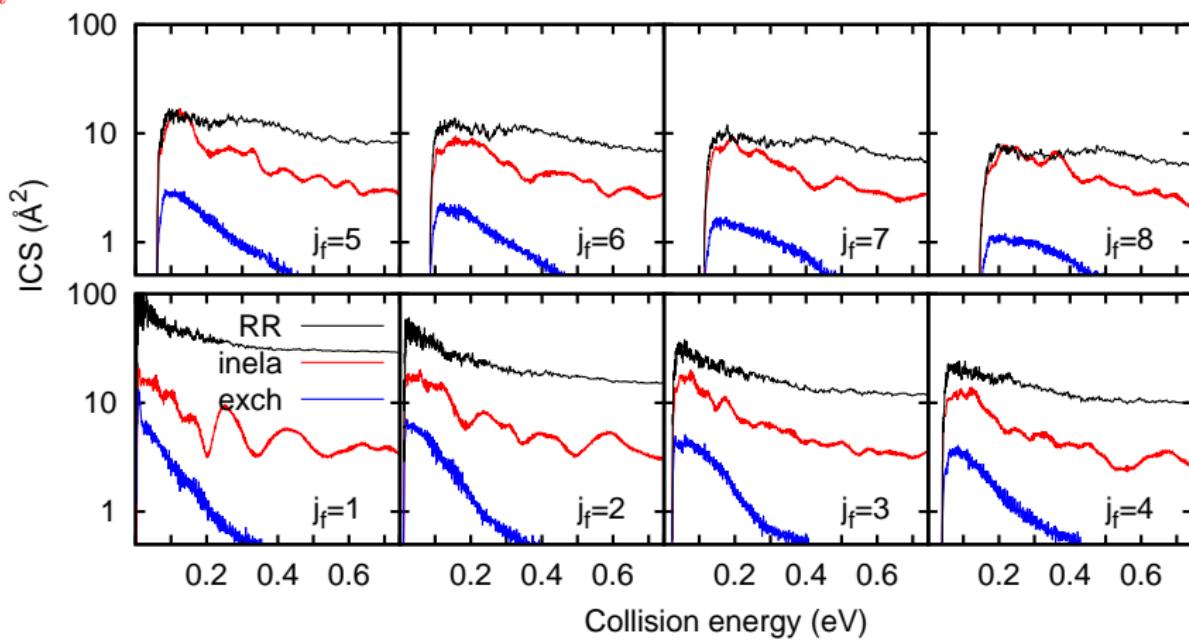
Cross sections: inelastic and exchange for quadruplet

 $j_i = 0$ 

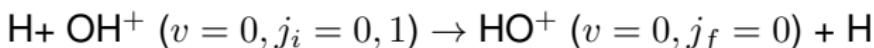
Cross sections: inelastic and exchange for quadruplet

 $j_i = 1$ 

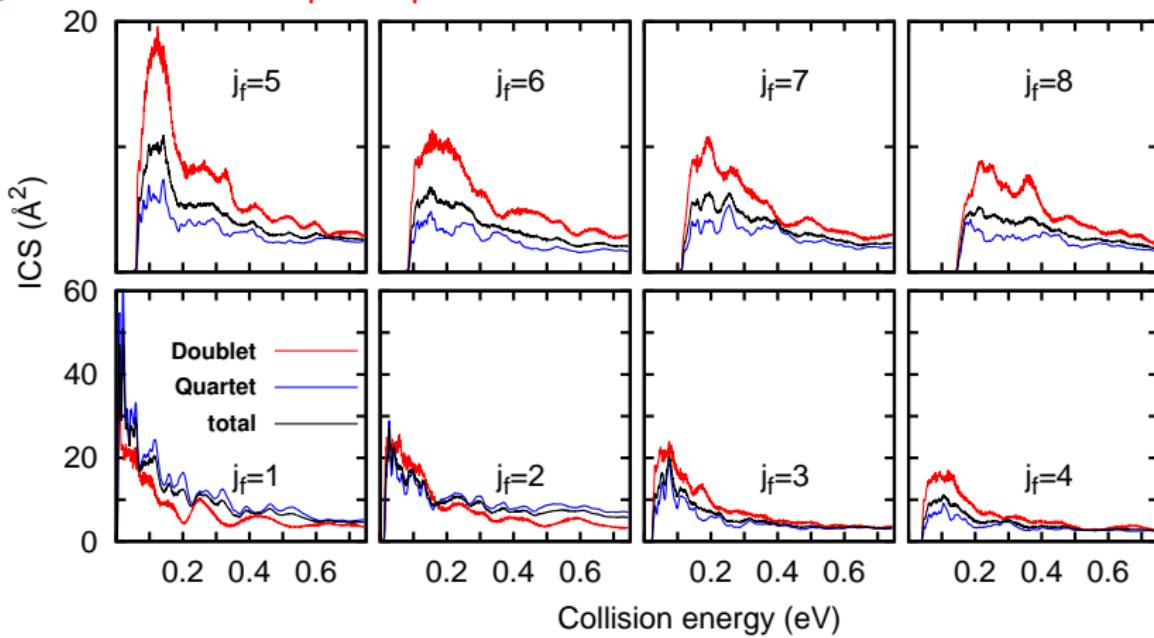
Cross sections: inelastic and exchange for doublet

 $j_i = 0:$ 

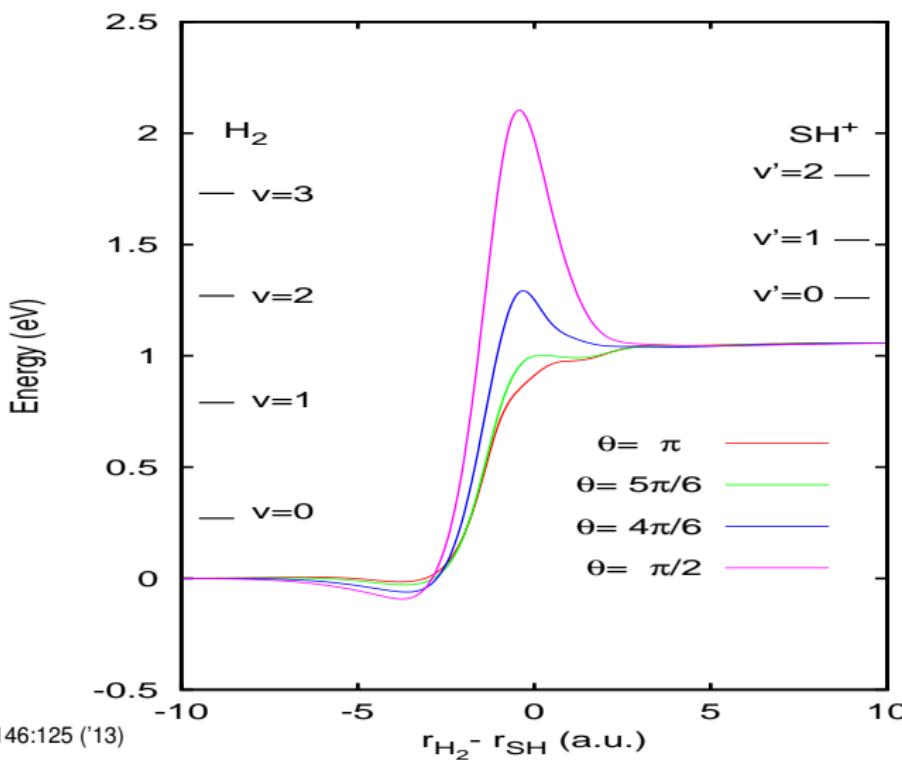
Cross sections: inelastic and exchange for doublet



