Lanthanide-based luminescent molecular thermometers

Carlos D. S. Brites, Patricia P. Lima, Nuno J. O. Silva, Angel Millán, Vitor S. Amaral, Fernando Palacio, and Luís D. Carlos

Received (in XXX, XXX) Xth XXXXXXXXXX 200X, Accepted Xth XXXXXXXXXX 200X
First published on the web Xth XXXXXXXXXX 200X
DOI: 10.1039/b000000x

Non-invasive accurate thermometers with high spatial resolution and operating at sub-micron scales, where the conventional methods are ineffective, is currently a very active field of research strongly stimulated in the last couple of years by the challenging demands of nanotechnology and biomedicine. This review offers a general overview of recent examples of accurate luminescent thermometers working at micrometric and nanometric scales, particularly those involving advanced Ln³⁺-based functional organic–inorganic hybrid materials.

Introduction

Sensing and mapping temperature in an accurate and non-invasive way with submicrometric resolution is critical to understand numerous features of micro and nanoscale electronic and photonic devices, such as thermal transport, heat dissipation, and profiles of heat transfer and thermal reactions. Furthermore, the precise discrimination of the temperature of a living cell, especially that of a cancer cell, strongly impacts on the consciousness of its pathology and physiology and, in turn, on the optimization of therapeutic processes (e.g. in hyperthermal tumour treatment and photodynamic therapy).

Conventional temperature sensors are the so-called contact thermometers, in which the measurement of temperature is achieved by heat flow to an invasive probe. The contact nature of such measurements makes traditional liquid-filled and bimetallic thermometers, thermocouples and thermistors not suitable for temperature measurements of fast-moving objects or in scales below 10 μm. Therefore, a new generation of nanoscaled thermometers is demanded to accurately measure temperature distributions down to the nanoscale regime, where conventional methods are not able to give satisfactory results (e.g. intracellular temperature fluctuations), temperature at molecular scale, or in microcircuits and microfluids.

The first approaches to nanothermometry tried to miniaturize the geometrical size of the conventional thermometers. For example, filled system thermometers are fabricated from nanotubes based on temperature-dependent thermal expansion of liquids and techniques using an atomic force microscope have succeeded in achieving spatial and temperature resolutions down to ~50 nm and ~1 mK, respectively. Whereas nanotube-based systems require a transmission electron microscope for the read-out and calibration and high vacuum conditions, a marked limitation for practical applications, scanning microscopy systems have a read-out rate limited by material surface conditions and probe motion (typically 30 ms per pixel) making real-time temperature mapping unfeasible.

These constraints lead to the development of non-contact thermometry techniques, such as, IR thermography, Raman spectroscopy, thermoreflectance and luminescence. Amongst non-invasive or semi-invasive spectroscopic methods for determining temperature, the temperature dependence of phosphor luminescence (band shape, peak energy, intensity and/or lifetimes) is an accurate alternative technique, working remotely, through an optical detection system, even in biological fluids, strong electromagnetic fields and fast-moving objects. This technique (often referred as thermographic phosphor thermometry) offers high detection sensitivity and spatial resolution in short acquisition times. Thermal probes, such as organic dyes, polymers, semiconductor nanocrystals and trivalent lanthanides are incorporated into inorganic, organic or organic/inorganic hybrid hosts and most of the luminescent thermometers work exploiting the drop of their emission intensity with increasing temperature, due to the thermal activation of nonradiative deactivation pathways. A pioneering example is the work of O. Zohar et al. on the thermal imaging of metabolic heat signals generated by ligand-receptor interactions in single cells using the temperature dependent luminescence intensity of a Eu³⁺-β-diketonate chelate.

The present review describes current progress of non-invasive (or semi-invasive) Ln³⁺-containing luminescent molecular thermometers. Particular emphasis is placed on the few examples that effectively illustrate the temperature sensing/mapping at the sub-micrometric scale: (i) a scanning thermal microscope with Er³⁺/Yb³⁺ co-doped fluoride glass co-doped with Eu³⁺ and Yb³⁺ NPs, (ii) 2D distribution maps using Er³⁺-Yb³⁺ Er³⁺/Yb³⁺ NPs, (iii) Yb³⁺-based organic–inorganic hybrid NPs incorporating a Eu³⁺ tris(β-diketonate) complex, and (iv) magnetic hybrid NPs co-doped with Eu³⁺ and Tb³⁺ tris(β-diketonate) chelates.

Examples of optical micro and nanothermometers

Temperature sensing with thermographic phosphors is commonly classified in decay time and intensity ratio algorithms. Decay time methods exploit the temperature dependence of the lifetime of an emitting level and the intensity ratio algorithms use directly the intensity of one or more transitions to detect temperature. The intensity of each transition is proportional to the total number of atoms (population) in a given excited state at temperature T:

\[
I \propto g A \hbar \nu \exp \left( -\frac{E}{k_B T} \right)
\]
where \( g \) is the degeneracy of the state, \( A \) is the spontaneous emission rate, \( v \) is the frequency, \( h \) and \( k_B \) are the Planck and the Boltzmann constants, respectively, and \( E \) is the energy of the level. Frequently two transitions assigned to the same phosphor are used and the ratio between their fluorescent intensities (\( I_1 \) and \( I_2 \)) taken as a measurement of absolute temperature, since:

\[
\frac{I_1}{I_2} = \frac{B_1 A_1 h v_1}{B_2 A_2 h v_2} \exp \left( -\frac{\Delta E_{1g}}{k_B T} \right) = B \times \exp \left( -\frac{\Delta E_{1g}}{k_B T} \right),
\]

where \( B \) is a constant. The main advantage of using the ratio between the fluorescent intensities of two transitions as a measurement of absolute temperature, a numerical technique known as fluorescence intensity ratio (FIR), is the fact that one single emission spectrum contains all the information needed to compute the absolute temperature. The FIR ratio (sometimes also called two-colour response) is based on the temperature dependent Boltzmann distribution of electrons between adjacent emitting energy levels and manifests as a temperature-dependent intensity ratio between emissions from two distinct spectral lines.

The FIR algorithm (as well as the decay time data) makes the fluorescence signal easier to acquire and analyse, thus overcoming the main drawbacks of the intensity-based measurements of only one transition.\(^{17,18}\)

### Scanning thermal microscopes

The scanning thermal microscope developed by Aigouy and collaborators\(^{15,33,36}\) can measure the heating of electrically excited stripes, micro and nanowires (i.e. \( 1 \) \( \mu \)m-wide, \( 40 \) \( \mu \)m-long and \( 40 \) nm thick nickel stripe) by gluing \( \text{Er}^{3+}/\text{Yb}^{3+} \) co-doped fluoride amorphous glass particles or \( \text{PbF}_2: \text{Er}^{3+}/\text{Yb}^{3+} \) nanocrystals at the extremity of an atomic force microscopy scanning tip (Fig. 1). The \( \text{Yb}^{3+} \) ions were excited by a \( 975 \) nm near infrared (NIR) laser, through an NIR-to-visible up-conversion (UC) process in which two low energy photons are “added up” to give one high energy photon.\(^{37}\) This NIR excitation is very efficient because, at this wavelength, the \( \text{Yb}^{3+} \) ions have a large absorption cross section. Once excited, they transfer their energy to adjacent \( \text{Er}^{3+} \) ions that will glow in the green (\( \text{^5}H_{11/2} \rightarrow \text{^4}I_{15/2} \) and \( \text{^4}S_{3/2} \rightarrow \text{^4}I_{15/2}, \) ca. \( 525 \) and \( 550 \) nm, respectively) and red (\( \text{^4}F_{9/2} \rightarrow \text{^4}I_{15/2}, \) ca. \( 660 \) nm) spectral regions. These \( \text{Er}^{3+} \) transitions (that are insensitive to photodegradation) are temperature dependent and their intensities were used to determine the temperature using the FIR algorithm.

By adjusting the electrical current flowing through the structure, the resulting temperature variations modulate the \( \text{Er}^{3+} \) particle fluorescence giving rise to the thermal contrast. Although the fluorescence is affected both by the near-field optical distribution and by temperature variations, this thermal contribution can be treated separately by comparing the obtained images with reference ones recorded when the current in the device is zero. The determination of the temperature is, therefore, performed by analyzing the thermal quenching of the \( \text{Er}^{3+} \) fluorescence. The results obtained are in good agreement with numerical simulations of the heat dissipation of the micro and nanowires. Since the fluorescent probe is not in direct contact with the micro or nanowires, the temperature determined by this method is actually the average temperature of the fluorescent probe that is controlled by the temperature of the interaction volume between the NP and the surface of the wires.