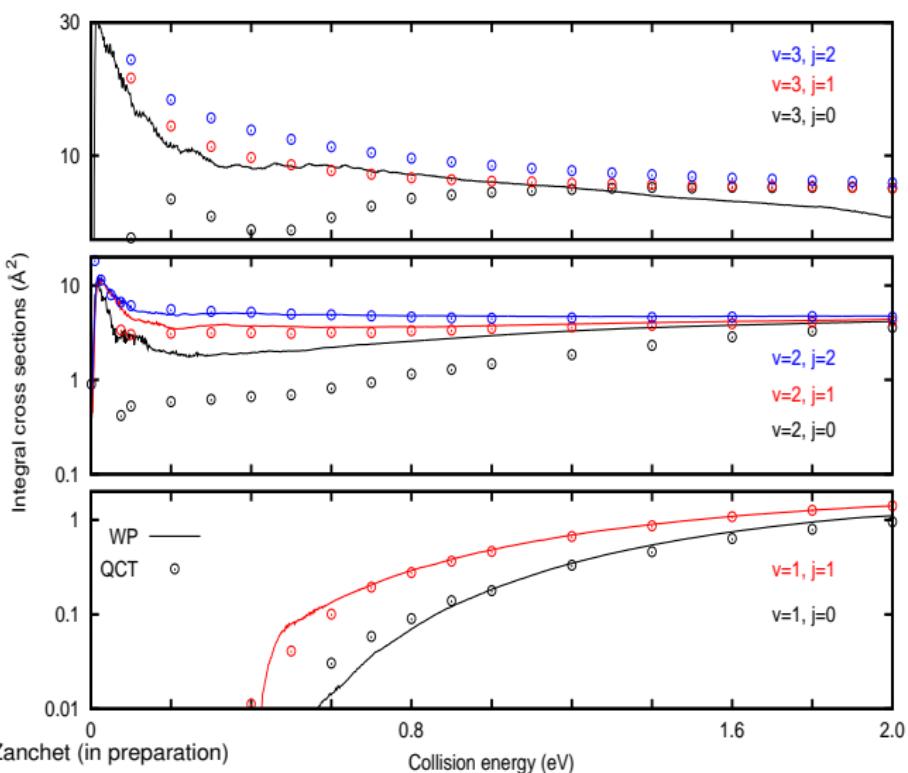
$j_i = 0$: Doublet vs. quadruplet

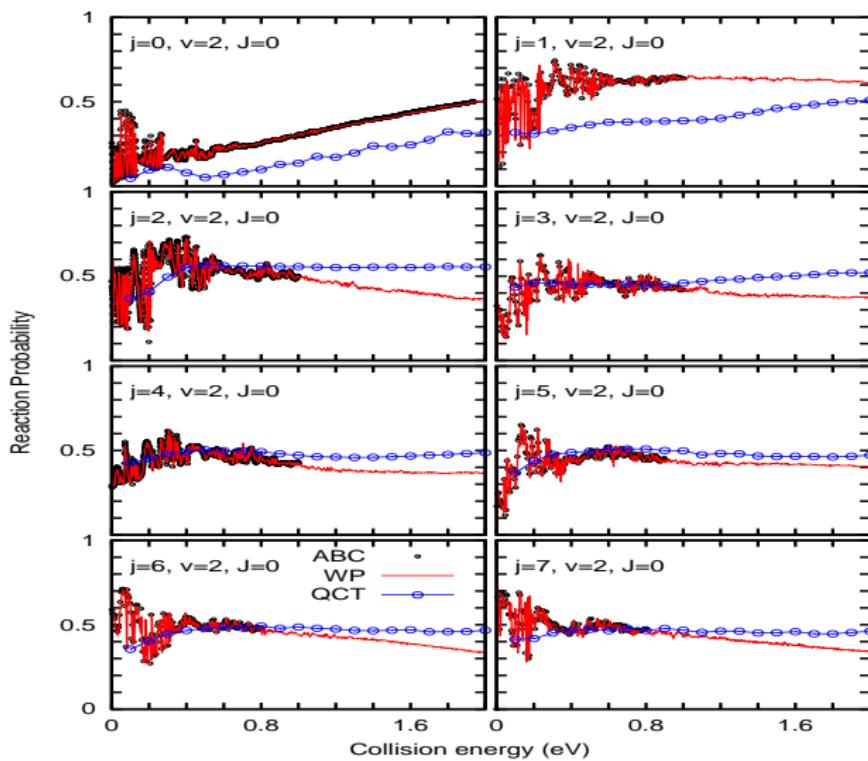


$S^+ + H_2(v, j)$ reaction: MEP

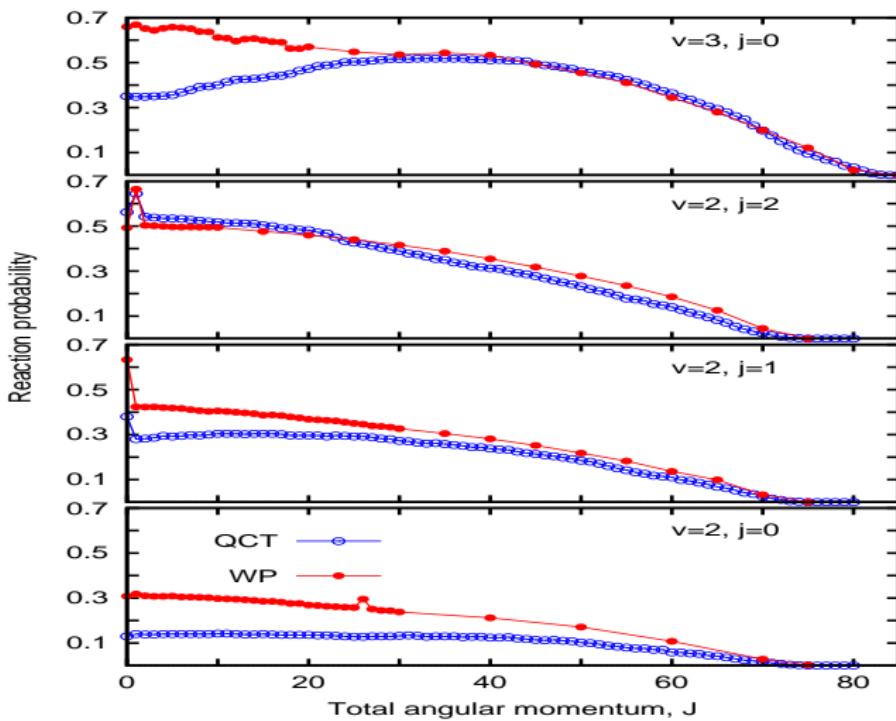


ICS: Wave Packet vs. QCT

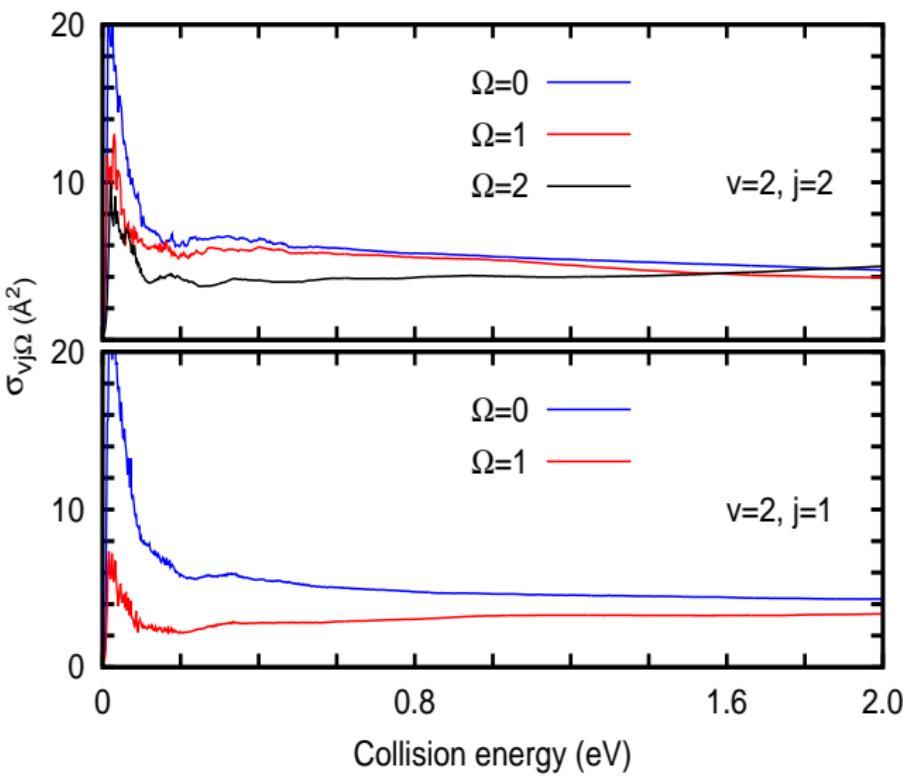


Reaction probabilities: $J = 0, v = 2, j$ 

Opacity functions: $J, v = 2, 3 j = 0, 1, 2$



Helicity : $\sigma_{vj\Omega}$

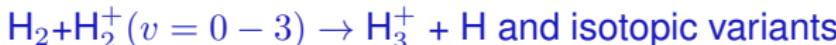


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Previous studies on H₄⁺.

- Experimental



Koyano & Tanaka ('80) ; Anderson et al. ('81) , Krenos et al. ('76) . . .

H_3^+ infrared spectra, produced in the previous reaction

Carrington & Kennedy ('84); ...

- Simulation of the potential

Diatomics-in-Molecules potential Muckerman *et al.* ('81)

PES with frozen H₃⁺ Alvarez-Collado *et al.* ('94)

Interpolated PES in the $H + H_3^+$ channel Moyano *et al.* ('04)

Stationary points and intersections along the MEP Varandas ('08)

No accurate global PES available

- $\text{H}_2^+ + \text{H}_2 \rightarrow \text{H}_3^+ + \text{H}$ collision simulations

Quasi-Classical trajectories, with Surface hopping using DIM

Stine & Muckerman (1976,...) ; Tully & coworkers (1976,...) ; Schatz & coworkers (1985,...) ; ...

$\text{H}_2 + \text{H}_3^+$ exchange collision: ortho/para conversion

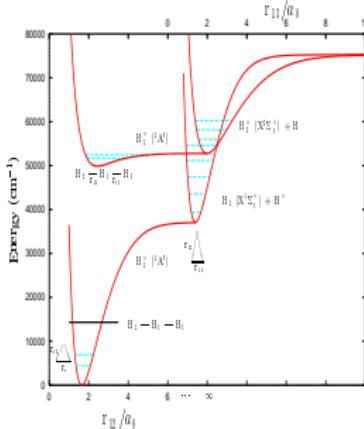
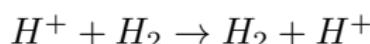
- QCT calculations: Moyano & Collins, ('03)
 - 7D Quantum WP with reduced dimensions: Bowman & co. ('10)
 - Statistical approaches including spin statistics:
Park & Light (2007) Hugo *et al.* (2009) McCall *et al.* (201

Objectives

- Full dimensional PES for H_4^+ and H_5^+
 - with good long range behavior
 - with analytical derivatives and non-adiabatic couplings
 - $H_2 + H_3^+ \rightarrow H_3^+ + H_2$ Is the reaction statistical?:
 - Going beyond pure statistical approach
 - Introduction of dynamical bias
 - Use global PES
 - Comparison with experimental $K^{hop}/K^{exchange}$
 - $H_2 + H_2^+ \rightarrow H_3^+ + H$

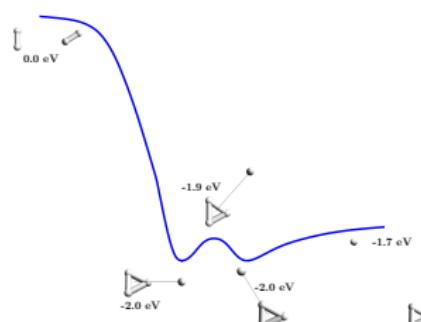
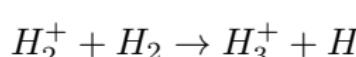
Non adiabatic couplings and Conical intersections Non adiabatic dynamics: MD with Quantum Jumps

H_n^+ : Potential energy surfaces



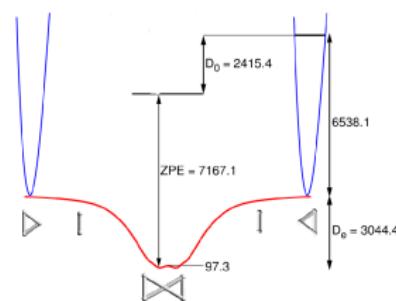
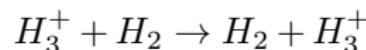
3 Dimensions

Aguado, et al. ('00)



6 Dimensions

Sanz-Sanz, et al. ('13)



9 Dimensions

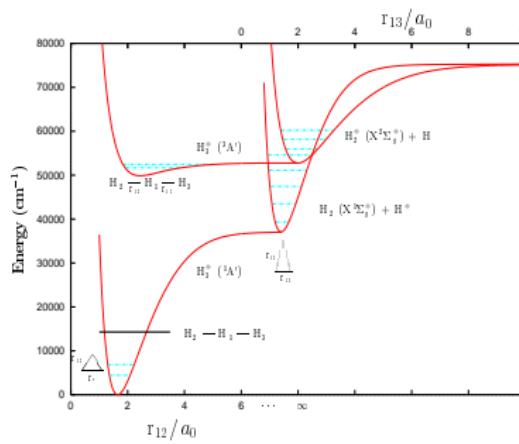
Aguado, et al. ('10)

- Proton exchange reaction: multistate problem
 - Long interaction to describe low temperature

Fitting procedure for H_n^+

Need of Multiple Surface to describe charge transfer

- $H_3^+ : V = H_{DIM} + V_{3B}$ Aguado *et al* (2000)
 - Long range behavior: Charge-quadrupole & induced dipole in diagonal of DIM matrix Velilla *et al* (2008)
 - Coupled diabatic PES's: 3-body terms in DIM matrix Varandas & co. (2007)



Fitting procedure for H_n^+

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 - Long range behavior: Charge-quadrupole & induced dipole in diagonal of DIM matrix Velilla *et al* (2008)
 - Coupled diabatic PES's: 3-body terms in DIM matrix Varandas & co. (2007)
 - $H_4^+ \& H_5^+ : V = H_{TRIM} + V_{NB}$

$$\hat{H}_{TRIM}(i) = \sum_{n>i, o>n} \hat{H}_{ino}^+(n-i, o-i) - 2 \sum_{p>i} \hat{H}_{ip}^+(p-i) \quad \text{Aguado et al.'10},$$

Sanz et al. ('13)

- High accuracy in fitting
 - Long range interaction
 - Symmetric under permutation
 - Accurate triatomic fragments
 - Non-adiabatic couplings

Analytical derivatives: Hellmann-Feynman theorem

Derivating electronic Schrödinger Eq. $\frac{\partial}{\partial \alpha} (H\phi_n) = \frac{\partial}{\partial \alpha} (W_n\phi_n)$

and premultiplying by $\phi_{n'}$ and integrating

$$\left\langle \phi_{n'} \left| \frac{\partial H}{\partial \alpha} \right| \phi_n \right\rangle = \delta_{nn'} \frac{\partial W_n}{\partial \alpha} + (1 - \delta_{nn'}) (W_n - W_{n'}) \left\langle \phi_{n'} \left| \frac{\partial \phi_n}{\partial \alpha} \right. \right\rangle$$

$$\phi_n = \sum_d T_n^d f_d \quad \rightarrow \quad \left\langle \phi_{n'} \left| \frac{\partial \phi_n}{\partial \alpha} \right. \right\rangle = \sum_{dd'} \frac{(T_{n'}^{d'})^\dagger T_n^d}{W_n - W_{n'}} \frac{\partial \langle f_{d'} | H | f_d \rangle}{\partial \alpha}$$

where $\langle f_{d'} | H | f_d \rangle$ are the matrix elements of the TRIM Hamiltonian.

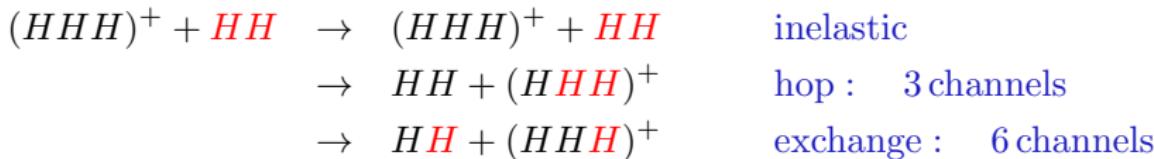
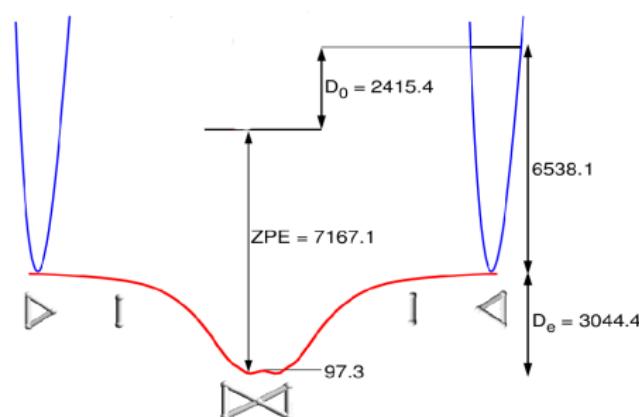
**Analytical: potential derivatives
non-adiabatic matrix elements
for H_3^+ , H_4^+ & H_5^+**

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$\text{H}_3^+ + \text{H}_2 \rightarrow \text{H}_2 + \text{H}_3^+$ exchange reaction

- Production of H_2D^+
 - Ortho/para conversion of H_3^+
 - Large Zero Point energy

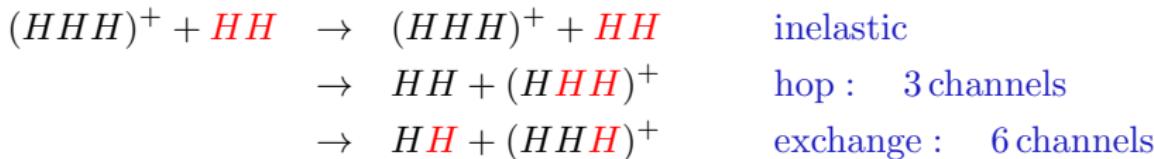
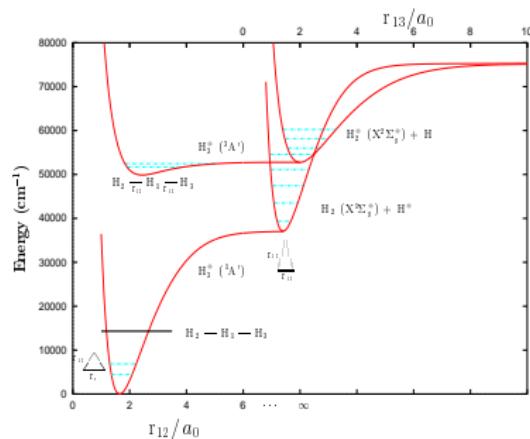