For the scanning thermal microscope using \( \text{PbF}_2: \text{Er}^{3+}/\text{Yb}^{3+} \) NPs as fluorescence probes\(^{15}\) the relative sensitivity is \( 1.1\%/\text{K} \) at \( \approx 310 \) K (Table 1). This thermometer presents a spatial resolution (defined as the minimum distance to move in order to get a temperature change greater than the sensitivity of the thermometer) in the range of the fluorescent particle size (< \( 500 \) nm). For the scanning thermal microscope with the \( \text{Er}^{3+}/\text{Yb}^{3+} \) co-doped fluoride amorphous glass particles,\(^{33}\) the relative sensitivity changes between \( 0.37 \) and \( 0.59\%/\text{K} \), in the temperature range \( 297-390 \) K (Table 1). While in this latter scanning thermal microscope the quenching of luminescence with temperature is calculated directly from the temperature dependence of the intensity of the two \( \text{^3}H_{11/2} \rightarrow \text{^4}I_{15/2} \) and \( \text{^4}S_{3/2} \rightarrow \text{^4}I_{15/2} \) transitions,\(^{33}\) in the former example,\(^{15}\) the quenching value is extracted from the normalized thermally modulated fluorescence images.

Tikhomirov et al. used an analogous UC mechanism to fabricate an optical nanoheater of fluoride \( \text{Yb}^{3+}/\text{Er}^{3+} \) co-doped NPs embedded into different host matrices.\(^{35}\) The NPs were excited by NIR radiation (laser diode at \( 975 \) nm) that is dissipated via phonon relaxation (up conversion process), resulting in a temperature increase, measured by FIR of transitions of \( \text{Er}^{3+} \) (\( \text{^3}H_{11/2} \rightarrow \text{^1}I_{15/2} \) and \( \text{^4}S_{3/2} \rightarrow \text{^1}I_{15/2} \)), until \( 800^\circ \text{C} \). The authors have suggested the use of the nanoheater for local hypothermal treatment of cells with simultaneous monitoring of temperature.

### Upconverting nanoparticles

The rise time response of the \( \text{^5}D_0 \) excited state in \( \text{Y}_2\text{O}_3: \text{Eu}^{3+} \) (3–4% mole concentration) phosphor was calibrated against a temperature range of 200–700 °C demonstrating 2D rise time thermal imaging over an area of \( \approx 500 \) \( \mu \)m\(^2 \) (diameter of the laser beam on the target).\(^{29}\) The uncertainty in the temperature measurements is relatively high, \( \approx 20\% \) at \( 673 \) K, although a decrease of the \( \text{Eu}^{3+} \) activator concentration could significantly lengthen the rise time and then the accuracy of the thermometer. The sensitivity and the spatial resolution were not reported and the temperature determination algorithm is rather complicated, relatively to those of decay time and intensity-based measurements, requiring more sophisticated equipment and post-processing computational analysis.\(^{29}\)

Other example that illustrates the sensing/mapping of temperature at the sub-micrometer scale was reported by Allison et al.\(^{25}\) The laser-induced fluorescence lifetime of YAG:Ce NPs (average size \( \approx 30 \) nm) was measured as a function of temperature from 280 to 350 K. The decay lifetimes for the NPs varied from 18 to 27 ns (\( 33\% \) relative to the longest lifetime measured) and this significant variation, coupled with the high signal strength that was observed, suggested the utilisation of YAG:Ce NPs as useful thermographic nanophosphors. From an instrumentation perspective, the short decay times of Ce-doped nanophosphors makes the fluorescence signal easier to acquire and analyse, relatively to phosphors with millisecond decay times (when roughly the same number of photons are emitted and detected in both cases), thus providing a noteworthy advantage. The reason
for this improvement is that the peak values of the signals are greater and the rise times are faster, thus reducing triggering errors in the detection electronics. Another attractive characteristic of Ce-doped nanoporphors is that they can typically be excited by inexpensive light emitting diodes, thus not only reducing the cost of the phosphor thermography system but also alleviating any concerns about heating of the phosphor substrate by the laser (with the resulting introduction of a systematic uncertainty in the temperature measurements). The upconverting Ln\(^{3+}\)-doped NPs can be stimulated using low power and inexpensive NIR lasers and are excited with wavelengths in the optical penetration window of cells and tissues (650–1300 nm), an important requirement in biological environments who suffer less damage under excitation in such range. Moreover, besides less scattered by the specimen, NIR excitation light will not induce background autofluorescence, leading to increased signal-to-noise ratio of the detected signal. Thermal profiles created when heating a colloidal solution of the NaYF\(_4\)Er,Yb NPs in water using a pump–probe experiment were obtained using the temperature dependence of the intensity ratio between the Er\(^{3+}\)-\(\text{H}_{11/2} \rightarrow \text{I}_{15/2}\) and \(\text{S}_{3/2} \rightarrow \text{I}_{15/2}\) transitions. The internalization of the NPs by human cervical carcinoma (HeLa) cells permits to exploit the thermal sensitivity of the two Er\(^{3+}\) bands to create a nanothermometer capable of measuring the internal temperature of a living cancer cell, from 298 K to its thermally induced death at 318 K. The cells were placed in a confocal fluorescence microscope and excited at 920 nm so that the upconverted Er\(^{3+}\) fluorescence of the NPs permitted the measurement of the inner HeLa cell temperature (Fig. 2). The NaYF\(_4\)Er\(^{3+}\)/Yb\(^{3+}\) NPs are one of the few cases reported so far of cellular thermometers. Despite the absence of a quantitative estimation of the thermometer sensitivity, the NPs are able to investigate (without using an exterior reference) the cellular changes occurring between 298 and 318 K, as a result of the external heating.

Organic-Inorganic hybrids

The last two examples illustrating the sensing/mapping of temperature at the nanoscale report the incorporation of Ln\(^{3+}\)-based \(\beta\)-diketonate complexes into organic-inorganic hybrid materials. Luminescent molecular thermometers based on Eu\(^{3+}\) tris(\(\beta\)-diketonate) complexes were proposed in the past decades involving isolated complexes or complexes embedded into polymer thin films and organic-inorganic hybrid materials (maximum sensitivities and temperature ranges listed in Table 1). The complexes act as temperature sensors through a thermally driven energy back transfer mechanism between the \(\text{D}_0\), Eu\(^{3+}\) emitting level and the lowest ligand triplet excited state. Then, its thermal sensitivity can be tuned by changing the ligands and/or the metal ion concentration. A serious drawback of Ln\(^{3+}\) \(\beta\)-diketonate complexes relates to their photodecomposition under UV irradiation (in certain cases, in only a few hours), which decreases luminescence intensity and, thus, makes molecular thermometers based on these complexes not suitable for long-term monitoring. However, the incorporation of the complexes into organic-inorganic hybrid hosts considerably improves the UV photostability, overcoming, therefore, the main drawback of the use of such complexes as long-term monitoring luminescent thermometers. Furthermore, the thermal stability and the mechanical properties of the resulting Ln\(^{3+}\)-containing organic-inorganic hybrids are considerably ameliorated, relatively to those of the parent complexes. This is an important advantage concerning the technological applicability of these materials, namely in lighting, sensing, and displays. Finally, the combination of the intrinsic characteristics of sol–gel derived hosts and the luminescence features of Ln\(^{3+}\) ions in multifunctional organic-inorganic hybrids offer excellent perspectives for designing new luminescent materials with enhanced desired characteristics, thus opening exciting new directions in materials science and related technologies, with noteworthy results in the ecofriendly integration, miniaturization, and multifunctionalization of devices.
bip(1-pyrenyl)propane dissolved in the 1-butyl-1-
allyl)propane bis(trifluoromethylsulfonyl)imide, 
[C$_{mpy}$][Tf$_2$N]), ionic liquid$^{10}$ for dye-doped materials, and the 
$^3_{1/2}H_{11/2} ightarrow ^1_{15/2}S_{3/2} ightarrow ^1_{15/2}$ intensity ratio, in Eu$^{3+}$/Yb$^{3+}$-based up-
converting NPs$^{5, 15, 30-32, 42, 43}$ and the $^4D_0 ightarrow ^2F_2^0/D_2^0 ightarrow ^2F_3$ intensity ratio in Eu$^{3+}$/Tb$^{3+}$-doped hybrid magnetic nanoclusters,$^{18}$ for 
Ln$^{3+}$-based materials.