Each pathway obeys strict **nuclear spin** selection rules.
 This allows to “infer” the hop/exchange ratio: $\alpha = k^H/k^E$

Cordonnier et al. ('00)

Crabtree et al. ('11)

$\text{H}_3^+ + \text{H}_2 \rightarrow \text{H}_2 + \text{H}_3^+$ exchange reaction

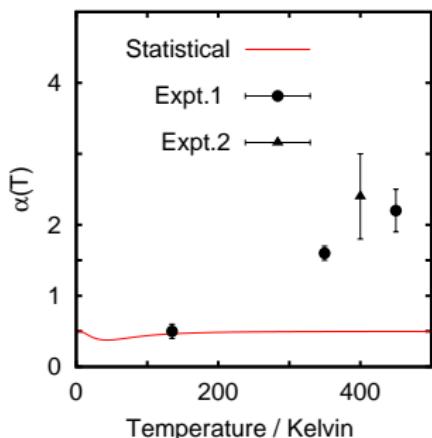


Each pathway obeys strict **nuclear spin** selection rules.
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Cordonnier et al. ('00)

Crabtree et al. ('11)

hop/exchange ratio, α , and statistical behaviour



α inferred from observed p-H₃⁺ fraction vs. p-H₂ enrichment

Expt. 1: Crabtree et al. JCP ('11)

Expt. 2: Cordonnier et al. JCP ('00)

H_5^+ complex lives enough to allow full scrambling

Statistical limits:

Parker & Light ('07),

Hugo *et al.* ('09)

- no scrambling $S_M = \{1/4, 3/4, 0\}$ $\alpha = \infty$
 - full scrambling $S_M = \{1/10, 3/10, 6/10\}$ $\alpha = 0.5$

Need of dynamics to describe transition

Beyond Pure Statistical Models

- Need of dynamics to describe α transition
 - Many degree of freedom (10D)
 - Low energy and long-lived resonances
 - All atoms are equivalent: NO reduced dimension models
 - Quantum and many semiclassical methods difficult
 - Introduce dynamical Bias through determination of S_M using Quassi-classical trajectories

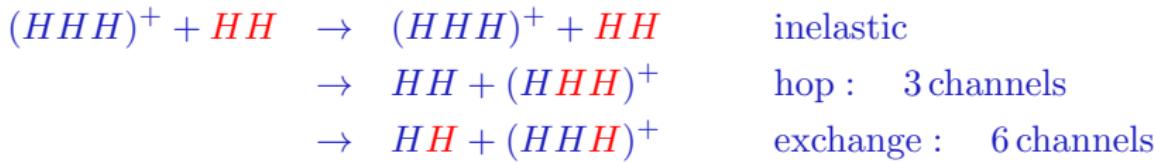
González-Carrasco et al. JCP 137, (12), 094303

$\text{H}_2 + \text{H}_3^+$: Statistical model of Park & Light, JCP ('07)

For total angular momentum, J , and nuclear spin, I , the reaction probability

$$P_{sr,M's'r'}^{J\Omega}(E) = \frac{g_{Is} \gamma_{sIs'}^{M'} W_{sr\nu\Omega_t}^{J\Omega}(E) W_{s'r'\nu'\Omega_t'}^{J\Omega}(E)}{\sum'' \gamma_{sIs''}^{M''} W_{s''r''\nu''\Omega_t''}^{J\Omega}(E)}$$

1. Nuclear spin statistical weight matrix, g
 2. Capture probability, $W_{sr\nu\Omega_t}^{J\Omega}$
 3. Spin branching ratio matrices, $\gamma_{sIs'}^{M'}$
 4. Scrambling matrix, $S_M = \{1/10, 3/10, 6/10\}$

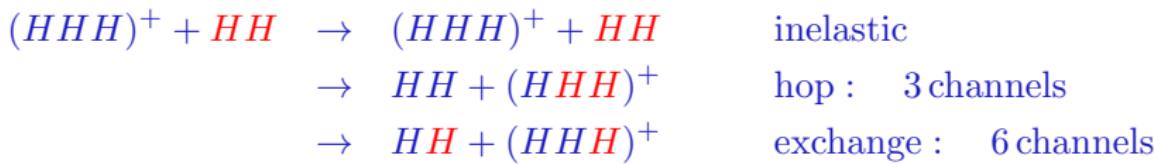


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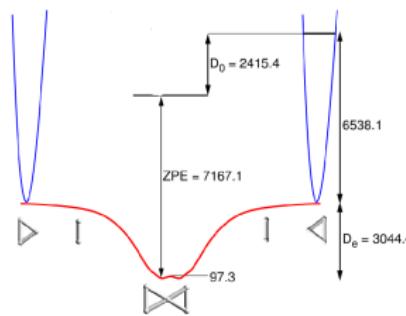
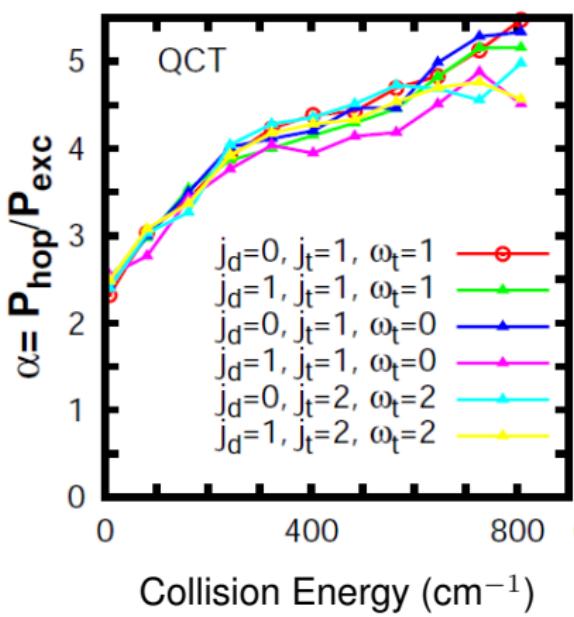
$$P_{sr,M's'r'}^{J\Omega}(E) = \frac{g_{Is} \gamma_{sIs'}^{M'} W_{sr\nu\Omega_t}^{J\Omega}(E) W_{s'r'\nu'\Omega_t'}^{J\Omega}(E)}{\sum'' \gamma_{sIs''}^{M''} W_{s''r''\nu''\Omega_t''}^{J\Omega}(E)}$$

1. Nuclear spin statistical weight matrix, g
 2. Capture probability, $W_{sr\nu\Omega_t}^{J\Omega}$
 3. Spin branching ratio matrices, $\gamma_{sIs'}^{M'}$
 4. Scrambling matrix, S_M from dynamical calculations



QCT Dynamically biased S_M and ZPE

$$S_M(E) = \text{QCT prob.}$$

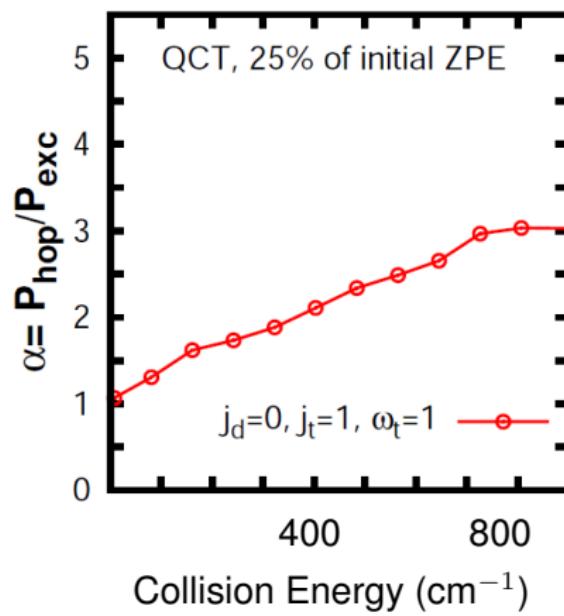
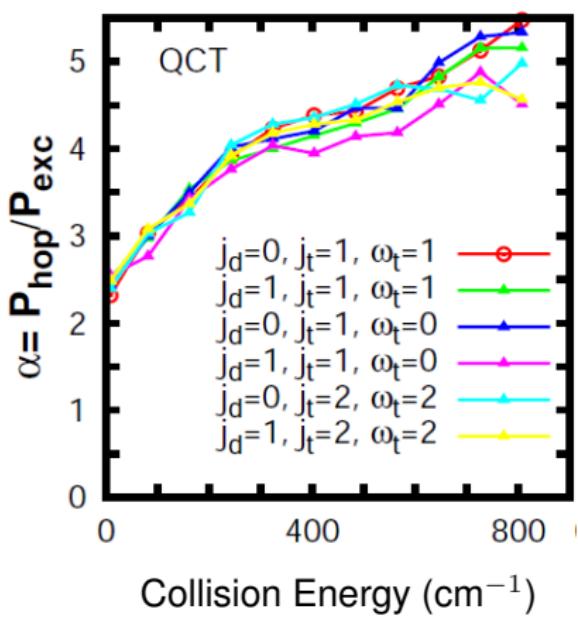