This latter example consists of a luminescent molecular 
thermometer based on hybrid magnetic nanoclusters (size ranging 
from 100 to 400 nm) formed by a γ-Fe$_2$O$_3$ maghemite core 
(hydrodynamic size of 21.0±4.0 nm) coated with a tetraethyl 
ortho silicate/aminopropyltriethoxysilane (TEOS/APTES) 
organosilica shell and co-doped with [Eu(btfa)$_3$(MeOH)(bpea)] 
and [Tb(btfa)$_3$(MeOH)(bpea)] β-diketonate chelates [where btfa$, 
bpea and MeOH represent, respectively, 4,4,4-trifluoro-1-
phenyl-1,3-butadienedionate, 1,2-bis(4-pyridyl)ethane and 
methanol]. Fig. 3. The Eu$^{3+}$/Tb$^{3+}$ co-doped NPs were 
prepared using Eu:Tb ratios of 2:1, 1:1, 1:2, 1:3 and 1:10 
followed the procedure described elsewhere, supporting 
information of ref. 17.

The Eu$^{3+}$/Tb$^{3+}$ luminescent nanothermometer is self-
referencing, allowing absolute measurements in the 10-350 K 
temperature range with a temperature uncertainty of 0.5 degree. 
Its emission colour is fine-tuned as a function of temperature 
(Fig. 4) and the temperature sensitivity is up to 4.9%-K$^{-1}$, the 
highest value reported so far for Ln$^{3+}$-based thermometers. 
Moreover, it exhibits high photositivity in long-term use.

Highly luminescent water suspensions of the hybrid Eu$^{3+}$/Tb$^{3+}$ 
NPs at room temperature (Fig. 5) indicate that the quenching of 
the Ln$^{3+}$ emission by nonradiative energy transfer from the intra-
df/exited states to the high-energy OH-vibrations (~3700 cm$^{-1}$) is 
not significant. The Eu$^{3+}$ and Tb$^{3+}$ first coordination ligands in 
these NPs efficiently shield the metal centres from the OH 
oscillations permitting, therefore, that the nanothermometer 
operates in aqueous suspensions, a critical advantage for 
biomedical applications.

The variation of the Eu$^{3+}$/Tb$^{3+}$ ratio affords tunability to the 
temperature working range. Alternatively, tunability is also 
accomplished by changing the host matrix, thus modifying the 
interaction between the Ln$^{3+}$ and the host matrix energy levels.$^{18}$ 
This is the case of Eu$^{3+}$/Tb$^{3+}$ co-doped di-ureasils$^{40}$ that permitted 
absolute temperature measurements at physiological 
temperatures. Moreover, the di-ureasil hybrid matrix enables 
processing the thermometer material as thin films for 
sensing/mapping large areas with a spatial resolution limited by 
the size of the optical detectors (~1-10 µm for commercial optical 
fibres and CCD cameras). Finally, the combination of the 
molecular thermometer with a nanometric magnetic/luminescent 
host matrix provides multifunctionality at the nanoscale to the 
device.

When comparing this thermometer with the single Ln$^{3+}$-based 
one, proposed previously, it is clear that it represents a step 
forward in thermometry at the nanoscale. The synergy resulting 
from combining temperature sensing/mapping and 
superparamagnetism opens the way for new exciting applications, 
especially in the biomedical field. In particular, such association 
will provide a unique instrument to map, in a non-invasive way, 
temperature distributions in biological tissues (e.g., in tumours) 
during heat release, due to the application of an ac field to 
magnetic NPs (magnetic hyperthermia), this being, undoubtedly, 
a powerful tool for the study of the development of micro-processes 
occurring within a cell.