- α too high!!
 - ZPE favors direct hop mechanism

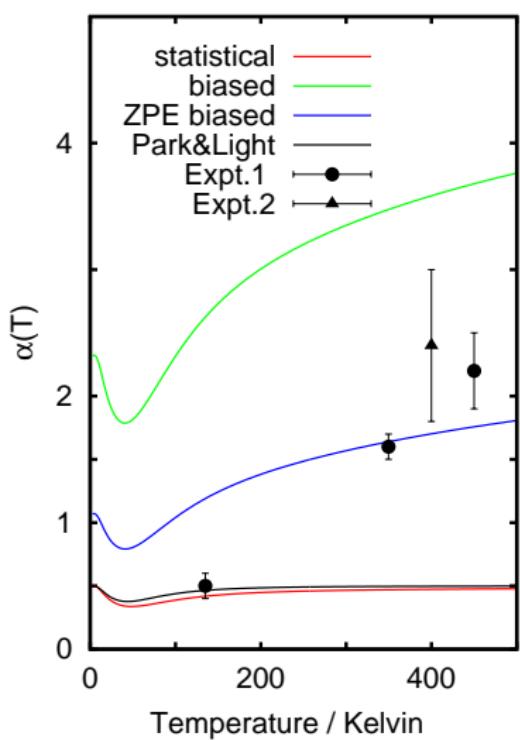
QCT Dynamically biased S_M and ZPE

$$S_M(E) = \text{QCT prob.} \quad \text{ZPE}_{H_2} + \text{ZPE}_{H_2^+} = 6538 \approx \text{ZPE}_{H_2^+} = 7167 \text{ cm}^{-1}$$

ZPE reduction: Manikandan & Hase ('12)



Comparison with Experiments



Expt. 1: Crabtree, *et al.*, J. Chem. Phys., 134, (11)
Expt. 2: Cordonnier, *et al.*, J. Chem. Phys., 113, (00)
Simul. Gómez-Carrasco *et al* J. Chem. Phys., (12)

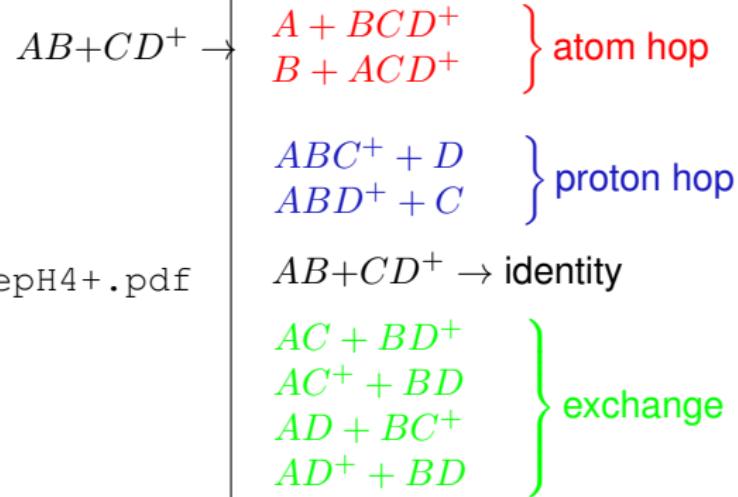
- ZPE plays a fundamental role
 - ZPE-biased “works” at $T > 300$ K
 - Statistical behaviour at $T < 200$ K
 - Need of quantum treatments
 - Direct experimental information in infrared predissociation spectra

Experiments by Cheng et al. JPCI ('10)

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- 1 Introduction
 - 2 $\text{H}_2 + \text{A}^+$ & $\text{HA}^+ + \text{H}$
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 - 6 Conclusions

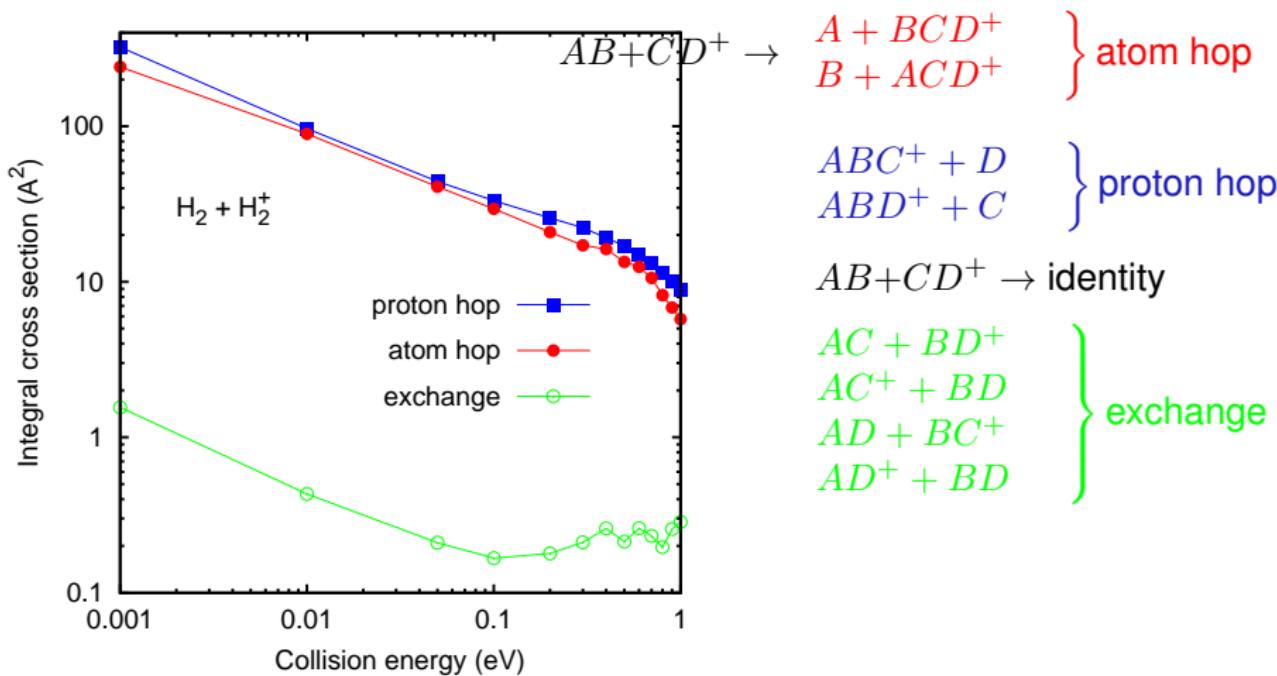
Direct $H_2 + H_2 \rightarrow H + H_3^+$ reaction