Conclusions

This review summarized the recent progress on non-invasive 
Ln$^{3+}$-containing luminescent molecular thermometers, with 
particular emphasis on the promising examples reported so far that 
effectively illustrate the temperature sensing/mapping at the sub-
micrometer scale: scanning thermal microscopes that uses emitter 
NPs glued at the end of the scanning tip as temperature probes; 
NPs (Y$_2$O$_3$:Eu$^{3+}$, YAG:Ce$^{3+}$ and NaYF$_6$:Er$^{3+}$/Yb$^{3+}$), and Eu$^{3+}$- or 
Eu$^{3+}$/Tb$^{3+}$-containing siloxane-based organic-inorganic hybrid 
NPs. The main challenges that are currently facing scientists in 
the field are: (i) synthesis of biocompatible and highly efficient 
luminescent up-converting NPs acting as thermometers capable 
of operate in the physiological temperature range with sub-
micron spatial resolution.; (ii) design of different host hybrid 
matrices and/or lanthanide ligands to make absolute 
measurements at the physiological temperature range feasible 
with hybrid NPs; (iii) internalization of the NPs (both the up-
converting and the hybrid ones) into cellular medium; (iv) 
quantitative characterization of the NP-to-surface heat transfer 
mechanisms, particularly for scanning thermal microscopes; and 
(v) combination and integration of different functionalities into 
more complex architectures, e.g. magnetic hyperthermia and the 
corresponding temperature monitoring. This sort of “lab-on-a-
particle” technology is clearly one of the next evolutionary steps 
in this research area with dramatic impact on information and 
communication technologies, global health and biomedicine.

Acknowledgements

The authors are grateful to Fundação para a Ciência e a 
Tecnologia (FCT, Portugal), COMPETE and FEDER 
programs (PTDC/CTM/101324/2008) and Integrated Spanish-
Portuguese Action PT2009-0131 for financial support. The 
work in Zaragoza has been supported by the grants 
MAT2007-61621 and CONSOLIDER CSD2007-00010 from the 
Ministry of Education. CD$^2$B (SFRH/BPD/38472/2007) 
and PPL (SFRH/BPD/34365/2006) thank FCT for grants. 
NJOS acknowledges FCT for Ciência 2008 program.

Notes and references

100 $^a$ Department of Physics, CICECO, University of Aveiro, Campus 
Universitário de Santiago, 3810–193 Aveiro, Portugal. Fax: 351 234 
378197; Tel: 351 234 370946; E-mail: lcarlos@ua.pt

$^b$ Departamento de Física de la Materia Condensada, Facultad de 
Ciencias y Ingeniería Química, Universidad de Zaragoza, 50009 Zaragoza, Spain.

2. S. Sadat, A. Tan, Y. J. Chua and P. Reddy, Nano Lett., 2010, 10, 
2613-2617.
Elbaum, D. L. Alkon and T. Yoshioka, Biophys. J., 1998, 74, 82-
89.
**Fig. 1** Example of the experimental setup used in scanning thermal microscopes. The device is powered by a square electrical current (brown curve) that induces a periodic Joule heating (red curve) and a fluorescence quenching of the particle (black curve). The inset is a scanning electron microscope image of the nanocrystal glued at the end of the tip. (Reproduced with permission from ref. 15, copyright 2009 AIP).

**Fig. 2** (Top) Optical transmission images of an individual HeLa cell at three inner temperatures. Cell death is observed at 45 °C. (Bottom) Temperature of the HeLa cell determined by Er$^{3+}$ luminescence in the NaYF$_4$:Er,Yb NPs as a function of the applied voltage. (Reproduced with permission from ref. 5, copyright 2010 ACS).

**Fig. 3** (a) Optical microscope image at 300 K (UV excitation at 365 nm) of NP3.1-2 NPs. The images were recorded in a Carl Zeiss Axiovert 200 M fluorescence microscope equipped with Apotome at 125 and 300 K. The red colour corresponds to real colour. (b) SEM image with a general view of the NPs. (c) Amplification TEM image of the NPs showing aggregation occurred probably during the evaporation of the dispersion on the grid.

**Fig. 4** (Top) Optical microscope images at 125 and 300 K (UV excitation at 365 nm) of NP3.1-2 NPs. The colours correspond to real colours. (Bottom) Emission spectra (excited at 365 nm) of the NPs recorded at 125 and 300 K.

**Fig. 5** Photograph of an water suspension (1 mL) of the NP3.1-2 NPs under (a) day light illumination and (b) UV excitation (365 nm) throughout an optical fibre bundle.
Table 1 – Sensitivity and temperature operation range of Ln$^{3+}$-based luminescent molecular thermometers.

<table>
<thead>
<tr>
<th>Phosphor</th>
<th>Maximum Sensitivity (% Intensity / K)</th>
<th>Temperature range (K)</th>
<th>Temperature of maximum sensitivity (K)</th>
<th>Temperature dependence algorithm</th>
<th>Observations</th>
</tr>
</thead>
<tbody>
<tr>
<td>BaTiO$_3$:Er NPs$^{43}$</td>
<td>0.5 (a)</td>
<td>333-466</td>
<td>333</td>
<td>$^4H_{11/2} \rightarrow ^4I_{15/2}$/$^4S_{3/2} \rightarrow ^4I_{15/2}$ intensity ratio</td>
<td>Effect of the size of the NP on sensitivity discussed</td>
</tr>
<tr>
<td>ZnO:Er NPs$^{42}$</td>
<td>0.6 (a)</td>
<td>283-475</td>
<td>~443</td>
<td>$^4H_{11/2} \rightarrow ^4I_{15/2}$/$^4S_{3/2} \rightarrow ^4I_{15/2}$ intensity ratio</td>
<td>-</td>
</tr>
<tr>
<td>Gd$_2$O$_3$:Er/Yb$^{30, 31}$</td>
<td>0.4 (b)</td>
<td>300-900</td>
<td>~300</td>
<td>$^4H_{11/2} \rightarrow ^4I_{15/2}$/$^4S_{3/2} \rightarrow ^4I_{15/2}$ intensity ratio</td>
<td>-</td>
</tr>
<tr>
<td>Er/Yb co-doped PbF$_2$ NPs$^{15}$</td>
<td>1.1 (a)</td>
<td>290-325</td>
<td>~310</td>
<td>$^4H_{11/2} \rightarrow ^4I_{15/2}$/$^4S_{3/2} \rightarrow ^4I_{15/2}$ intensity ratio</td>
<td>Accuracy of 0.8 K. For the Er$^{3+}$/Yb$^{3+}$ co-doped fluoride amorphous glass NPs$^{33}$ the relative sensitivity changes between 0.37 and 0.59%/K</td>
</tr>
<tr>
<td>Eu β-diketonate complex$^{28}$</td>
<td>~2.0 (a)</td>
<td>280-330</td>
<td>~283</td>
<td>$^5D_0 \rightarrow ^7F_2$ intensity</td>
<td></td>
</tr>
<tr>
<td>Eu/Tb hybrid NPs$^{38}$</td>
<td>4.7 (a)</td>
<td>10-350</td>
<td>134</td>
<td>$^5D_0 \rightarrow ^7F_2$/$^5D_4 \rightarrow ^7F_2$ intensity ratio</td>
<td>The Eu/Tb-containing di-ureasil film presents a sensitivity of 0.6%/K from 298 to 318 K</td>
</tr>
<tr>
<td>Eu hybrid NPs$^{17}$</td>
<td>3.1 (a)</td>
<td>283-323</td>
<td>~283</td>
<td>$^5D_0 \rightarrow ^7F_2$ intensity and $^5D_0$ lifetime</td>
<td>Temperature resolution of ~0.3 K, assuming a precision of 1% in emission intensity</td>
</tr>
<tr>
<td>Eu and Tb complexes$^{34}$</td>
<td>1.2 (a)</td>
<td>295-335</td>
<td>-</td>
<td>$^7D_2 \rightarrow ^7F_3$ (Tb$^{3+}$) and $^5D_0 \rightarrow ^7F_2$ (Eu$^{3+}$) intensities; $^5D_4$ and $^7D_0$ lifetimes</td>
<td>The maximum sensitivity is reported for the Tb complex in a solution of dimethyl sulfoxide. Maximum $^7D_1$ lifetime sensitivity of 13.8μs/degree (3 wt% in PMMA); Pressure sensor for the Eu complex</td>
</tr>
<tr>
<td>Eu β-diketonate and Pt porphyrin type complexes$^{34, 44}$</td>
<td>0.9-4.4 (b)</td>
<td>278-323</td>
<td>~278</td>
<td>$^5D_0 \rightarrow ^7F_2$ intensity and $^5D_0$ lifetime</td>
<td>The two luminophors were jointly incorporated into a polymer single paint allowing temperature and pressure measurements</td>
</tr>
</tbody>
</table>

(a) Sensitivity defined as $\partial Q/\partial T$, where $Q$ corresponds to the quenching of luminescence with temperature.  
(b) Relative sensitivity defined as $\partial \ln(Q)/\partial T$
Figure 1
Figure 2
Figure 4

(a) 125 K

(b) 300 K