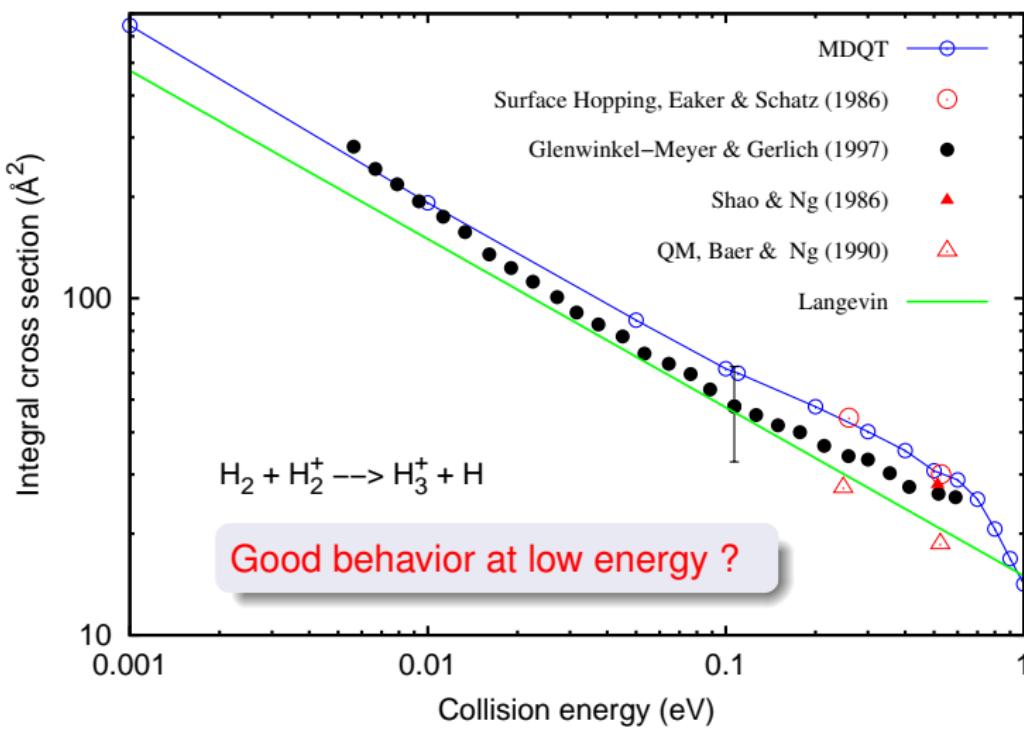
.../.../.../figuras/H4+/mepH4+.pdf

Direct $\text{H}_2 + \text{H}_2^+ \rightarrow \text{H} + \text{H}_3^+$ reaction

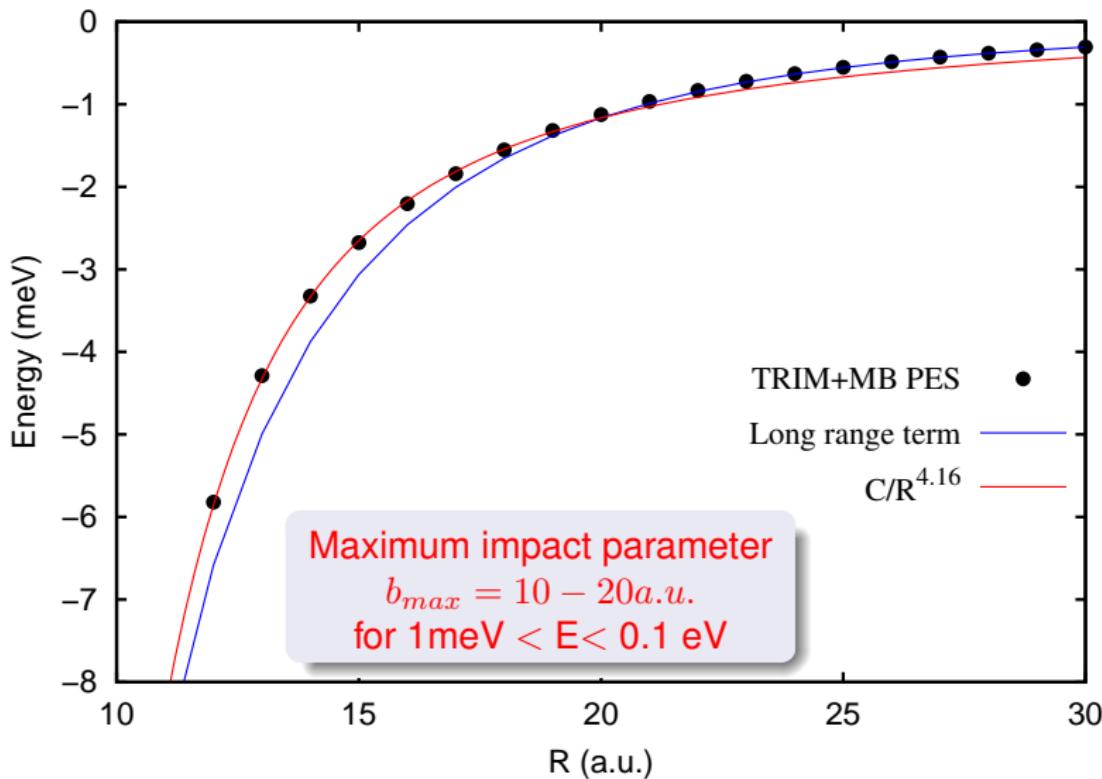
Total reaction cross section using QCT



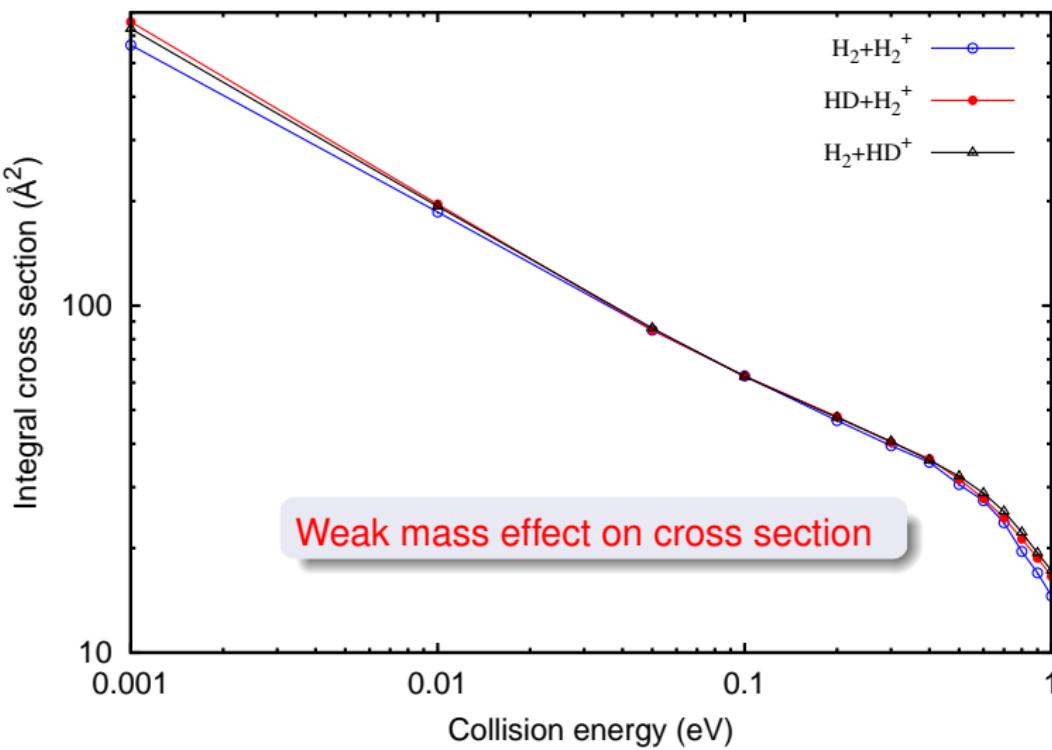
Comparison with experiments



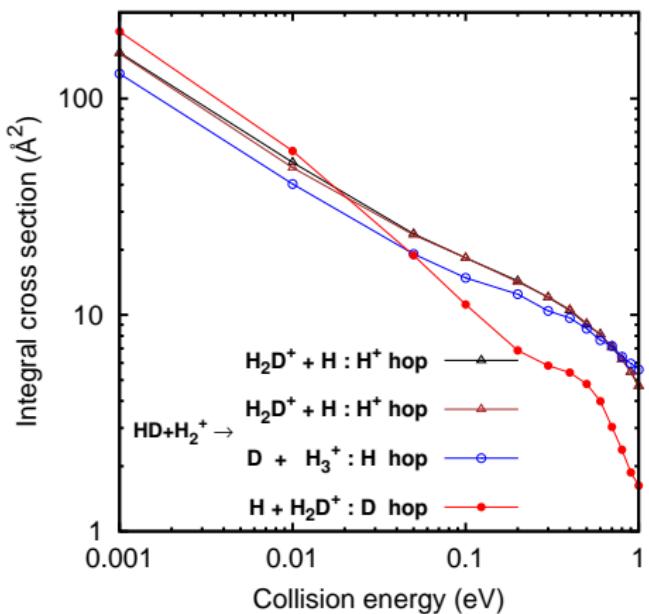
Long range behavior



Deuterium effect on cross sections



$\text{H}_2\text{D}^+ / \text{H}_3^+$ ratio: D or H hop?



- More reaction channels $\rightarrow \text{H}_2\text{D}^+$
 - H_2^+ breaks more easily
 - Variation of H-hop vs. D-hop

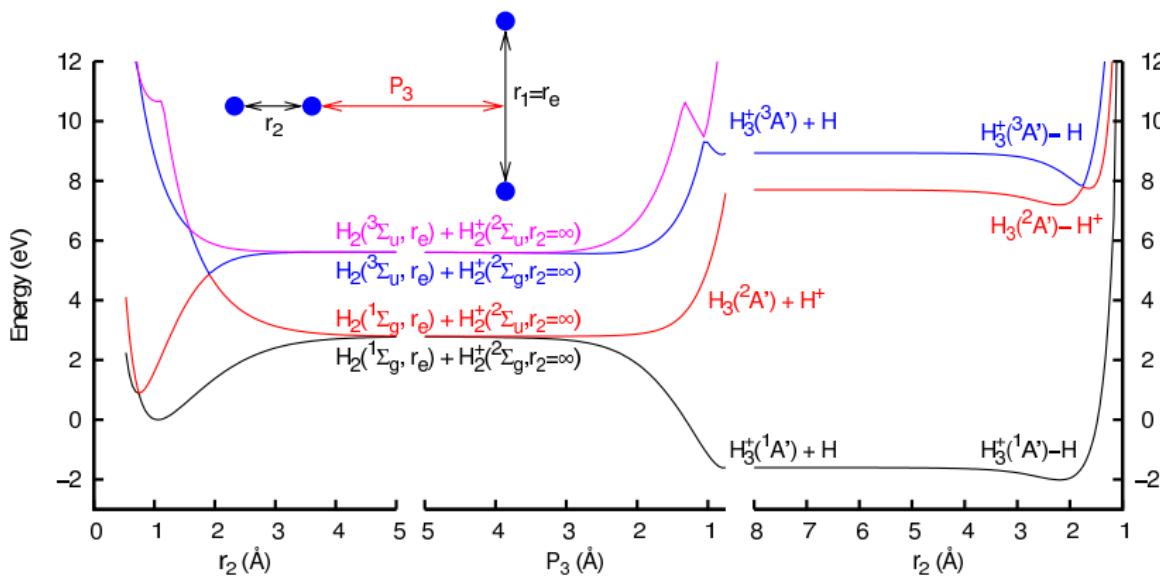
density of states:
 $\rho(\text{H}_2\text{D}^+) > \rho(\text{H}_3^+)$

Quantum effects ?

 - Nuclear spin and symmetry effects ?

Non-adiabatic vs charge transfer processes

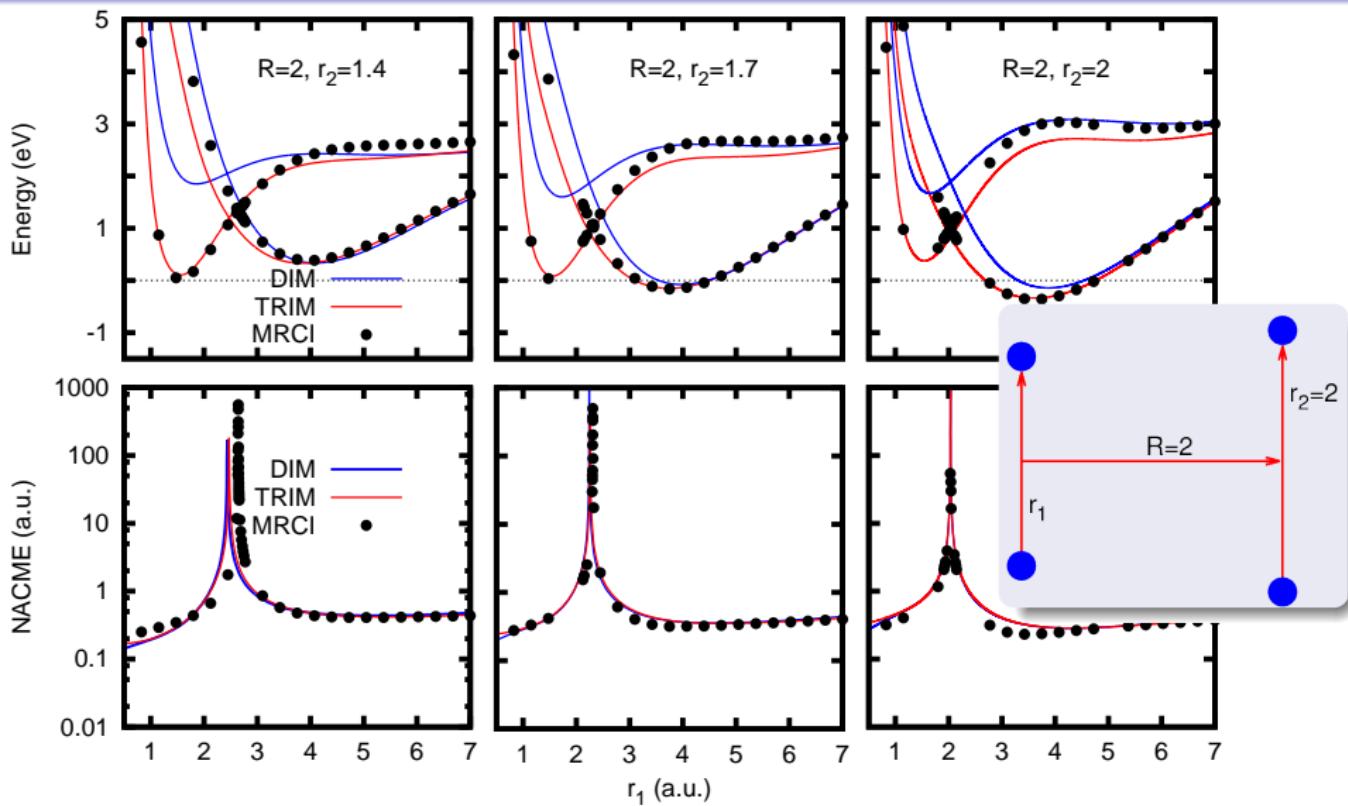
Crossings: when hydrogen atoms are at equivalent positions



Sanz-Sanz et al. JCP ('13)

Stein & Muckerman, JCP, 68 ('78)

Non-adiabatic couplings & Conical intersections

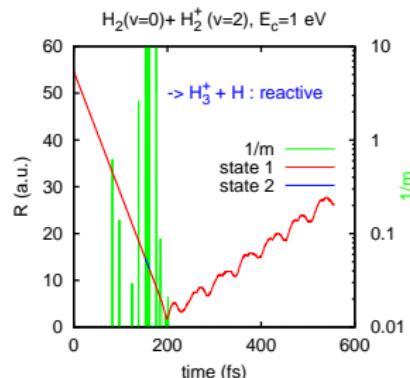
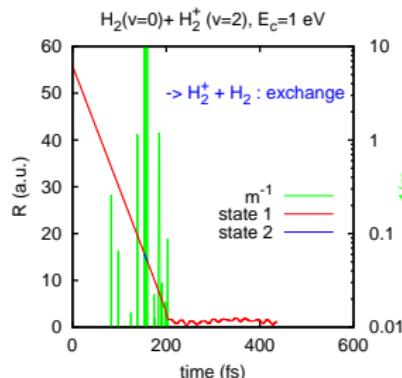
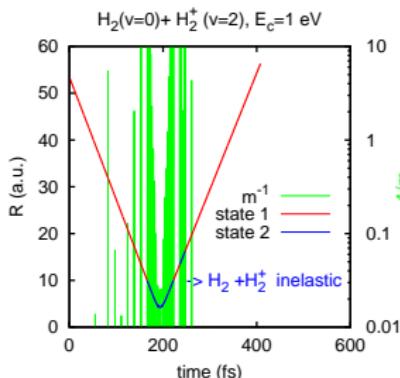


Non-adiabatic dynamics

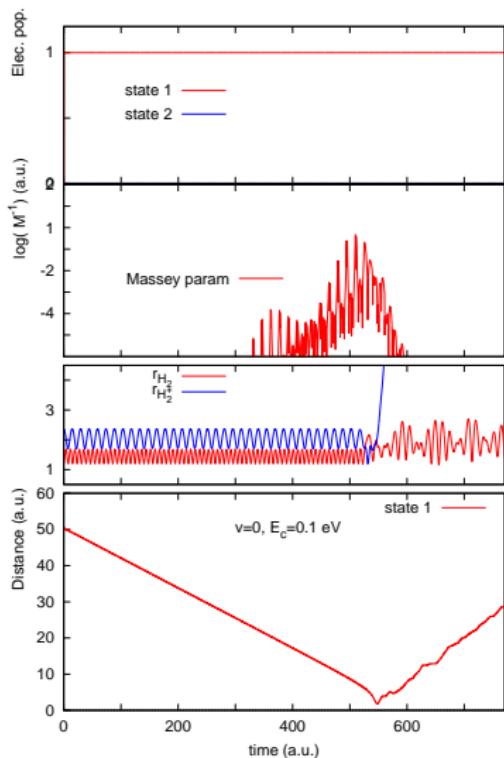
Using the “Molecular dynamic with quantum jumps” method (Tully '90)
with 2 adiabatic states

Transitions proportional to the inverse of Masey parameter:

$$m_{ij}^{-1} = \frac{\vec{v} \cdot \vec{\nabla}_{ij}}{2(W_i - W_j)}$$



Surface hopping method (Tully)



- #### - Electronic Quantum Wave packet

$$\Phi(t) = a_1(t)|1\rangle + a_2(t)|2\rangle$$

Split Operator propagator

- “Electronic jumps” probability

Proportional $\partial a_i(t) / \partial t$

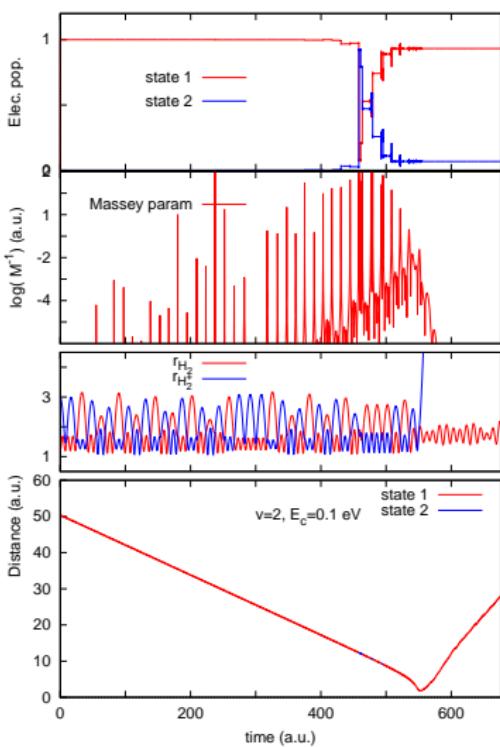
- Classical trajectory on $|1\rangle$ or $|2\rangle$

Energy change along ∇_{ij}

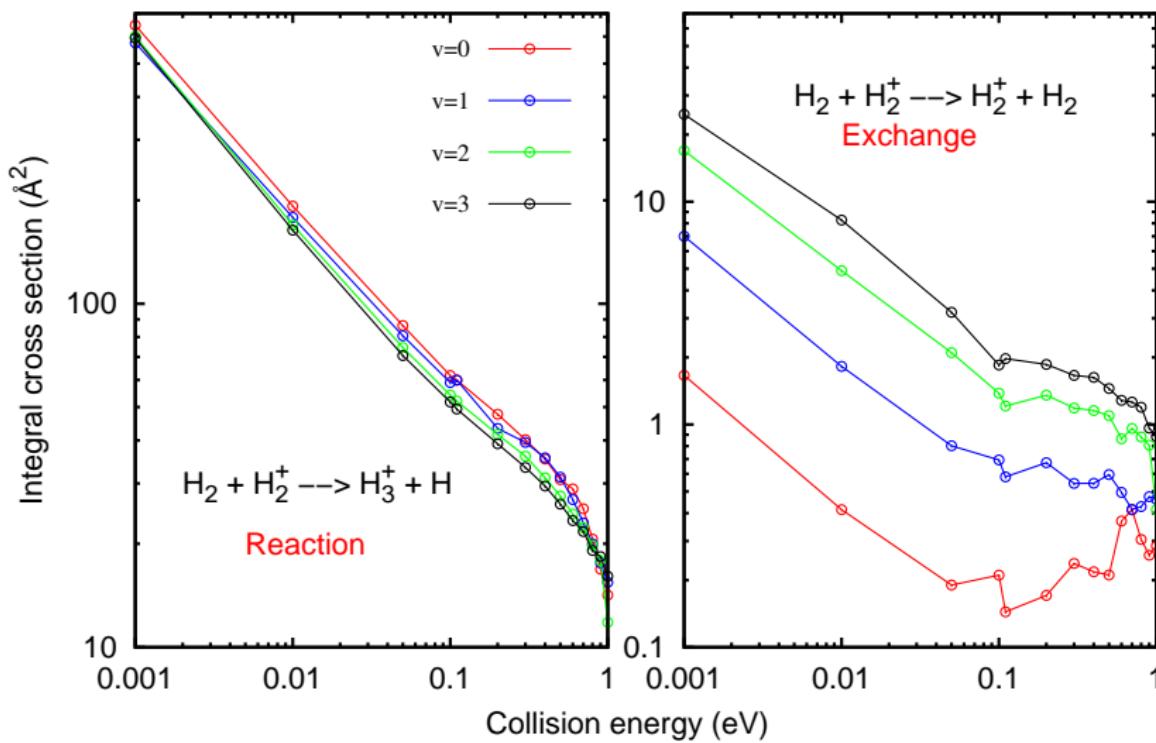
Example v=0 and 2 for $E_c=0.1$ eV

Surface hopping method (Tully)

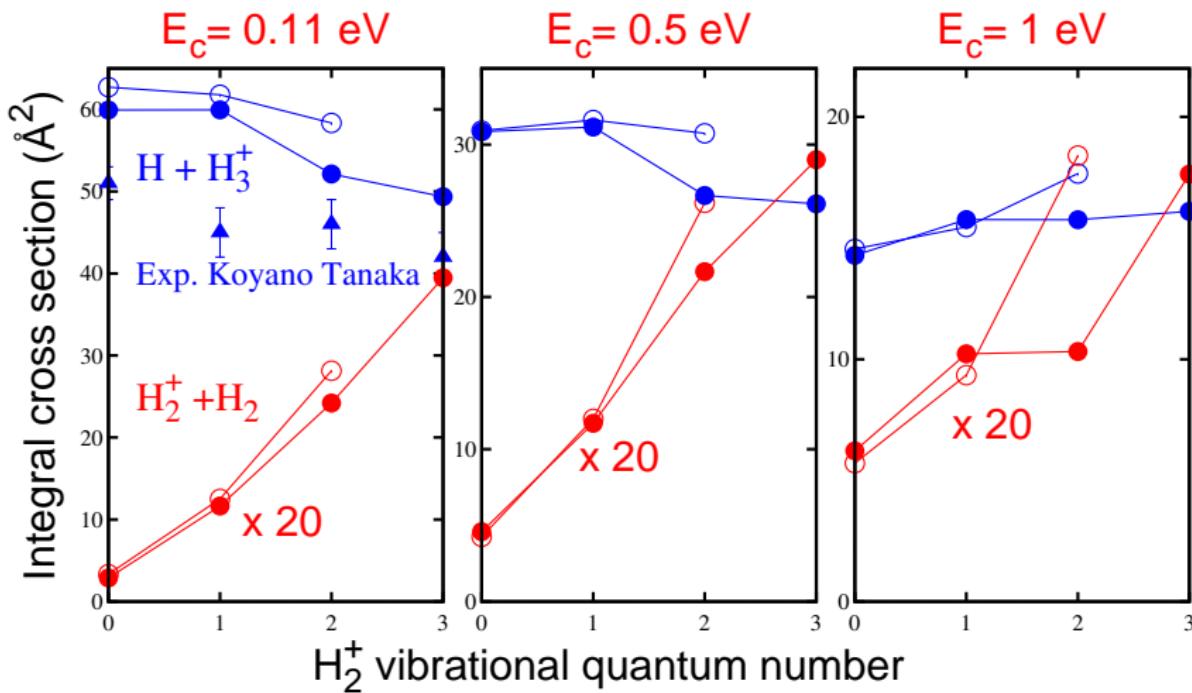
Is adiabatic approach good?



Initial vibration effect



Initial vibration effect

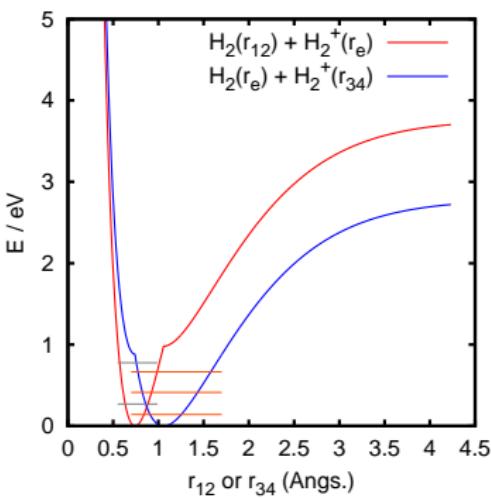
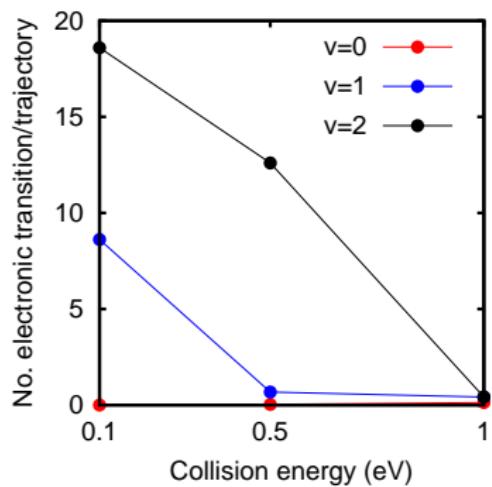


Triangles: Experiments at 0.11 eV by Koyano & Tanaka, JCP('80)

open circles; QCT

full circles: Surface hopping with 2 states

Collision energy effect



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Conclusions ?

- Accurate analytical PES of H_4^+ and H_5^+ based on a TRIM description, with analytic derivatives and NACME
 - $H_2 + H_3^+$ collisions presents a transition from statistical behaviour at $T < 200K$ and a more direct mechanism at $T > 300 K$.
Need of Quantum methods to study this transition & to understand the statistical behaviour at low temperatures, due to long-lived resonances
 - $H_2 + H_2^+(v = 0) \rightarrow H_3^+ + H$ is a direct reaction, well described on a single surface for $E_{col} < 1 \text{ eV}$.
 - For $v > 0$, non-adiabatic processes are important, specially at low E_c
 - The TRIM based PES describes accurately H_3^+ fragments: need of quantum methods and nuclear spin statistics to describe final H_3^+ rovibrational states

Acknowledgments

Collaborators

Cristina Sanz-Sanz, CSIC & Univ. Autónoma Madrid

Alfredo Aquado, Univ. Autónoma Madrid

Susana Gómez-Carrasco - Univ. Salamanca

Fedor Y. Naumkin, Univ. Toronto

Financial Support

Ministerio Ciencia y Tecnología (SPAIN)

CSD2009-00038, "Molecular Astrophysics: the Herschel and Alma era"

FIS2011-29596-C02